

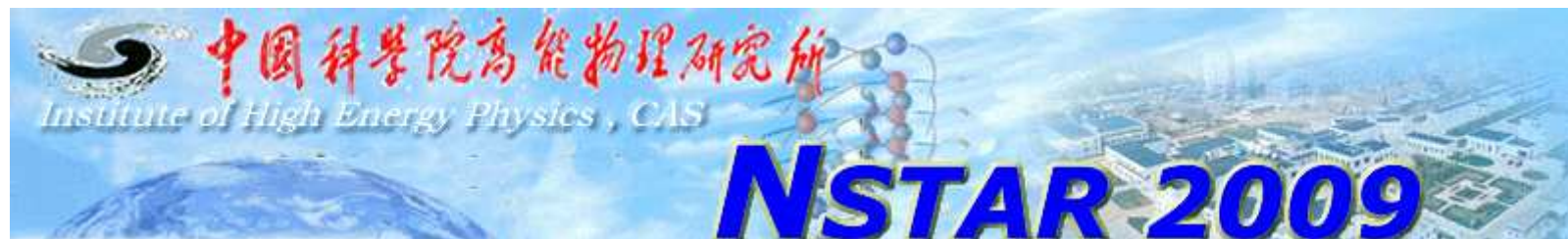


Decay of Roper Resonance in Hybrid Baryon Model

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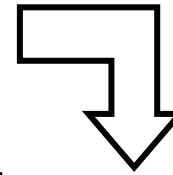


Outline

- **Introduction**
- **Hybrid baryon model**
- **Interaction operator**
- **Transition amplitude**
- **Results on decay widths**
- **Conclusion**

Baryon excitation states

understanding



Nucleon internal structure

N, Nucleon(939)

$$J = 1/2, I = 1/2, P = +1$$

$\Delta(1232)$

$$J = 3/2, I = 3/2, P = +1$$

$N^*(1440)$, Roper Resonance

$$J = 1/2, I = 1/2, P = +1$$

$N^*(1440)$

Intrinsic excitation of Nucleon
at energy around 500 MeV

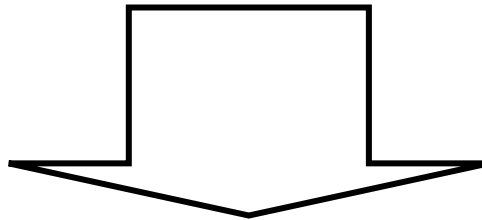
Simple Model for $N^*(1440)$

$N^*(1440)$ is three-quark with the **first radial excitation**

$N^*(1440) \Rightarrow$ too **low mass** to be **radial excitation**

Various quarks models

**Met difficulties to explain its mass
and electromagnetic couplings**



problems of the models

Hybrid Baryon Model

Bag Model

$N^*(1440)$

Hybrid

qqq

G

TE (transverse electric)

TM (transverse magnetic)

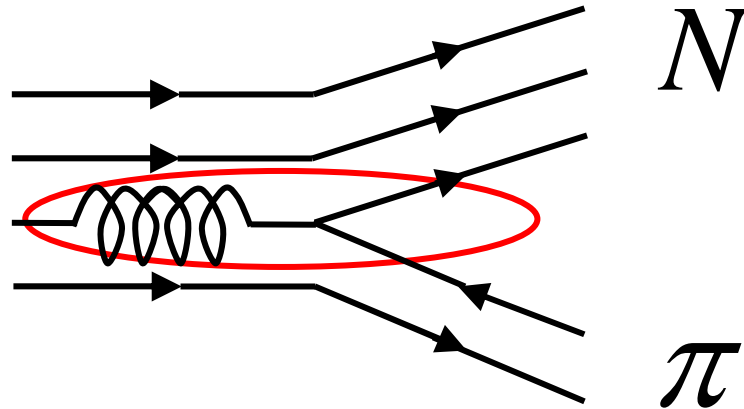
low mass

TE mode is the lowest eigenmode

$N^*(1440)$ \Rightarrow 3 quarks and a TE gluon

q^3G

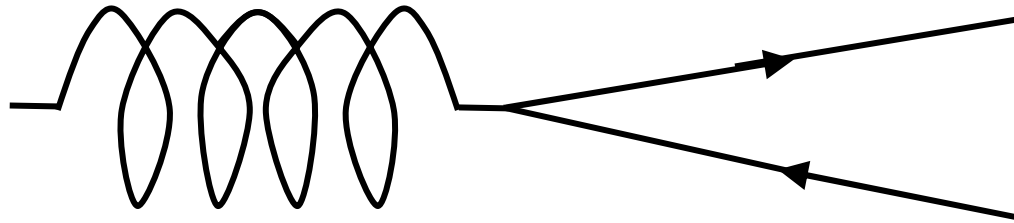
$$N^*(1440) \rightarrow N\pi$$

 $N^*(1440)$


Wave functions

 $\psi_{N^*(1440)}$
 ψ_N
 ψ_π

Interaction Operator



Nucleon wave function

$$\psi_N = \psi_N^{\text{spatial}} \psi_N^{\text{spin-flavor}} \psi_N^{\text{color}}$$

qqq \Rightarrow harmonic oscillator potential

$$\psi_N = N_N \exp\left[-\frac{1}{4}a^2 (\vec{p}_1 - \vec{p}_2)^2\right] \exp\left[-\frac{1}{12}a^2 (\vec{p}_1 + \vec{p}_2 - 2\vec{p}_3)^2\right]$$

$$\frac{1}{\sqrt{2}} \sum_{J=0,1} \left| \left(\frac{1}{2}^{(1)} \otimes \frac{1}{2}^{(2)} \right)_J \otimes \frac{1}{2}^{(3)} \right\rangle_{\frac{1}{2}, S_Z}^{\text{Spin}} \left| \left(\frac{1}{2}^{(1)} \otimes \frac{1}{2}^{(2)} \right)_J \otimes \frac{1}{2}^{(3)} \right\rangle_{\frac{1}{2}, T_Z}^{\text{Flavor}}$$

$$\frac{1}{\sqrt{6}} \sum_{i,j,k} \varepsilon_{ijk} |q_1\rangle_i |q_2\rangle_j |q_3\rangle_k$$

Pion wave function

$$\psi_{\pi} = \psi_{\pi}^{\text{spatial}} \psi_{\pi}^{\text{spin-flavor}} \psi_{\pi}^{\text{color}}$$

$q\bar{q}$ \Rightarrow harmonic oscillator potential

$$\psi_{\pi} = N_{\pi} \exp\left[-\frac{1}{8}b^2(\bar{p}_1 - \bar{p}_2)^2\right] \left| \left(\frac{1}{2}^{(1)} \otimes \frac{1}{2}^{(2)} \right) \right\rangle_{0,0}^{\text{Spin}}$$

$$\left| \left(\frac{1}{2}^{(1)} \otimes \frac{1}{2}^{(2)} \right) \right\rangle_{1,t_z}^{\text{Flavor}} \frac{1}{\sqrt{3}} \sum_{l=1}^3 |\bar{q}_1\rangle_l |q_2\rangle_l$$

N*(1440) wave function

$$\Psi_{N^*(1440)} = \Psi_{N^*(1440)}^{\text{spatial}} \Psi_{N^*(1440)}^{\text{spin-flavor-color}}$$

$\Psi_{N^*(1440)}^{\text{spatial}}$ q^3G \Rightarrow harmonic oscillator potential

$$\Psi_{N^*}^{\text{Spatial}} = N_{N^*} \exp\left[-\frac{1}{4}a^2(\bar{p}_1 - \bar{p}_2)^2\right] \exp\left[-\frac{1}{12}a^2(\bar{p}_1 + \bar{p}_2 - 2\bar{p}_3)^2\right] \exp\left[-\frac{2}{3} \frac{a^2}{(3R+1)^2} (\bar{p}_1 + \bar{p}_2 + \bar{p}_3 - 3R\bar{p}_4)^2\right]$$

where $R = \frac{m_q}{m_g}$

$$\psi_{N^*(1440)}^{\text{spin-flavor-color}} = \mathbf{A} \left| {}^2N_g \right\rangle + \mathbf{B} \left| {}^4N_g \right\rangle$$

$2S + 1$ S : total spin of 3 quarks

$$\left| {}^2N_g \right\rangle = \frac{1}{2} \left[(\phi^\rho \chi^\rho - \phi^\lambda \chi^\lambda) \varphi^\rho - (\phi^\rho \chi^\lambda - \phi^\lambda \chi^\rho) \varphi^\lambda \right] \otimes |G\rangle$$

$$\left| {}^4N_g \right\rangle = \frac{1}{\sqrt{2}} \left[(\phi^\lambda \varphi^\rho - \phi^\rho \varphi^\lambda) \chi^S \right] \otimes |G\rangle$$

ϕ flavor wave function

φ color wave function

χ spin wave function

$|G\rangle$ Gluon wave function

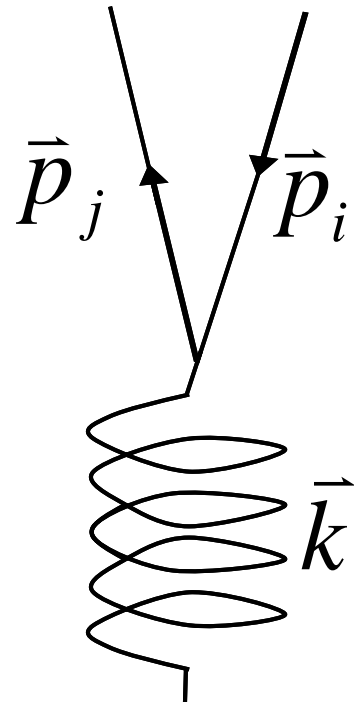
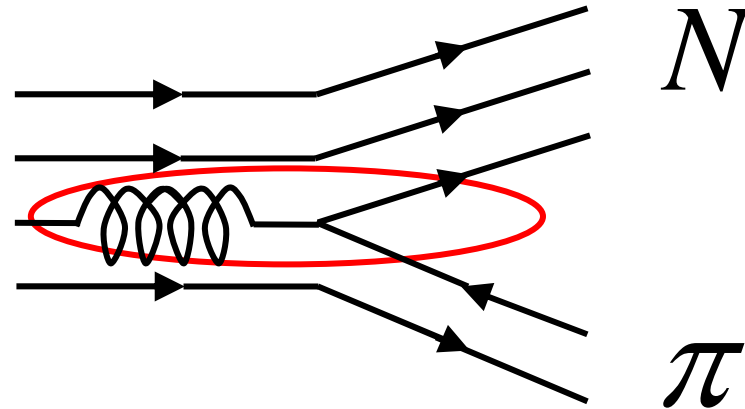
S totally symmetric

λ mix symmetric

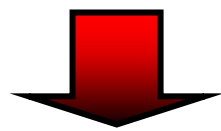
ρ mix antisymmetric

Interaction Operator

$$N^*(1440) \rightarrow N\pi$$

$$N^*(1440)$$


$$L_{\text{int}} = \bar{\psi} g \gamma^\mu A_\mu^\alpha \frac{\lambda^\alpha}{2} \psi$$

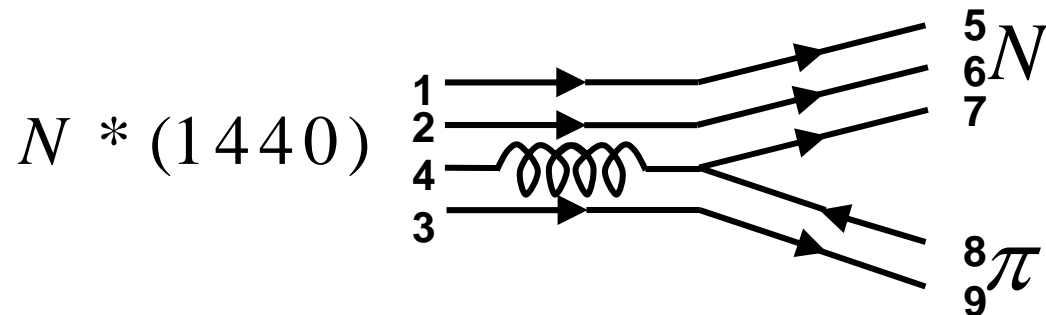


Nonrelativistic approximation

$3S_1$

$$V_{ij} = g \sum_{\mu} (-1)^{\mu+1} \sigma_{ij}^{\mu} \mathbf{1}_{ij}^F \varepsilon_{-\mu} A^\alpha \frac{\lambda^\alpha}{2} \delta(\vec{p}_i + \vec{p}_j - \vec{k})$$

Transition amplitude



$$T = \langle N\pi | V_{78} | N^* \rangle$$

$$V_{78} = g \sum_{\mu} (-1)^{\mu+1} \sigma_{78}^{\mu} \mathbf{1}_{78}^F \varepsilon_{-\mu} A^{\alpha} \frac{\lambda^{\alpha}}{2} \delta(\vec{p}_7 + \vec{p}_8 - \vec{p}_4)$$

Properties

$$\langle 0,0 | \sigma_{i,j}^{\mu} | [\bar{\chi}_i \otimes \chi_j]_{JM} \rangle = (-1)^M \sqrt{2} \delta_{J,1} \delta_{M,-\mu}$$

$$\langle 0,0 | \mathbf{1}_{78}^F | T, T_Z \rangle = \sqrt{2} \delta_{T,0} \delta_{T_Z,0}$$

$$\langle 0,0 | \varepsilon_{-\mu} | e_{sm_s} \rangle = \delta_{-\mu, m_s}$$

Decay width



$$\Gamma_{N^*(1440) \rightarrow N\pi} \propto \frac{E_N E_\pi}{M_{N^*}} g^2 p^3 \beta \exp(-2\alpha p^2)$$

depend on
 a, b

$$\times \sum_{S_z} \sum_{m_l} \left| (-1)^{m_l} \left(\frac{A}{2} Q_1 + \frac{B}{\sqrt{2}} Q_2 \right) \right|^2$$

depend on
 a, b, R

$$\psi_{N^*(1440)}^{\text{spin-flavor-color}} = A \left| {}^2N_g \right\rangle + B \left| {}^4N_g \right\rangle$$

$$Q_1 = \left\langle \psi_{N\pi}^{\text{spin-flavor-color}} \left| \mathcal{O}_{\text{spin-flavor-color}} \right| {}^2N_g \right\rangle$$

$$Q_2 = \left\langle \psi_{N\pi}^{\text{spin-flavor-color}} \left| \mathcal{O}_{\text{spin-flavor-color}} \right| {}^4N_g \right\rangle$$

Where a and b are determined by the nucleon and meson sizes,
 $NN, N\bar{N}, \pi N$ interactions

$$a = 3.1 \text{ GeV}^{-1}$$

$$b = 4.1 \text{ GeV}^{-1}$$

Decay channels**Interested channels**

$$N^*(1440) \rightarrow N\pi$$

$$N^*(1440) \rightarrow \Delta\pi$$

$$N^*(1440) \rightarrow N\rho$$

$$N^*(1440) \rightarrow N\eta$$

Results on decay widths

PRD 46 (1992) p70-74

	Hybrid Baryon Model				Experimental Data (PDG)
	A=0 B=1	A=1 B=0	A/B=1 $A=\sqrt{\frac{1}{2}}$	A/B=-1 $A=-\sqrt{\frac{1}{2}}$	
$\frac{\Gamma_{N^*(1440) \rightarrow N\rho}}{\Gamma_{N^*(1440) \rightarrow N\pi}}$	0.027	0.018	0.016	0.021	< 0.15
$\frac{\Gamma_{N^*(1440) \rightarrow N\eta}}{\Gamma_{N^*(1440) \rightarrow N\pi}}$	0.784	0.196	0.577	0.124	-
$\frac{\Gamma_{N^*(1440) \rightarrow \Delta\pi}}{\Gamma_{N^*(1440) \rightarrow N\pi}}$	1.812	0.065	0.179	0.392	0.27 - 0.55

Conclusions

- $N^*(1440)$ is treated as a hybrid baryon, $q^3 G$, in nonrelativistic regime.
- The vertex where a gluon is destroyed and form a pair of quark-antiquark is described in the effective 3S_1
- The decay width ratios of $N^*(1440)$ have been shown. It is found that our theoretical results are fairly consistent with experimental data.

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Expand the total transition amplitude in **partial wave**

$$T_{N^*(1440) \rightarrow N\pi} = \sum_{l, m_l} T_{lm_l} Y_{lm_l}^*(\hat{p})$$

$$T_{lm_l} = g \delta_{S_Z, S_Z''} \delta_{T_Z, T_Z''} P_{1m_l} \left[\frac{A}{2} Q_1 + \frac{B}{\sqrt{2}} Q_2 \right]$$

depend on m_l, a, b, c, R

$$Q_1 = \left\langle \psi_{N\pi}^{\text{spin-flavor-color}} \left| \mathcal{O}_{\text{spin-flavor-color}} \right| {}^2 N'_g \right\rangle$$

$$Q_2 = \left\langle \psi_{N\pi}^{\text{spin-flavor-color}} \left| \mathcal{O}_{\text{spin-flavor-color}} \right| {}^4 N'_g \right\rangle$$

$$N^*(1440) \rightarrow N\pi$$

m_l	S_z	Q_1	Q_2
-1	-1/2	-0.148	0.131
	1/2	-0.444	0.079
0	-1/2	-0.296	-0.052
	1/2	0.296	0.052
1	-1/2	0.444	-0.079
	1/2	0.148	-0.131

Values of Q_1 and Q_2 for different initial and final states

$$\begin{aligned}
 |^4N_g\rangle = \frac{1}{\sqrt{2}} & \left[\left| \left(\frac{1}{2}^{(1)} \otimes \frac{1}{2}^{(2)} \right)_1 \otimes \frac{1}{2}^{(3)} \right\rangle_{\frac{3}{2}, m_{123}} \otimes \left(Y_{1m_l} (\vec{p}_1 + \vec{p}_2 + \vec{p}_3 - 3R\vec{p}_4) \otimes e_{1,m_s} \right)_{1,m_2'} \right]_{\frac{1}{2}, S_Z''} \text{Spin} \\
 & \left\{ \left[\left| \left(\frac{1}{2}^{(1)} \otimes \frac{1}{2}^{(2)} \right)_1 \otimes \frac{1}{2}^{(3)} \right\rangle_{\frac{1}{2}, T_Z''} \text{Flavor} \frac{1}{\sqrt{8}} \sum_{\alpha} \psi_{\alpha}^{\rho} |g\rangle_{\alpha}^{\text{color}} \right] \right. \\
 & \left. - \left[\left| \left(\frac{1}{2}^{(1)} \otimes \frac{1}{2}^{(2)} \right)_0 \otimes \frac{1}{2}^{(3)} \right\rangle_{\frac{1}{2}, T_Z''} \text{Flavor} \frac{1}{\sqrt{8}} \sum_{\alpha} \psi_{\alpha}^{\lambda} |g\rangle_{\alpha}^{\text{color}} \right] \right\}
 \end{aligned}$$

$$|{}^2N_g\rangle = \frac{1}{2} \sum_{J_{12}} \left\{ \left[\left(\left(\frac{1}{2}^{(1)} \otimes \frac{1}{2}^{(2)} \right)_{J_{12}} \otimes \frac{1}{2}^{(3)} \right)_{\frac{1}{2}, m_{123}} \otimes \left(Y_{1, m_l} (\bar{p}_1 + \bar{p}_2 + \bar{p}_3 - 3R\bar{p}_4) \otimes e_{1, m_s} \right)_{1, m_2'} \right] \right\} \left. \begin{array}{l} \text{Spin} \\ \frac{1}{2}, S_Z'' \end{array} \right.$$

$$\left[\left(\left(\frac{1}{2}^{(1)} \otimes \frac{1}{2}^{(2)} \right)_{J_{12}} \otimes \frac{1}{2}^{(3)} \right)_{\frac{1}{2}, T_Z''} \text{Flavor} \frac{1}{\sqrt{8}} \sum_{\alpha} \psi_{\alpha}^{\rho} |g\rangle_{\alpha}^{\text{color}} \right. \\ \left. - \left[\left[\left(\left(\frac{1}{2}^{(1)} \otimes \frac{1}{2}^{(2)} \right)_{1-J_{12}} \otimes \frac{1}{2}^{(3)} \right)_{\frac{1}{2}, m_{123}} \otimes \left(Y_{1, m_l} (\bar{p}_1 + \bar{p}_2 + \bar{p}_3 - 3R\bar{p}_4) \otimes e_{1, m_s} \right)_{1, m_2'} \right] \right] \right. \left. \left. \left[\left(\left(\frac{1}{2}^{(1)} \otimes \frac{1}{2}^{(2)} \right)_{J_{12}} \otimes \frac{1}{2}^{(3)} \right)_{\frac{1}{2}, T_Z''} \text{Flavor} \frac{1}{\sqrt{8}} \sum_{\alpha} \psi_{\alpha}^{\lambda} |g\rangle_{\alpha}^{\text{color}} \right] \right] \right\} \left. \begin{array}{l} \text{Spin} \\ \frac{1}{2}, S_Z'' \end{array} \right.$$