### Study of N\* in $\chi QM$ via $\eta$ productions

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 $N^*$  in  $\eta$  production Extracting information of resonances

#### Why we study the $\eta$ production

• N\*: Nonperturbative QCD

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 $N^*$  in  $\eta$  production Extracting information of resonances from data Reach the subnucleonic degrees of freedom

### Why we study the $\eta$ production

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- The problems about N\* spectrum compared with CQM.



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### Why we study the $\eta$ production

- N\*: Nonperturbative QCD
- The problems about N\* spectrum compared with CQM.
- Ways to deepen our understanding of N\*
  - Photo-/electroproduction: CLAS,ELSA,MAMI...
  - $\pi$  N scattering: CB...
  - NN collision: COSY, HPLUS@CSR...
  - $J/\psi$  decay in  $e^+e^-$  collision: BES
  - . . . . . .

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### Why we study the $\eta$ production

- N\*: Nonperturbative QCD
- The problems about N\* spectrum compared with CQM.
- Ways to deepen our understanding of N\*
- Merit of  $\eta N$  channel
  - No  $\Delta$ :  $I_{\eta} + I_N \rightarrow 1/2$ , no 3/2
  - "η-mesic nuclei"
  - Important *S*<sub>11</sub>(1535)

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We will study N\* through two  $\eta$  production processes  $\gamma p \rightarrow \eta p$  and  $\pi^- p \rightarrow \eta n$ 

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### Extracting information of resonances from data

Extracting the information of resonances, such as, mass, decay width, from the observables

- SAID
- MAID
- EBAC
- Bonn-Gatchina groups
- Geissen group
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### Extracting information of resonances from data

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Do not reach the subnucleonic degrees of freedom!

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### Reach the subnucleonic degrees of freedom

Approaches with the subnucleonic degrees of freedom

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## Reach the subnucleonic degrees of freedom

Approaches with the subnucleonic degrees of freedom

- Approaches based on fundamental theory QCD
  - Lattice QCD: great technical difficulties for resonances.
  - QCD sum rule: low energy ones,  $\Delta(1232)$ ,  $S_{11}(1535)$ ....

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### Reach the subnucleonic degrees of freedom

Approaches with the subnucleonic degrees of freedom

- Approaches based on fundamental theory QCD
  - Lattice QCD
  - QCD sum rule
- Constituent quark model
  - spectrum:  $SU(6) \otimes U(3)$ .
  - transition amplitudes:  $A_{1/2}$ ,  $A_{3/2}$ ,  $\Gamma_{R \rightarrow MB}$

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## Reach the subnucleonic degrees of freedom

Approaches with the subnucleonic degrees of freedom

- Approaches based on fundamental theory QCD
  - Lattice QCD
  - QCD sum rule
- Constituent quark model
  - spectrum
  - transition amplitudes

Those approaches did not investigate reaction mechanisms.

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**Our approach**: study of  $\eta$  productions in constituent quark model

Starting Point: effective chiral Lagrangian,

$$\mathcal{L} = \bar{\psi}[\gamma_{\mu}(i\partial^{\mu} + V^{\mu} + \gamma_5 A^{\mu}) - m]\psi + \cdots,$$

where  $V^{\mu} = \frac{1}{2} (\xi \partial^{\mu} \xi^{\dagger} + \xi^{\dagger} \partial^{\mu} \xi)$ ,  $A^{\mu} = \frac{1}{2i} (\xi \partial^{\mu} \xi^{\dagger} - \xi^{\dagger} \partial^{\mu} \xi)$  with  $\xi = \exp(i\phi_m/f_m)$ .

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**Our approach**: study of  $\eta$  productions in constituent quark model

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For the pseudoscalar meson productions with the Hamiltonian

$$\mathcal{M}_{fi} = \langle N_f | H_{f,i} | N_i \rangle + \sum_j \left\{ \frac{\langle N_f | H_f | N_j \rangle \langle N_j | H_i | N_i \rangle}{E_i + \omega_i - E_j} + \frac{\langle N_f | H_i | N_j \rangle \langle N_j | H_f | N_i \rangle}{E_i - \omega_j - E_j} \right\} + \mathcal{M}_{\mathcal{T}},$$



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#### Formulism: for s-channel

Connect observables with CQM

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Connect observables with CQM

$$\begin{split} \mathcal{M}_{N^*}^{\gamma} & \to f_{1/\pm} \\ \mathcal{M}_{N^*}^{\gamma} &= \mathit{i} \mathit{f}_{1/\pm} \sigma \cdot \epsilon + \mathit{f}_{2/\pm} \sigma \cdot \hat{\mathbf{q}} \sigma \cdot (\hat{\mathbf{k}} \times \epsilon) + \mathit{i} \mathit{f}_{3/\pm} \sigma \cdot \hat{\mathbf{k}} \hat{\mathbf{q}} \cdot \epsilon + \mathit{i} \mathit{f}_{4/\pm} \sigma \cdot \hat{\mathbf{q}} \epsilon \cdot \hat{\mathbf{q}}, \\ \mathcal{M}_{N^*}^{\eta_*} &= \mathit{f}_{1/\pm} + \sigma \cdot \hat{\mathbf{q}} \sigma \cdot \hat{\mathbf{k}} \mathit{f}_{2/\pm}. \end{split}$$

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Connect observables with CQM

 $\mathcal{M}_{N^*}^{\gamma} \rightarrow f_{1/\pm}$ 

$$\begin{split} f_{1/\pm} & \longrightarrow A_{3/2}^{\gamma}, A_{1/2}^{m} \\ f_{1/\pm} & = f_{0} [\mp A_{1/2}^{\gamma} - \sqrt{\frac{l+1/2 \mp 1/2}{l+1/2 \pm 3/2}} A_{3/2}^{\gamma}] P_{\ell\pm1}^{\prime}, \\ f_{2/\pm} & = f_{0} [\mp A_{1/2}^{\gamma} - \sqrt{\frac{l+1/2 \pm 3/2}{l+1/2 \mp 1/2}} A_{3/2}^{\gamma}] P_{\ell}^{\prime}, \qquad f_{1} & = \sum_{l=0}^{\infty} [f_{l+}P_{l+1}^{\prime} - f_{l-}P_{l-1}^{\prime}], \\ f_{3/\pm} & = \pm f_{0} \frac{2A_{3/2}^{\gamma}}{\sqrt{(l-1/2 \pm 1/2)(l+3/2 \pm 1/2)}} P_{\ell\pm1}^{\prime\prime}, \qquad f_{2} & = \sum_{l=0}^{\infty} [f_{l-} - f_{l+}] P_{l}^{\prime}. \\ f_{4/\pm} & = \mp f_{0} \frac{2A_{3/2}^{\gamma}}{\sqrt{(l-1/2 \pm 1/2)(l+3/2 \pm 1/2)}} P_{\ell}^{\prime\prime}, \\ \text{where } A_{\lambda}^{\gamma} \text{ is the helicity amplitudes and } f_{0} \equiv \frac{1}{(2J+1)2\pi} [\frac{M_{N}E_{N}}{M_{N^{*}}^{N}} k]^{1/2} A_{1/2}^{m} \text{ with } A_{1/2}^{m} \text{ the } N^{*} \rightarrow \eta N \text{ decay amplitude, appearing in the partial decay width} \end{split}$$

Connect observables with CQM

 $\mathcal{M}_{N^*}^{\gamma} \rightarrow f_{1/\pm}$ 

 $f_{1/\pm} \rightarrow A^{\gamma}_{3/2}, A^m_{1/2}$ 

 $A_{3/2}^{\gamma}, A_{1/2}^{m} \rightarrow \langle N|H|N^{*} \rangle$ 

$$\begin{split} \mathbf{A}_{\lambda} &= \sqrt{\frac{2\pi}{k}} \langle \mathbf{N}^*; J\lambda | H_e | \mathbf{N}; \frac{1}{2}\lambda - 1 \rangle, \\ \mathbf{A}_{\nu}^m &= \langle \mathbf{N}; \frac{1}{2}\nu | H_m | \mathbf{N}^*; J\nu \rangle. \end{split}$$

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Connect observables with CQM

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 $A_{3/2}^{\gamma}, A_{1/2}^{m} \rightarrow \langle \textit{N} | \textit{H} | \textit{N}^{*} \rangle$ 

#### $\langle N|H|N^* \rangle$ in CQM

- Wave function:  $|N^*\rangle$  from potential model: OGE, GBE...
- Hamiltonian: H<sub>e</sub>, H<sub>m</sub>: χQM

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Connect observables with CQM

 $\mathcal{M}_{N^*}^{\gamma} \rightarrow f_{1/\pm}$ 

 $f_{1/\pm} \rightarrow A^{\gamma}_{3/2}, A^m_{1/2}$ 

 $A_{3/2}^{\gamma}, A_{1/2}^{m} \rightarrow \langle \textit{N} | \textit{H} | \textit{N}^{*} \rangle$ 

 $\langle N|H|N^* \rangle$  in CQM

#### Done

Advantage: the breaking of  $SU(6) \otimes O(3)$  symmetry is introduced through potential model to avoid a strength parameter (coupling constant) for each resonances

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#### Results

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#### Channels considered

- s-channel
  - n=1: S<sub>11</sub>(1535), S<sub>11</sub>(1650), D<sub>13</sub>(1520), D<sub>13</sub>(1700), and D<sub>15</sub>(1675);  $n=2: P_{11}(1440), P_{11}(1710), P_{13}(1720),$  $P_{13}(1900), F_{15}(1680), F_{15}(2000), and$  $F_{17}(1990).$ n>2: degenerated New resonances: S<sub>11</sub>, D<sub>13</sub>, D<sub>15</sub>
- u-channel : Degenerated
- t-channel : Neglected

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#### Data

#### Results

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#### Results

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## Experiments: $d\sigma/d\Omega$ and $\Sigma$ , T for $\gamma p \rightarrow \eta p$

Exp.	year	Obs.	Angular	$P_\gamma$	W	N <sub>dp</sub>
MAMI	(1994)	$d\sigma/d\Omega$	25-154	0.716-0.790	1.49-1.54	100
CLAS	(2002)	$d\sigma/d\Omega$	45-134	0.775-1.925	1.53-2.12	190
ELSA	(2003)	$d\sigma/d\Omega$	31-138	0.775-2.900	1.53-2.51	631
LNS	(2006)	$d\sigma/d\Omega$	25-154	0.718-1.142	1.49-1.74	180
GRAAL	(2006)	$d\sigma/d\Omega$	31-160	0.714-1.477	1.49-1.91	487
ELSA	(2006)	Σ	50-148	0.843-1.343	1.57-1.84	34
GRAAL	(2006)	Σ	40-160	0.724-1.472	1.50-1.91	150
BONN	(1997)	Т	33-145	0.717-1.105	1.49-1.72	50

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## Experiments: $d\sigma/d\Omega$ for $\pi^- p \rightarrow \eta n$

Ref.	Year	Angular	$P_{\pi}$	W	$\delta_{sys}$
Deinet	(1969)	32-123	0.718-1.050	1.51-1.70	11%
Richards	(1970)	26-154	0.718-1.433	1.51-1.90	10% to 14%
Debenham	(1975)	162-172	0.697-0.999	1.49-1.67	10% + 0.02 $\mu$ b
Brown	(1975)	18-160	0.724-2.724	1.51-2.45	10% or 0.01 $\mu$ b
Prakhov	(2005)	23-157	0.687-0.747	1.49-1.52	6%
Feltesse	(1975)	20-160	0.757	1.53	
Crouch	(1980)	14-167	1.395-3.839	1.88-2.85	
Morrison	(1999)	46-134	0.701-0.747	1.50-1.52	

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## Experiments: $d\sigma/d\Omega$ for $\pi^- p \rightarrow \eta n$

Ref.	Year	Angular	$P_{\pi}$	W	$\delta_{sys}$	N <sub>dp</sub>	N <sub>dp</sub>
Deinet	(1969)	32-123	0.718-1.050	1.51-1.70	11%	83	80
Richards	(1970)	26-154	0.718-1.433	1.51-1.90	10% to 14%	70	66
Debenham	(1975)	162-172	0.697-0.999	1.49-1.67	10% + 0.02 $\mu$ b	111	27
Brown	(1975)	18-160	0.724-2.724	1.51-2.45	10% or 0.01 $\mu$ b	379	51
Prakhov	(2005)	23-157	0.687-0.747	1.49-1.52	6%	84	70
Feltesse	(1975)	20-160	0.757	1.53		16	-
Crouch	(1980)	14-167	1.395-3.839	1.88-2.85		731	-
Morrison	(1999)	46-134	0.701-0.747	1.50-1.52		34	-

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#### Results

 $\chi^2 = \sum \frac{(V_{th} - V_{ex})^2}{(E_{ex}^V)^2 + (V_{th}' E_{ex}^E)^2}.$  Here  $V_{th}$ ,  $V_{ex} E_{ex}^V$  and  $E_{ex}^E$  are the values from theoretical calculation and experiment and the uncertainty of observable and energy, and  $V_{th}'$  are the derivative of observable with energy.

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# Spectrum



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## Results: $\sigma$ for $\gamma p \rightarrow \eta p$



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## Results: $d\sigma/d\Omega$ and $\Sigma$ for $\overline{\gamma p \rightarrow \eta p}$



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Results:  $\sigma$  for  $\pi^- p \rightarrow \eta n$ 



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## Results: $d\sigma/d\Omega$ for $\pi^- p \rightarrow \eta n$



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#### Results: Helicity amplitudes and decay widths

Resonances	A <sub>1/2</sub>	$A_{1/2}^{PDG}$	A <sub>3/2</sub>	$A_{3/2}^{PDG}$	$\sigma \sqrt{\Gamma_{\eta N}}$	$(\sigma)\sqrt{\Gamma_{\eta N}^{PDG}}$	$\sqrt{\Gamma_{\pi N}}$	$\sqrt{\Gamma_{\pi N}^{PDG}}$
S <sub>11</sub> (1535)	73	$90 \pm 30$			7.18	$8.87^{+1.37}_{-1.37}$	6.78	$8.22^{+1.59}_{-1.60}$
$S_{11}(1650)$	66	$53\pm16$			-2.42	$1.95^{+0.94}_{-1.57}$	8.85	$11.31^{+1.95}_{-1.98}$
$P_{11}(1440)$	-23	-65 $\pm$ 4			-2.42		17.16	$13.96^{+4.41}_{-3.48}$
$P_{11}(1710)$	-53	$9\pm22$			-1.05	$2.49^{+1.75}_{-0.88}$	4.12	$3.87^{+3.20}_{-1.64}$
P <sub>11</sub>	18				-2.79		6.59	
P <sub>11</sub>	3				-1.20		4.51	$5.34^{+2.16}_{-2.16}$
$P_{13}(1720)$	177	$18\pm30$	-69	-19 $\pm$ 20	2.91	$2.83^{+1.04}_{-0.71}$	20.15	$5.48^{+2.27}_{-1.60}$
$P_{13}(1900)$	30		2		-1.33	$8.35^{+2.11}_{-2.20}$	11.02	$11.38^{+2.20}_{-2.21}$
P <sub>13</sub>	28		0		2.44		3.06	
P <sub>13</sub>	12		2		0.03		5.54	
P <sub>13</sub>	-3		3		-1.01		3.12	
$D_{13}(1520)$	-7	$-24 \pm 9$	158	$166\pm5$	0.44	$0.51\substack{+0.07 \\ -0.06}$	14.77	$8.31^{+0.71}_{-0.53}$
$D_{13}(1700)$	-4	$\textbf{-18}\pm\textbf{13}$	4	$-2 \pm 24$	-0.81	$0.00^{+1.22}_{-0.00}$	4.92	$3.16^{+1.58}_{-1.58}$
$D_{15}(1675)$	-6	$19\pm8$	-8	$15\pm9$	-2.50	$0.00^{+1.28}_{-0.00}$	7.59	$7.75^{+0.87}_{-1.00}$
$F_{15}(1680)$	24	-15 $\pm$ 6	136	$133 \pm 12$	0.58	$0.00^{+1.18}_{-0.00}$	13.71	$9.37^{+0.53}_{-0.54}$
F <sub>15</sub>	-9		4		0.97		0.35	
$F_{15}(2000)$	-1		10		-0.47		3.60	$4.00^{+6.20}_{-2.18}$
$F_{17}(1990)$	5	1	6	4	-1.55	$0.00\substack{+2.17 \\ -0.00}$	6.84	$4.58^{+1.55}_{-1.55}$

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# Conclusion

- Study of  $\eta$  productions in quark level (CQM).
- Study of "missing" resonances in  $\eta$  productions directly in CQM.
- Study of spectrum and observables simultaneously in CQM.
- For known resonances:
  - Both :  $S_{11}(1535), S_{11}(1650), D_{13}(1520), F_{15}(1680), P_{13}(1720)$
  - $\pi^- p \rightarrow \eta n$  :  $P_{11}(1440)$ ,  $D_{15}(1675)$
- For "missing" resonances:
  - Both : Negligible
- For New resonances:
  - $\gamma \quad p \rightarrow \eta p$ : New  $S_{11}(1715)$  and  $D_{15}(2090)$
- For  $\pi^- p \rightarrow \eta n$ 
  - P<sub>13</sub>(1720) : second bump for the cross section
  - Near threshold: Large  $d(d\sigma/d\Omega)/dE \rightarrow \Delta E$  is more important!