

A new N^* resonance as a hadronic molecular state

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collaboration with
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new N^* resonance at around 1910 MeV

20~40 MeV blow $KK^{\text{bar}}N$ threshold (1930 MeV)

hadronic molecular state

Kaons are constituents



References

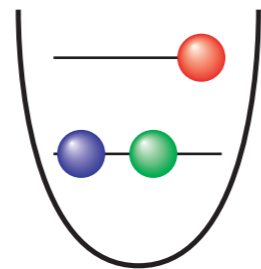
D. Jido and Y. Kanada-En'yo, **Phys. Rev. C78, 035203 (2008)**

See also for Ξ^* : Y. Kanada-En'yo and D. Jido, **Phys. Rev. C78, 025212 (2008)**

Introduction

So far, the structure of baryons has been investigated in quark models.

symmetries of quarks play a major role

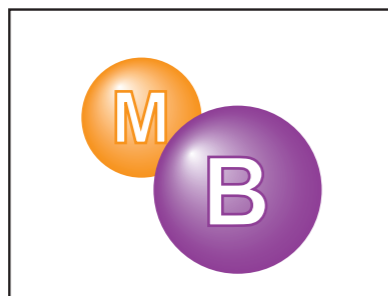


shell model picture

Baryon resonances : decay with strong interactions

large meson-baryon components

meson-baryon dynamics is also important



cluster picture

- inter-hadron distance is larger than quark dynamics
- difference in range of dynamics

What is hadronic molecular state ?

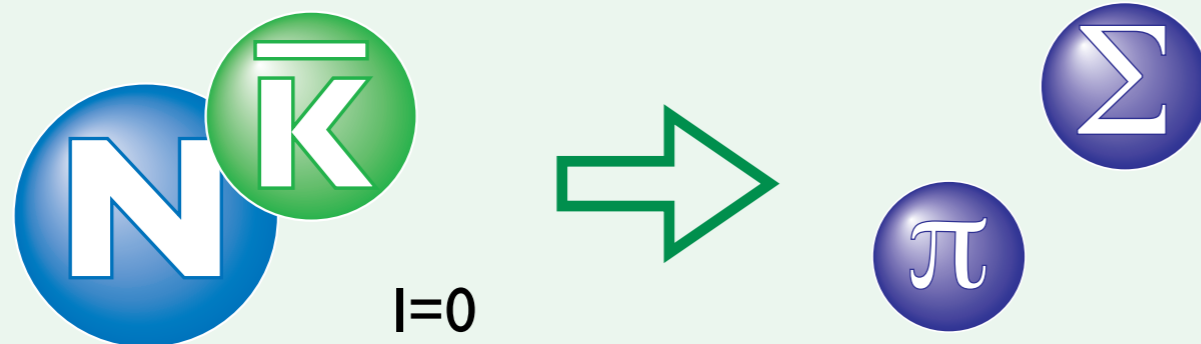
- **system of multiple hadrons described by hadron dynamics**

typical constituents are ground states hadrons

octet meson: π, K, η

octet baryons: N, Λ, Σ, Ξ

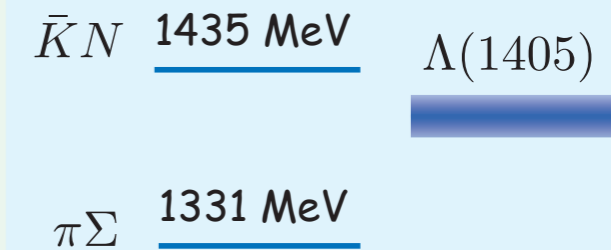
$\Lambda(1405)$: one of the historical examples



have been considered as a quasi-bound state of $K^{\text{bar}}N$ since late 50's

Recently this idea has been developed to light nuclear system

$\Lambda(1405)$ is described as a superposition of two states.



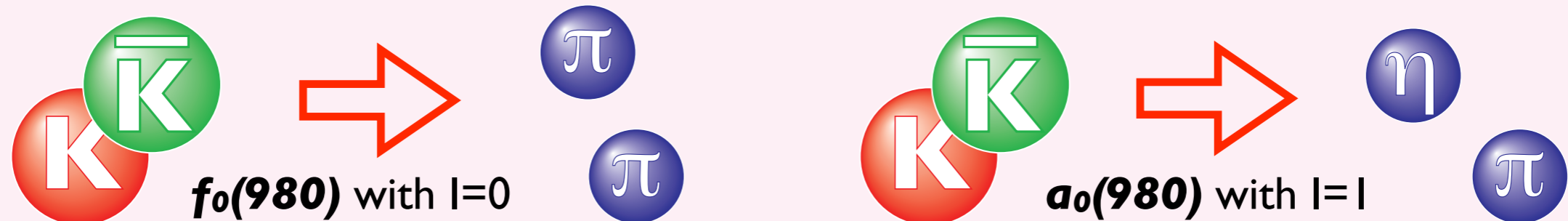
Dalitz, Tuan, PRL 2, 425 ('59)

Akaishi, Yamazaki, PRC 65, 044005 ('02)

DJ, Oller, Oset, Ramos, Meissner
NPA725, 181 ('03)

$f_0(980), a_0(980)$: the candidates of KK^{bar} QBS

scalar mesons



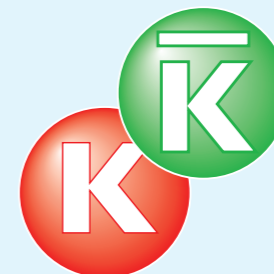
Weinstein, Isgur, PRL 48, 659 ('90)

What is hadronic molecular state ?

$\Lambda(1405)$



$f_0(980), a_0(980)$



- **weakly bound system with large width**

typical binding energy $\sim 10\text{-}30$ MeV

decay width ~ 50 MeV (strong interactions)

quasi-bound state

- **constituents keep their identity**

spatially extended (large size)

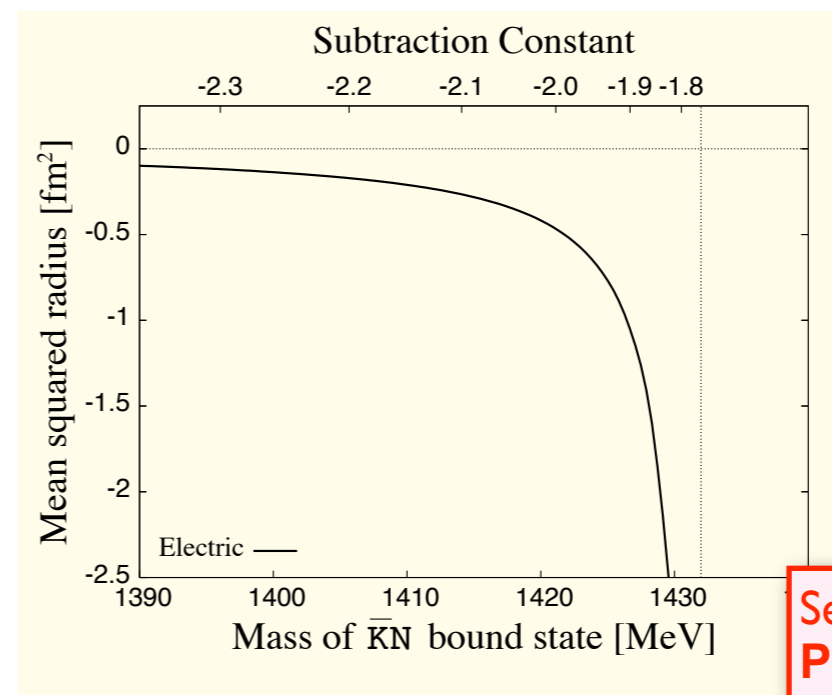
typically more than 1 fm

- **fragile system**

softer form factors

strong energy dependence

in production



charge radius
of $\Lambda(1405)$ as
 $K^{\text{bar}}N$ bound state

Sekihara, Hyodo, DJ
PLB669, 133 (2008)

quark degrees of freedom may be less important

Peculiarities of K meson

small binding energy $\sim 10\text{-}30$ MeV small kinetic energy

- **heavy particle** compared with kinetic energy

half of nucleon mass

cf. pion $m_\pi \approx 140$ MeV

non-relativistic potential model

isospin averaged mass

$$m_K = 495.7 \text{ MeV}$$

$$m_N = 938.9 \text{ MeV}$$

- **Nambu-Goldstone boson**

smaller mass compared with typical hadron mass scale

strong s-wave attraction in $K^{\text{bar}}N$

chiral effective theory momentum expansion

s-wave int. proportional to K energy

Kaons are different from pions in the energies of our interest !!

$K \bar{K} N$ system with $I=1/2, J^P=1/2^+$

simplest multi-kaon baryonic system:

assumption

non-relativistic treatment of kaons

$\Lambda(1405)$ is a quasi-bound state of $K^{\text{bar}}N$

$f_0(980)$ and $a_0(980)$ are quasi-bound states of KK^{bar}



s-wave

Binding energy is important for “fate” of hadronic molecular state

A) 3-body System $>$ 2-body BS + hadron

fall apart, large width

B) 3-body BS $<$ 2-body BS + hadron

quasi-stable

C) 3-body BS \ll 2-body BS + hadron

molecular picture broken down
large width of two-body decay

→ different approaches are necessary
3-body \rightarrow 2-body

Two-body interactions

$\Lambda(1405)$



$f_0(980), a_0(980)$



Interactions in $KK^{\text{bar}}N$ system

		$I=0$	$I=1$	threshold
N^* 	$\bar{K}N$	attraction $\Lambda(1405)$	weak attraction	1434.6 MeV
	$K\bar{K}$	$f_0(980)$	$a_0(980)$	991.4 MeV
	KN	repulsion very weak	strong repulsion	1434.6 MeV

Interactions in $K^{\text{bar}}K^{\text{bar}}N$ system

		$I=0$	$I=1$	threshold
Ξ^* 	$\bar{K}N$	attraction $\Lambda(1405)$	weak attraction	1434.6 MeV
	$\bar{K}\bar{K}$	repulsion weak	strong repulsion	991.4 MeV

Formulation for three-body system

non-relativistic potential model

Hamiltonian

$$H = T + V_{\bar{K}N}(r_1) + V_{KN}(r_2) + V_{K\bar{K}}(r_3),$$

two-body effective interactions

local potentials obtained by s-wave two-body scattering

Gaussian potential $V(r) = U \exp[-(r/b)^2]$

complex potentials to implement coupled-channels effects

no three-body interactions ex: $\bar{K}N \rightarrow \pi\Sigma$

no transitions to two hadrons $\bar{K}\bar{K}N \rightarrow MB$

will be suppressed in hadronic molecular states

recipe

1st: solve Schrödinger eq. without imaginary potential in variational method

obtain wavefunction Ψ and real part of energy

2nd: estimate imaginary part of energy $E^{\text{Im}} = \langle \Psi | \text{Im}V | \Psi \rangle.$

Effective interactions

Gaussian potential $V(r) = U \exp[-(r/b)^2]$

$\bar{K}N$
attractive

HW-HNJH and AY potentials

binding energy : 11 MeV (HW), 31 MeV (AY)

K^{bar} -N distance : **1.9 fm** (HW), **1.4 fm** (AY)

$K\bar{K}$
attractive

reproduce masses and widths of f_0 and a_0

mass: 980 MeV, width: 60 MeV

reproduced

binding energy : 11 MeV

K - K^{bar} distance : **2.1 fm**

PDG

mass: 980 ± 10 MeV

width: 40~100 MeV

mass: 984 ± 1.2 MeV

width: 50~100 MeV

KN
repulsive

reproduce scattering lengths

experimental data

$$a_{KN}^{I=0} = -0.035 \text{ fm}$$

$$a_{KN}^{I=1} = -0.310 \pm 0.003 \text{ fm}$$

Results of $KK^{\text{bar}}N$ system

N^* at 1910 MeV

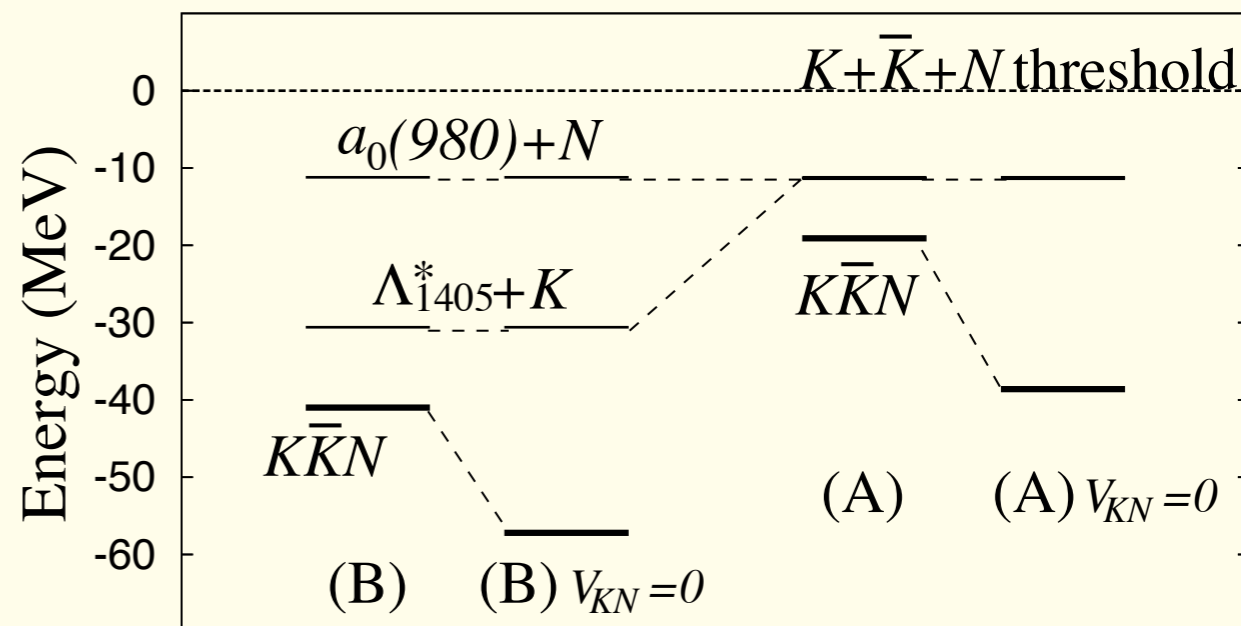
DJ, Y. Kanada-En'yo, PRC78, 035203 (2008)

$K\bar{K}N$ is bound below thresholds of $\Lambda(1405)+K$, $a_0(f_0)+N$

- loosely bound system

binding energy	width
HW: 19 MeV	88 MeV
AY: 39 MeV	98 MeV

sum of those of isolated two-particle systems



threshold of $KK^{\text{bar}}N$ 1930 MeV

$\Lambda(1405)$



binding energy	width
HW: 11 MeV	44 MeV
AY: 31 MeV	40 MeV

$f_0(980), a_0(980)$



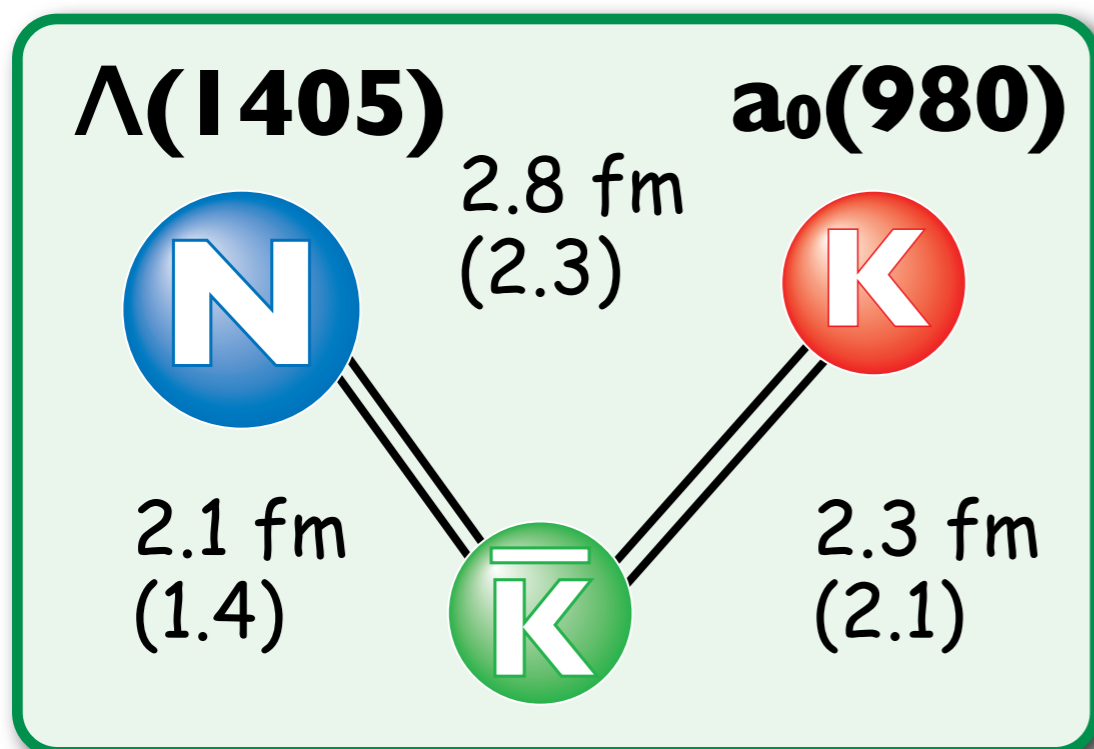
binding energy	width
11 MeV	60 MeV

Faddeev calculation also reproduces this resonance

Martinez Torres, Khemchandani, Oset, arXiv:0812.2235 [nucl-th]

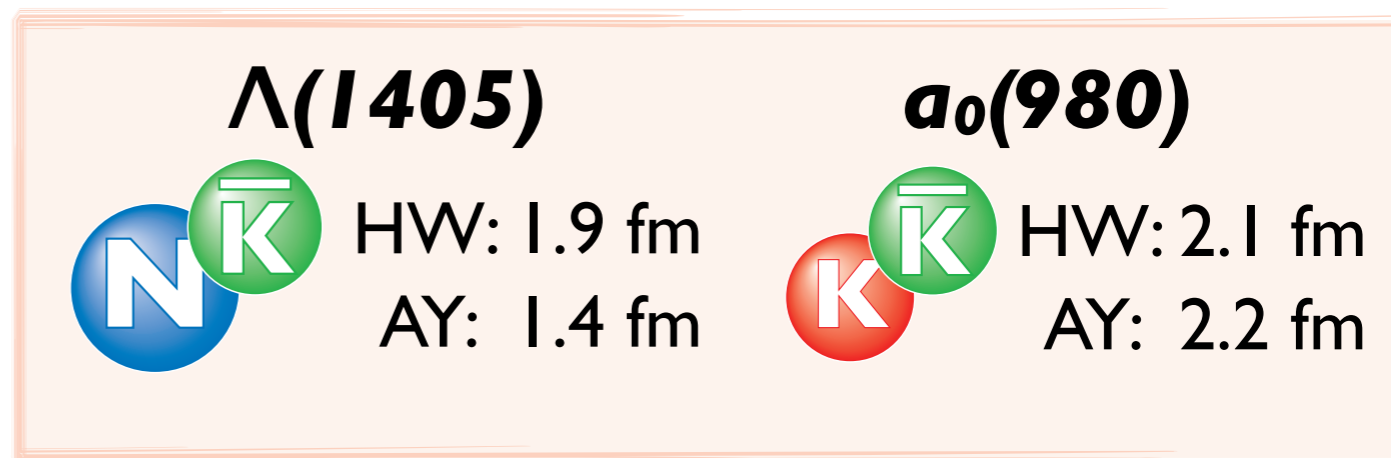
Structure of $N^*(1910)$

spatial structure



r.m.s radius: **1.7 fm** cf. 1.4 fm for ${}^4\text{He}$

mean hadron density: **0.07 hadrons/fm³**



hadron-hadron distances are comparable with nucleon-nucleon distances in nuclei

- **coexistence of two quasi-bound states keeping their characters**

$\Lambda(1405)+K$ $a_0(980)+N$

- **main decay modes**

$\pi\Sigma K$ from $\Lambda(1405)$

$\pi\eta N$ from $a_0(980)$

Two-body decay of N^*

possible two-body decay modes

based on geometrical argument

$$(K\bar{K}N) \rightarrow \pi N, \eta N, K\Lambda, K\Sigma$$

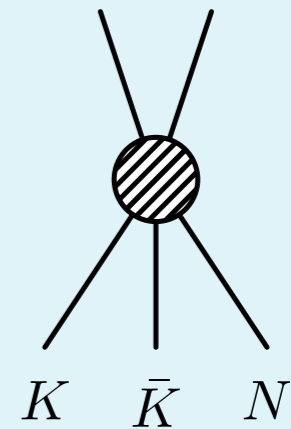
radius of system ~ 1.7 fm

spatially extended

two-body decays are strongly suppressed

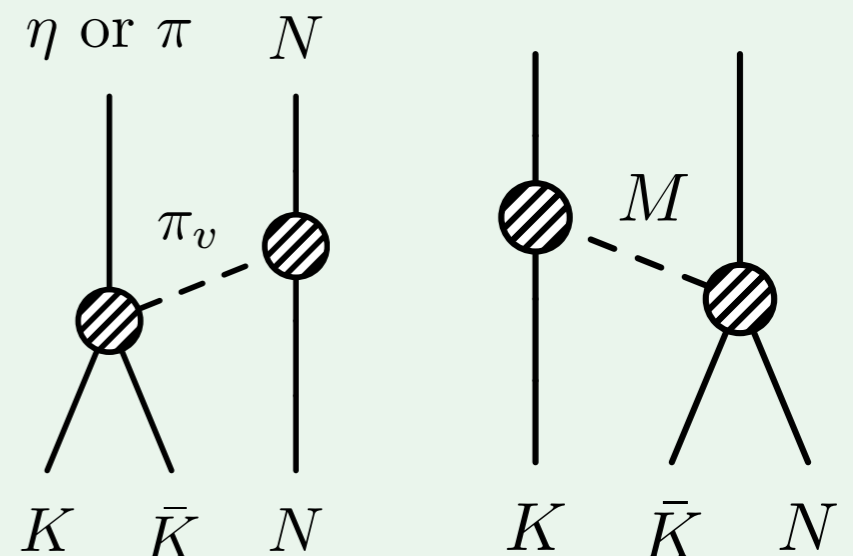
(a) contact interaction

- three particles at a point \rightarrow (density)²
- suppose radius of typical quark-model-like resonance to be 0.8 fm
- suppression factor: $(0.8 \text{ fm} / 1.7 \text{ fm})^6 \sim 0.01$



(b) virtual meson exchanges

- two-body decays without meson-exchange is impossible due to energy conservation.
- two-body transitions
- suppression factor: $(0.8 \text{ fm} / 1.7 \text{ fm})^3 \sim 0.1$
- finite range of virtual meson exchange



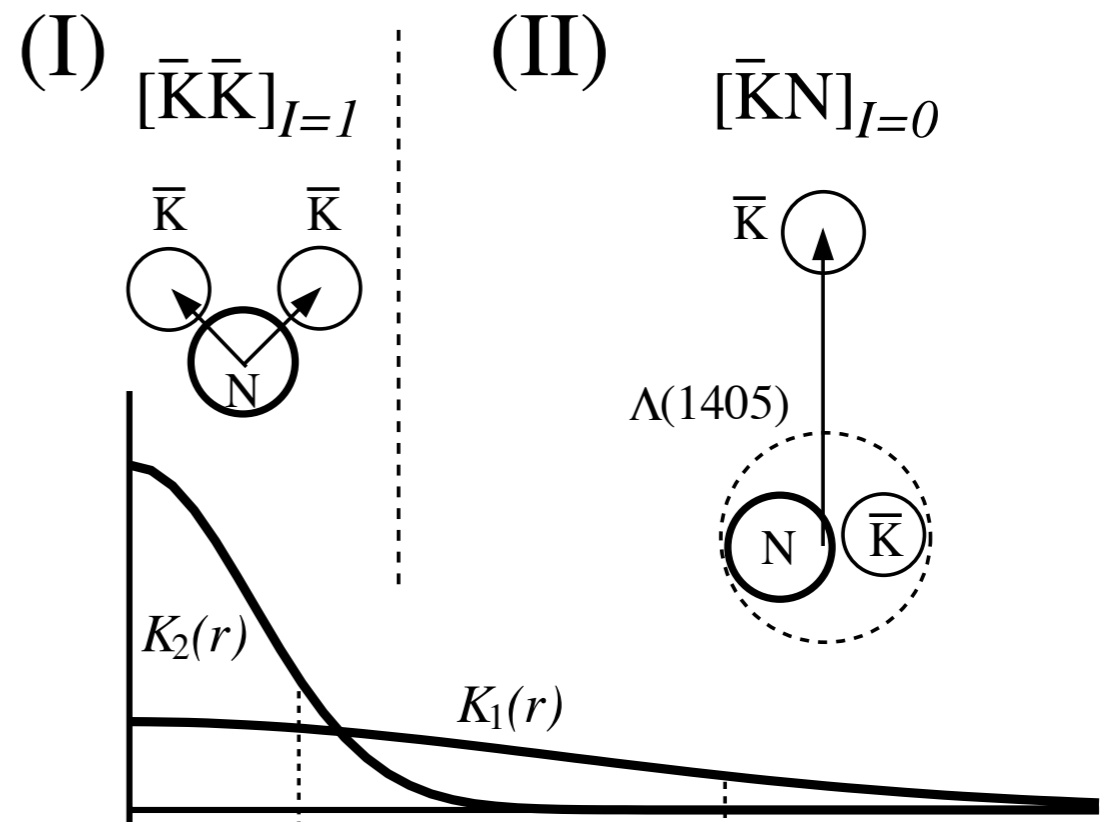
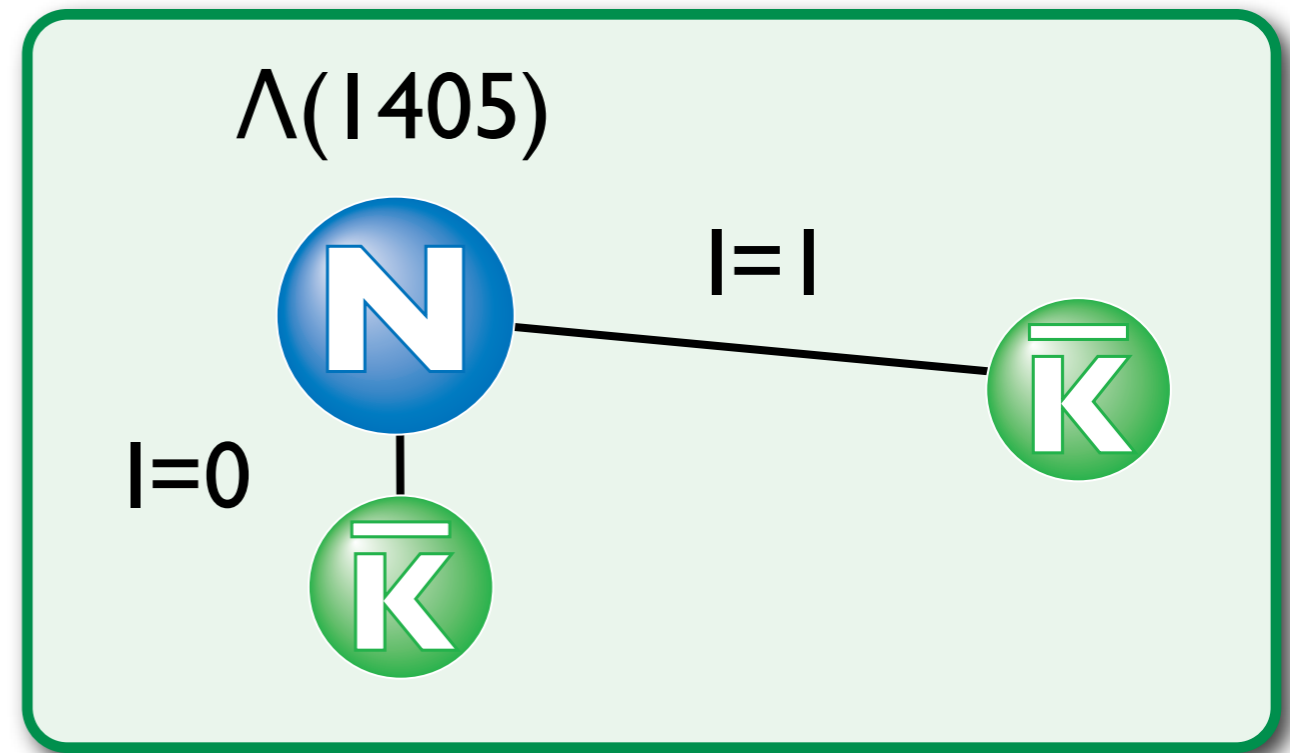
main two-body decay mode: ηN

$\bar{K}\bar{K}N$ system with $I=1/2, J^P=1/2^+$ Ξ^*

Once $\Lambda(1405)$ forms in a $K^{\text{bar}}N$ system with $I=0$, another K^{bar} and N has dominantly $I=1$ component, which is weak attraction. This is not enough to overcome the repulsive $K^{\text{bar}}K^{\text{bar}}$ interaction.

very weak binding

binding energy ~ 2 MeV



Y. Kanada-En'yo, DJ, PRC78, 025212 (2008)

Summary

hadronic molecular state

- system of multiple hadrons described by hadron dynamics
- weakly bound and spatially extended system

$\Lambda(1405)$, $f_0(980)$, $a_0(980)$

explore possibilities for three-body HMS

- non-relativistic potential model
- $KK^{\text{bar}}N$ system as $N^*(1910)$

binding energy : 20~40 MeV, width : 90~100 MeV
radius : 1.7 fm
interhadron distance : ~ 2 fm



keeping properties of the subsystems

- main decay modes
- | | |
|----------------|----------------------|
| $\pi \Sigma K$ | from $\Lambda(1405)$ |
| $\pi \eta N$ | from $a_0(980)$ |
| ηN | small two-body decay |

Double pole structure of $\Lambda(1405)$

DJ, Oller, Oset, Ramos, Meissner
NPA725, 181 ('03)

$\Lambda(1405)$ is a superposition of two states.

there are two attractive channels

group theoretically

SU(3) singlet and octet

physically

$\bar{K}N$, $\pi\Sigma$

$\Lambda(1405)$

below threshold of $\bar{K}N$

pole 1 : 1390 - 66i

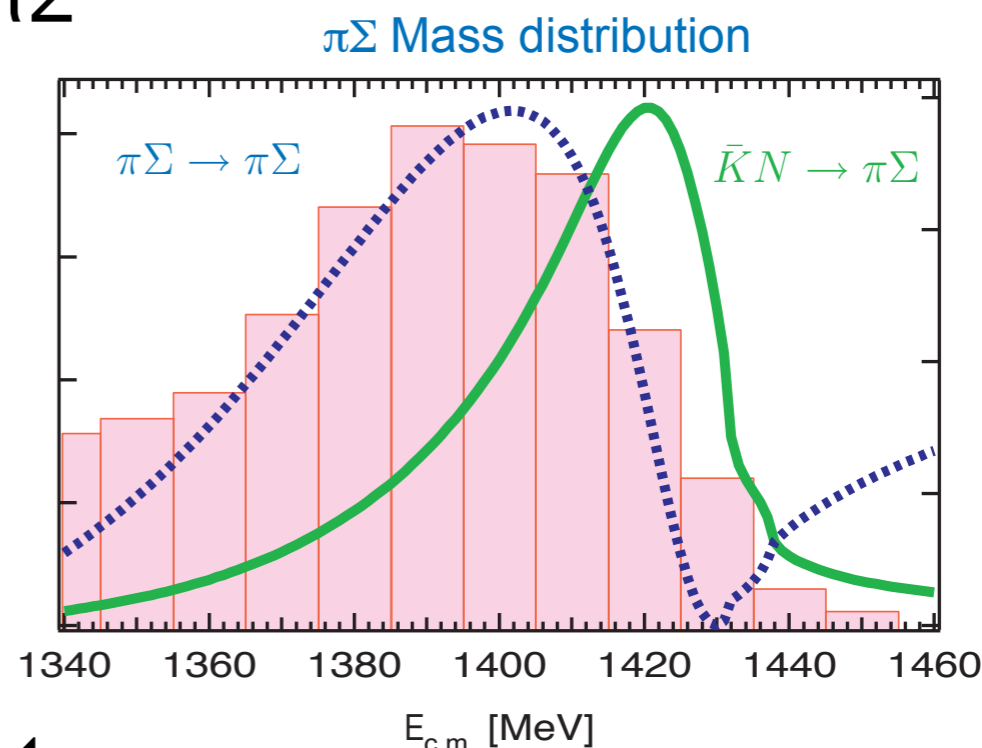
- wider width

- strongly couples to $\pi\Sigma$ state

pole 2 : 1426 - 16i

- narrower width

- dominantly couples to $\bar{K}N$ state



Implication of double pole structure

$\Lambda(1405)$ spectrum is dependent on channels

Resonance position in $\bar{K}N$ channel ~ 1420 MeV with narrower width
not 1405 MeV

This 15 MeV difference is important for $\bar{K}N$ interactions

Double pole structure of $\Lambda(1405)$

DJ, Oller, Oset, Ramos,
Meissner, NPA725, 181 ('03)

$\Lambda(1405)$ is a superposition of two states **having different properties.**

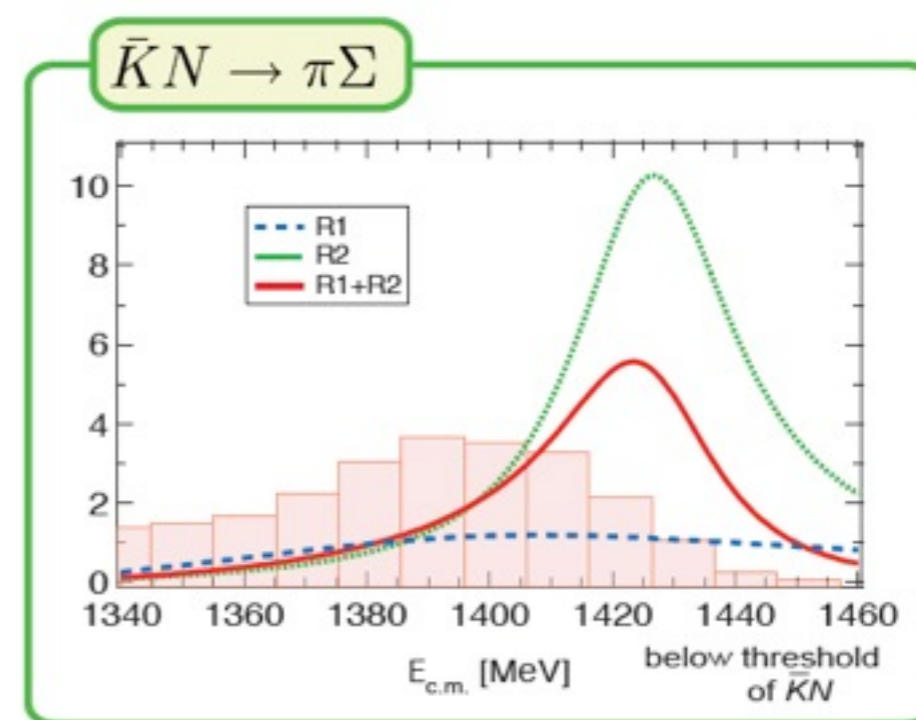
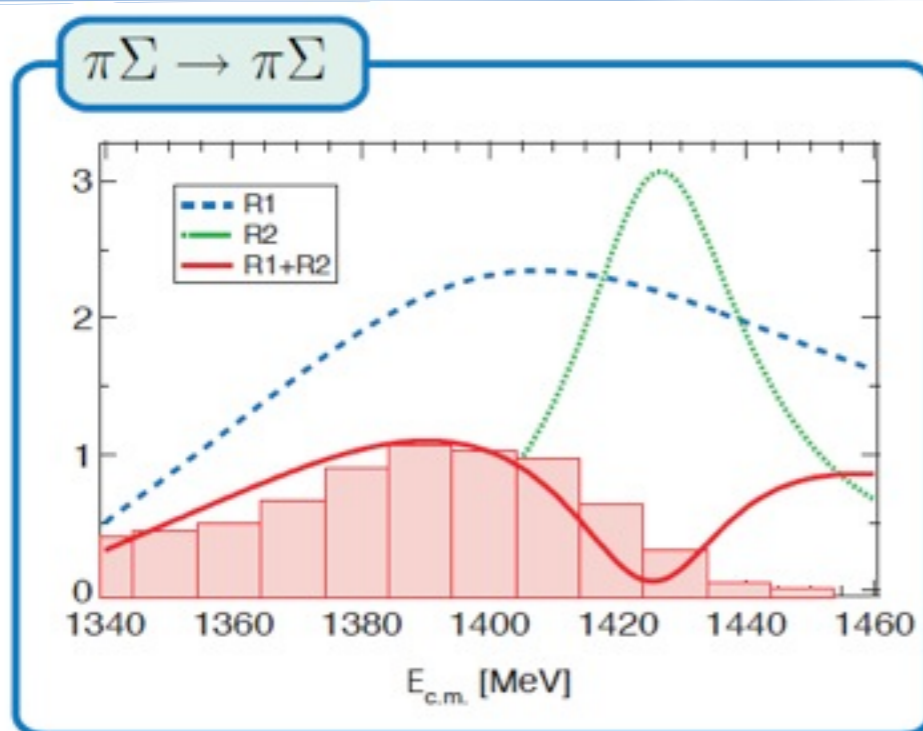
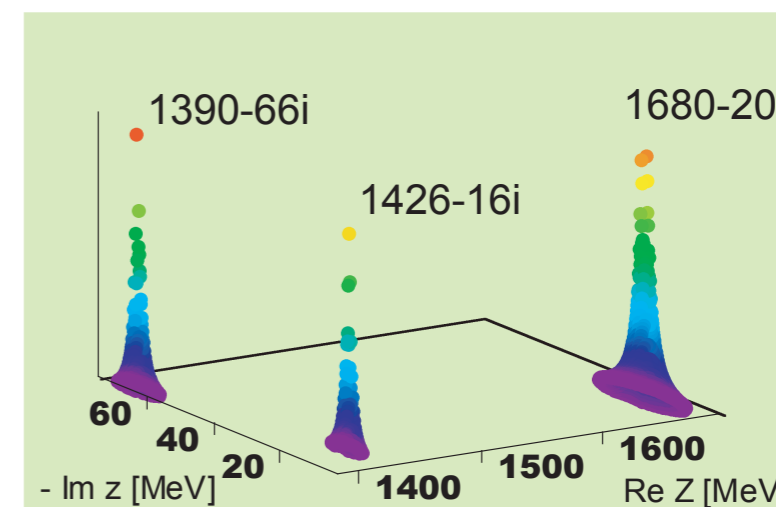
pole 1: **1390 MeV, width 132 MeV**

strongly couples to **$\pi\Sigma$ state**

pole 2: **1426 MeV, width 32 MeV**

dominantly couples to **$K^{\text{bar}}N$ state**

Modulus of T-matrix



two attractive channels: $K^{\text{bar}}N$ and $\pi\Sigma$

Hyodo, Weise, PRC77, 035204 ('08)

essentially, pole 1: **$\pi\Sigma$ resonance**, pole 2: **$K^{\text{bar}}N$ bound state**

Subthreshold properties of $K^{\text{bar}}N$

DJ, Oset, Sekihara, in progress

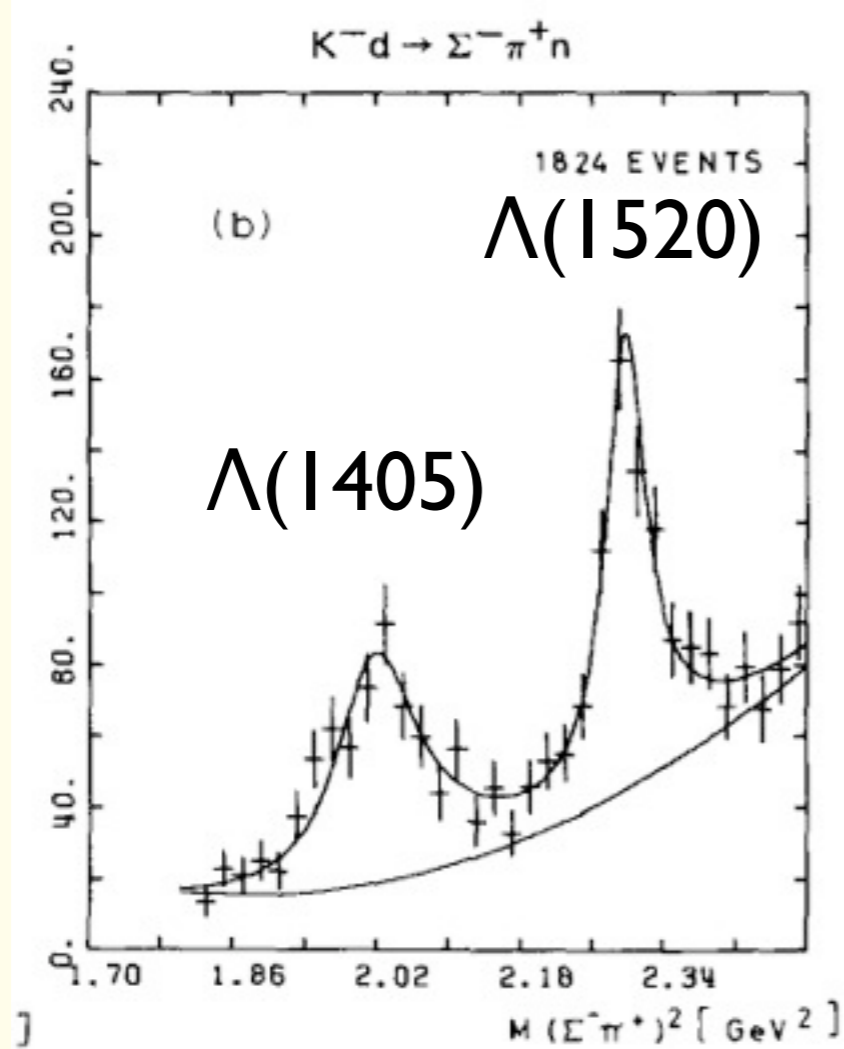
$\Lambda(1405)$ spectra in $K^{\text{bar}}N$ channel

$\bar{K}N \rightarrow \Lambda(1405)$

$$K^- d \rightarrow \Lambda(1405)n$$

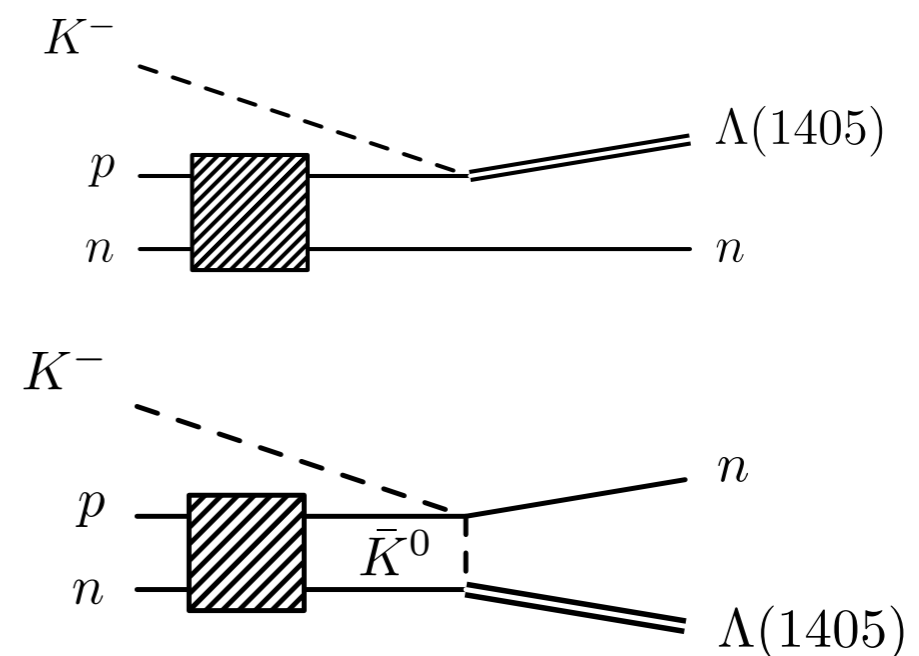
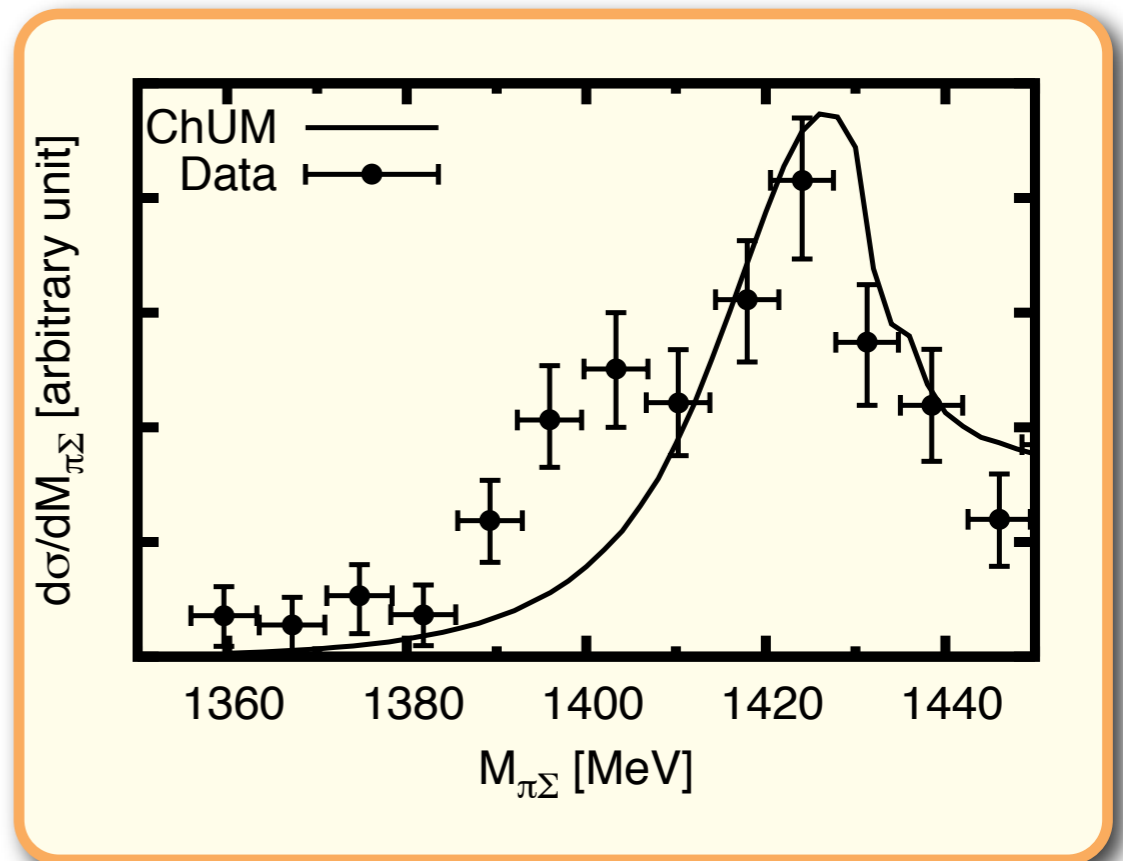
Braun et al. NPBI29,1,('77)

bubble chamber



$$K^- d \rightarrow \pi^+ \Sigma^- n$$

peak position 1420 MeV



Structure of Baryon Resonances

Baryon resonances : decay with strong interactions

in understanding the structure of baryon resonances

dynamical aspect (*hadron dynamics*)

decaying resonance \rightarrow large hadronic components

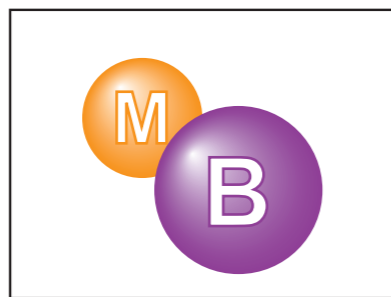
hadron dynamics is important

symmetry aspect (*quark dynamics*)

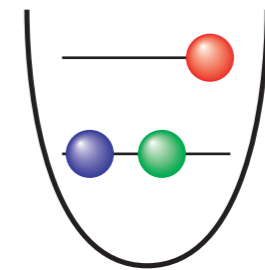
symmetry in terms of quarks

ex. chiral partners: N(1535) chiral partner of nucleon ??

cluster picture



shell model picture



- difference in range of dynamics
- meson cloud effects

Effective interactions

$$\bar{K}N \quad V_{\bar{K}N} = U_{\bar{K}N}^{I=0} \exp[-(r/b)^2] + U_{\bar{K}N}^{I=1} \exp[-(r/b)^2]$$

Hyodo-Weise potential (HW-HNJH)

PRC77,035204 (08)

derived from chiral unitary approach

energy dependent, but small in energy of interest

resonance position ~ **1420 MeV** (double pole structure)

interaction range $b = 0.47$ fm

Akaishi-Yamazaki potential (AY)

PRC64,044005 (02)

obtained phenomenologically

$I=0$: reproduce $\Lambda(1405)$ as quasi-bound state of $K^{\text{bar}}N$

mass: **1405 MeV**, width: 40 MeV

$I=1$: scattering and Konic atom data

interaction range $b = 0.66$ fm

$\Lambda(1405)$ as Quasi-bound state of $\bar{K}N$

$\Lambda(1405)$ $J^\pi = 1/2^-$, $I = 0$, $S = -1$, $Q = 0$

- most established resonance, seen in many exp. clearly
- mass : 1406.5 ± 4.0 MeV, (below K^-p threshold)
- width : 50 ± 2 MeV (PDG)
- decay mode $\Lambda(1405) \rightarrow (\Sigma\pi)_{I=0}$ 100 % **s-wave**

$\bar{K}N$	<u>1435 MeV</u>	$\Lambda(1405)$
$\pi\Sigma$	<u>1331 MeV</u>	

chiral unitary model

model parameters tuned so as to

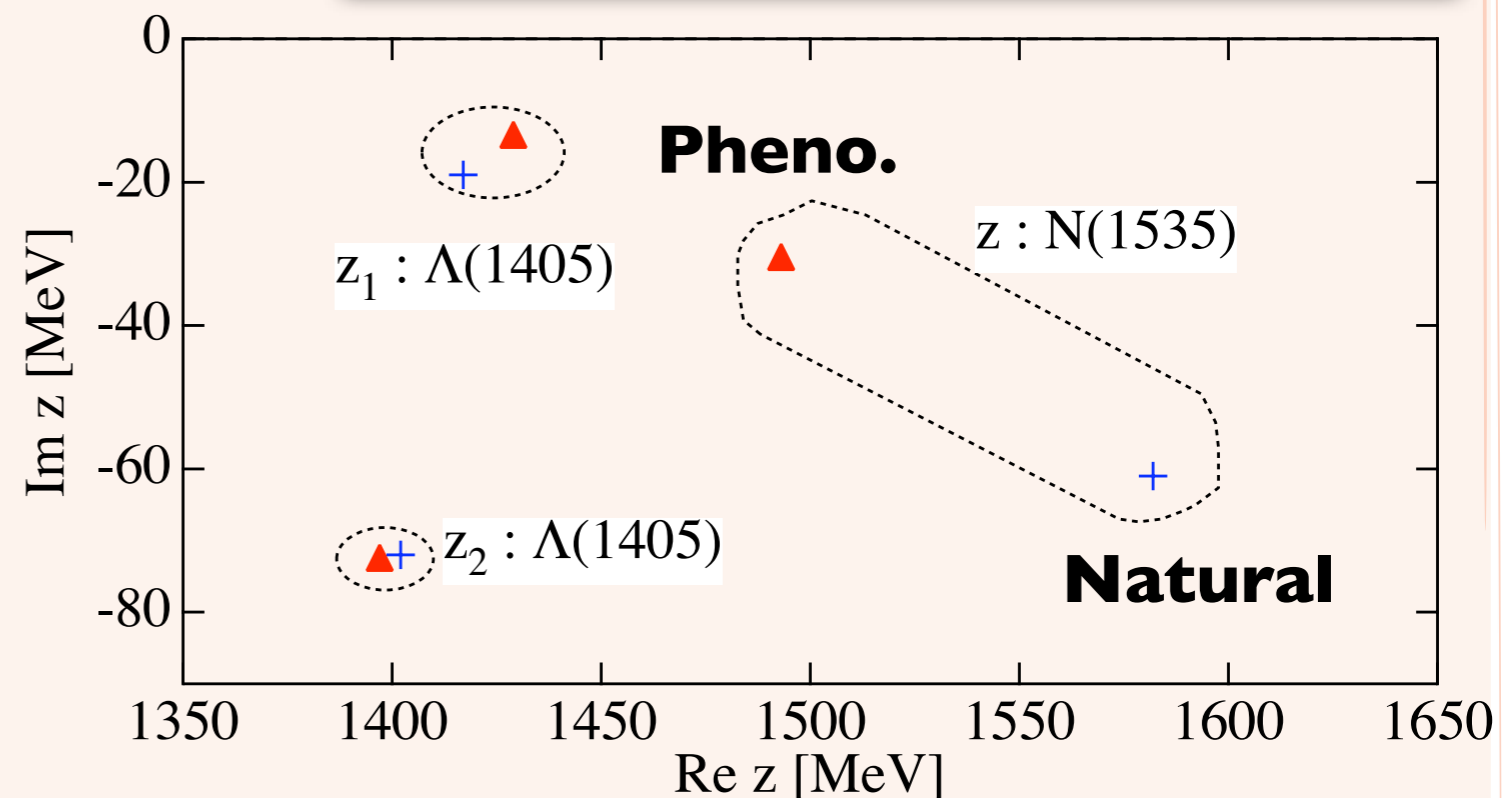
- exclude quark-originated states theoretically

+ Natural

- reproduce scattering data

▲ Pheno.

Hyodo, Jido, Hosaka, PRC78, 025203 ('08)



$\Lambda(1405)$ has mostly meson-baryon components.