

Left-right asymmetry for pion and kaon production in SIDIS process

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1. Introduction
2. Theoretical description
3. Numerical calculation
4. Summary

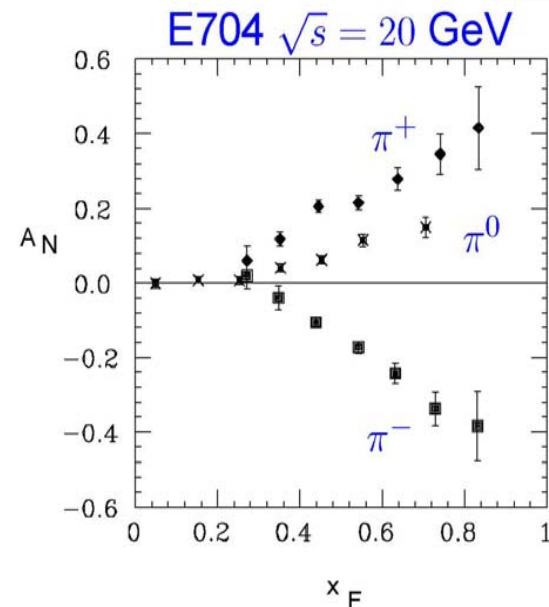
Left-right asymmetry

- several experiments have been carried to measure the “single spin asymmetry A_N ”, large asymmetry have been observed, whereas the theoretical prediction lead to $A_N = 0$.
- Left-right asymmetry reported by E704 experiment:

$$p^\uparrow p \rightarrow \pi X$$

$$A = - \frac{1}{P \cos \phi} \frac{N_\uparrow - N_\downarrow}{N_\uparrow + N_\downarrow}$$

Adama D, et al. Phys. Lett. B, 1991, 264: 462-466.



Theoretical understanding

- To describe the left-right asymmetry, two mechanisms are available.

Sivers effect: originated from quark intrinsic transverse momentum.

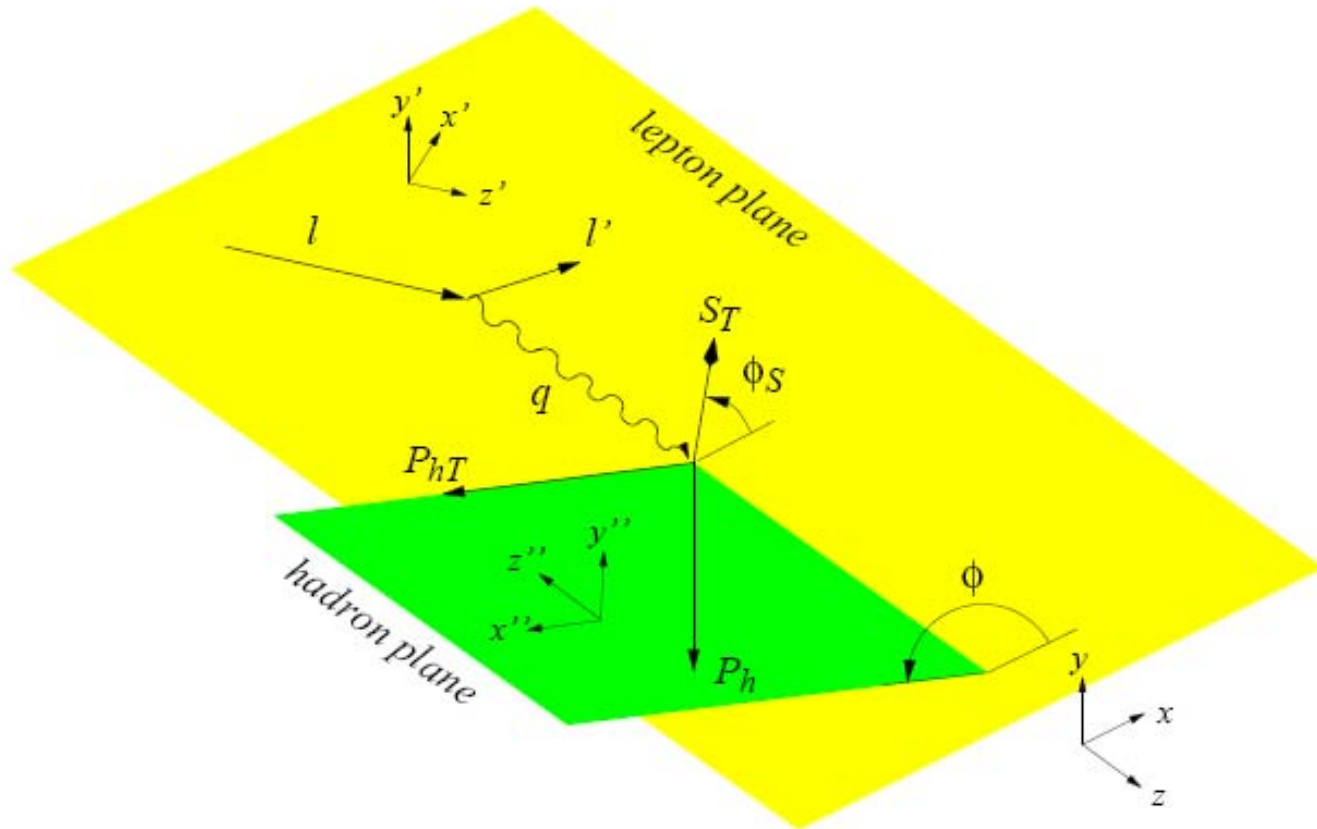
Sivers, Phys. Rev. D 43, 261 (1991).

Collins effect: generate the asymmetry in hadronization processes.

Collins, Nucl. Phys. B 396, 191 (1993).

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Coordinate system



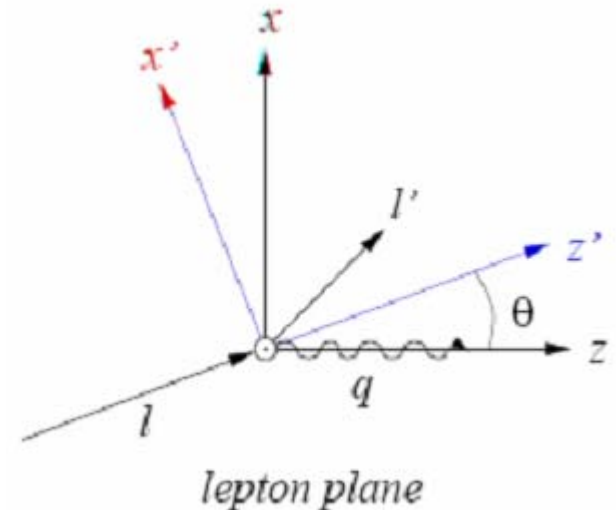
lp and $\gamma^* p$ frame

Transform from $\gamma^* p$ fram to lp fram

$$A_{UT}^\ell = \frac{\cos \theta A_{UT}^{\gamma^*} + \sin \theta \cos \phi_S A_{UL}^{\gamma^*}}{\sqrt{1 - \sin^2 \theta \sin^2 \phi_S^\ell}}$$

$$\sin \theta = \gamma \sqrt{\frac{1 - y - \frac{1}{4} y^2 \gamma^2}{1 + \gamma^2}},$$

$$\gamma = \frac{2xM_p}{Q}$$



•An estimation of θ :

HERMES: $\langle x \rangle = 0.09$, $\langle y \rangle = 0.54$, $\langle Q^2 \rangle = 2.41 \text{ GeV}^2$, $\langle \sin \theta \rangle = 0.073$

JLab(6GeV): $\langle x \rangle = 0.23$, $\langle y \rangle = 0.6$, $\langle Q^2 \rangle = 1.8 \text{ GeV}^2$, $\langle \sin \theta \rangle = 0.19$

JLab(12GeV): $\langle x \rangle = 0.23$, $\langle y \rangle = 0.57$, $\langle Q^2 \rangle = 2.5 \text{ GeV}^2$, $\langle \sin \theta \rangle = 0.17$

Indeed small

Cross section

Diehl and Sapeta, Eur. Phys. J. C 41, 515 (2005).

Bacchetta, et al., JHEP 0702:093 (2007).

$$\begin{aligned} \frac{d\sigma}{dx dy d\phi^\ell dz d\phi_h dP_{h\perp}^2} &= \frac{\alpha^2}{2sx(1-\varepsilon)} \frac{\cos\theta}{1 - \sin^2\theta \sin^2\phi_s^\ell} \times \left\{ F[f_1 D_1] \right. \\ &\quad - \frac{S_T \cos\theta}{\sqrt{1 - \sin^2\theta \sin^2\phi_s^\ell}} \sin(\phi_h^\ell - \phi_s^\ell) F \left[\frac{\vec{h} \cdot \vec{p}_\perp}{M_p} f_{1T}^\perp D_1 \right] \\ &\quad \left. - \frac{S_T \cos\theta}{\sqrt{1 - \sin^2\theta \sin^2\phi_s^\ell}} \sin(\phi_h^\ell + \phi_s^\ell) F \left[\frac{\vec{h} \cdot \vec{k}_\perp}{M_h} h_1 H_1^\perp \right] \right\} \\ &= d\sigma_{UU} + d\sigma_{Siv} + d\sigma_{Col} \end{aligned}$$

$$F[\omega fD] = \sum_a e_a^2 \int d^2\vec{p}_\perp d^2\vec{k}_\perp \delta^2(\vec{p}_\perp - \vec{k}_\perp - \frac{\vec{P}_{h\perp}}{z}) \omega(\vec{p}_\perp, \vec{k}_\perp) f_a(x, \vec{p}_\perp^2) D_a(z, z^2 \vec{k}_\perp^2)$$

Conventional treatment

- Deconvolution

$$W = \begin{cases} \frac{|P_{h\perp}|}{M_N} & : \text{Sivers} \\ \frac{|P_{h\perp}|}{M_h} & : \text{Collins} \end{cases}$$

$$\int d^2 P_{h\perp} \left(\left[\frac{|P_{h\perp}|}{M_h} \right] \cdot d\sigma_C \right) = -\frac{2\alpha^2}{sxy^2} A(y) \sum_q e_q^2 z h_1^q(x) H_1^{\perp(1)q}(z)$$

A factorized expression in x and z obtained

- Projection

$$\int d\phi \cdot \sin(\phi \pm \phi_s) \dots = 0$$

$$\langle w \cdot \sin(\phi - \phi_s) \rangle_{UT} \propto d\sigma_{Siv} \frac{1}{2} (2\pi)^2$$

$$\langle w \cdot \sin(\phi + \phi_s) \rangle_{UT} \propto d\sigma_{Col} \frac{1}{2} (2\pi)^2$$

- Why do we choose $\sin(\Phi - \Phi_s)$ as weighting function for Sivers effect, and $\sin(\Phi + \Phi_s)$ as weighting function for Collins effect.
- The choice of the weighting functions show **a bias on a certain theory.**

An universal way

J. She, Y. Mao, B.-Q. Ma, Phys. Lett. B 666, 355 (2008).

- The weighting functions are not put into use.
- **Only** the hadrons produced **within a range**, rather than the whole space are selected.

e.g., $-\frac{\pi}{4} \leq \phi_h \leq \frac{\pi}{4}$, left; $\frac{3\pi}{4} \leq \phi_h \leq \frac{5\pi}{4}$, right.

$$\begin{aligned} A_{LR} &= \frac{1}{S_T} \frac{\int d\phi^\ell d\phi_h dP_{h\perp}^2 (d\sigma(\phi_S) - d\sigma(\phi_S + \pi))}{\int d\phi^\ell d\phi_h dP_{h\perp}^2 (d\sigma(\phi_S) + d\sigma(\phi_S + \pi))} \\ &= \frac{1}{S_T} \frac{\int d\phi^\ell d\phi_h dP_{h\perp}^2 (d\sigma(\phi_h) - d\sigma(\phi_h + \pi))}{\int d\phi^\ell d\phi_h dP_{h\perp}^2 (d\sigma(\phi_h) + d\sigma(\phi_h + \pi))} \end{aligned}$$

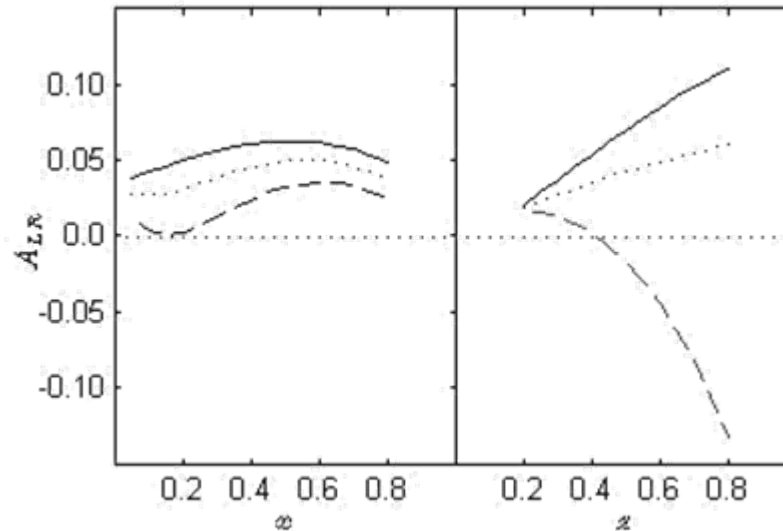
- Integral over Φ^1 ,

$$\frac{1}{2\pi} \int_0^{2\pi} d\phi^\ell \frac{\sin(\phi_h^\ell - \phi_s^\ell)}{(1 - \sin^2 \theta \sin^2 \phi_s^\ell)^{3/2}} = \sin(\phi_h - \phi_s) \left(1 + \frac{3}{4} \sin^2 \theta + o(\sin^4 \theta)\right)$$

$$\frac{1}{2\pi} \int_0^{2\pi} d\phi^\ell \frac{\sin(\phi_h^\ell + \phi_s^\ell)}{(1 - \sin^2 \theta \sin^2 \phi_s^\ell)^{3/2}} = -\sin(\phi_h - \phi_s) \left(\frac{3}{8} \sin^2 \theta + o(\sin^4 \theta)\right)$$

Sivers effect is $o(1)$, but other terms such as the Collins effect are $o(\sin^2 \theta)$. So Sivers effect is dominant and other effects are suppressed.

- Unweighted Asymmetry



- only u and d quark are considered here

- We reanalyze the pions production extend our calculation on the kaons production using the latest parameterization which contain information of both valance and sea quarks.
- Collins effect is not known so clearly yet and is not considered in our analysis as it is strongly suppressed

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Sivers function

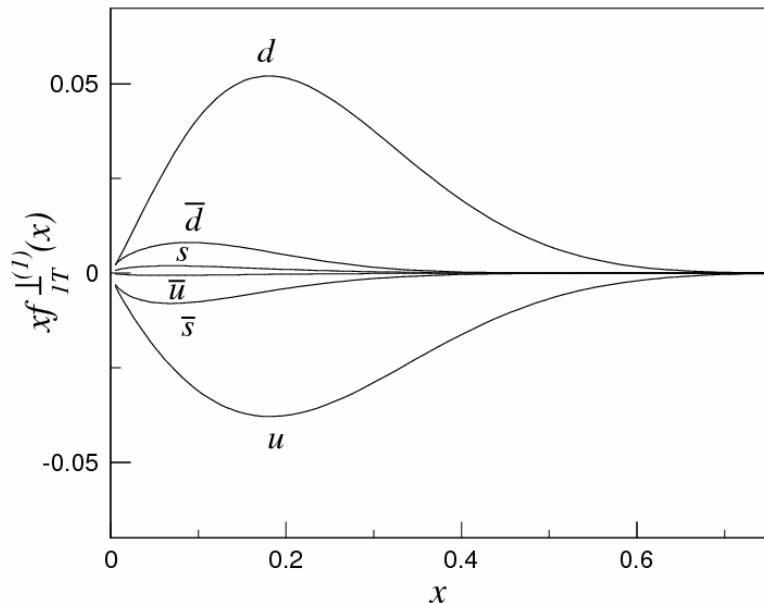
- Parameterization for Sivers function:

$$f_{1T}^{\perp q}(x, p_{\perp}^2) = -\frac{M_p}{p_{\perp}} \mathcal{N}_q(x) f_q(x) g(p_{\perp}^2) h(p_{\perp}^2),$$

$$\mathcal{N}_q(x) = N_q x_q^{\alpha} (1-x)_q^{\beta} \frac{(\alpha_q + \beta_q)^{(\alpha_q + \beta_q)}}{\alpha_q^{\alpha_q} \beta_q^{\beta_q}},$$

$$g(p_{\perp}^2) = \frac{e^{-p_{\perp}^2 / \langle p_{\perp}^2 \rangle}}{\pi \langle p_{\perp}^2 \rangle}, \quad h(p_{\perp}^2) = \sqrt{2} e \frac{p_{\perp}}{M'} e^{-p_{\perp}^2 / \langle M'^2 \rangle}.$$

$N_u = 0.35_{-0.08}^{+0.08}$	$N_d = -0.90_{-0.10}^{+0.43}$	$N_s = -0.24_{-0.50}^{+0.62}$
$N_{\bar{u}} = 0.04_{-0.24}^{+0.22}$	$N_{\bar{d}} = -0.40_{-0.44}^{+0.33}$	$N_{\bar{s}} = 1_{-0.0001}^{+0}$
$\alpha_u = 0.73_{-0.58}^{+0.72}$	$\alpha_d = 1.08_{-0.65}^{+0.82}$	$\alpha_{sea} = 0.79_{-0.47}^{+0.56}$
$\beta = 3.46_{-2.90}^{+4.87}$	$M_1^2 = 0.34_{-0.16}^{+0.30} (GeV/e)^2$	

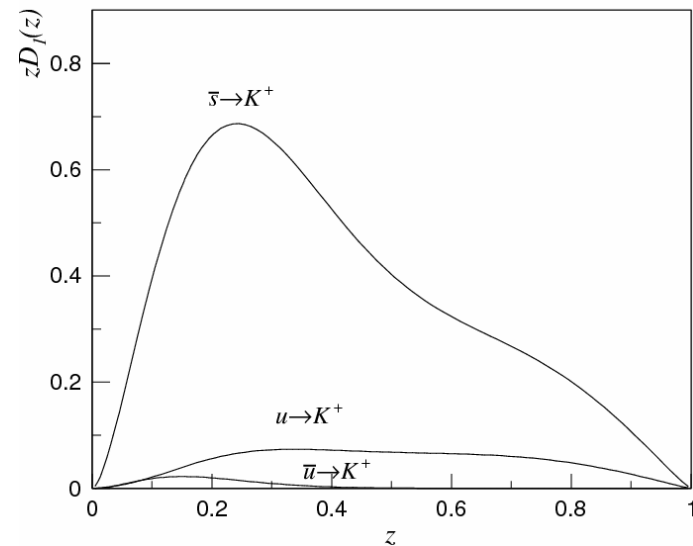
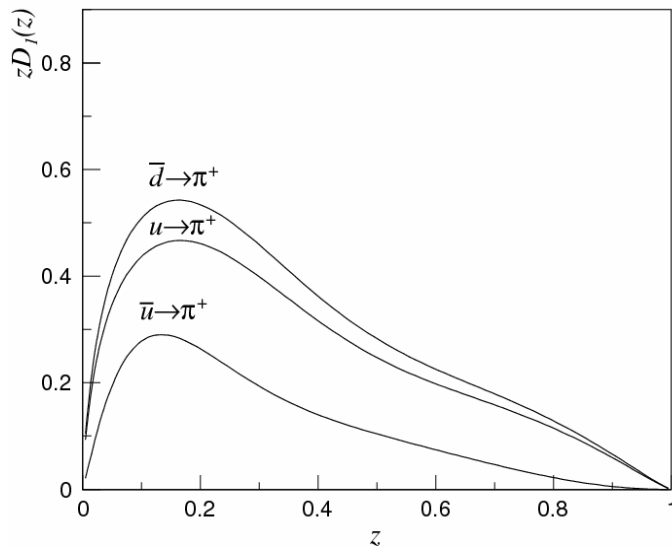


M. Anselmino et al., Eur. Phys. J. A 39 89 (2009).

Fragmentation function

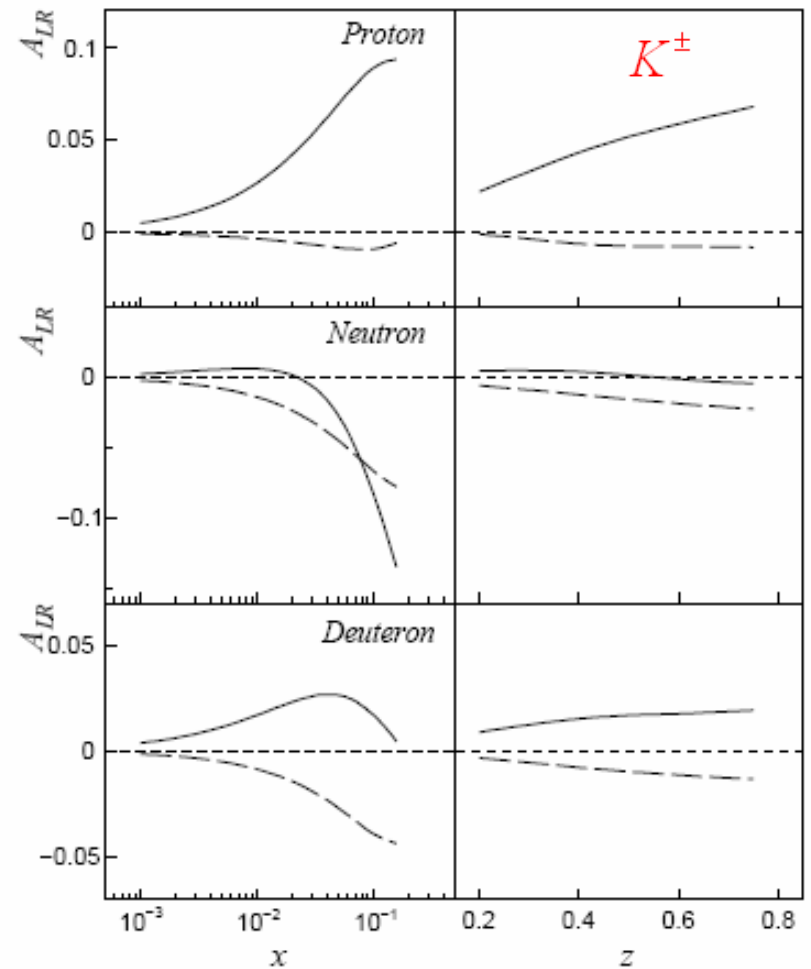
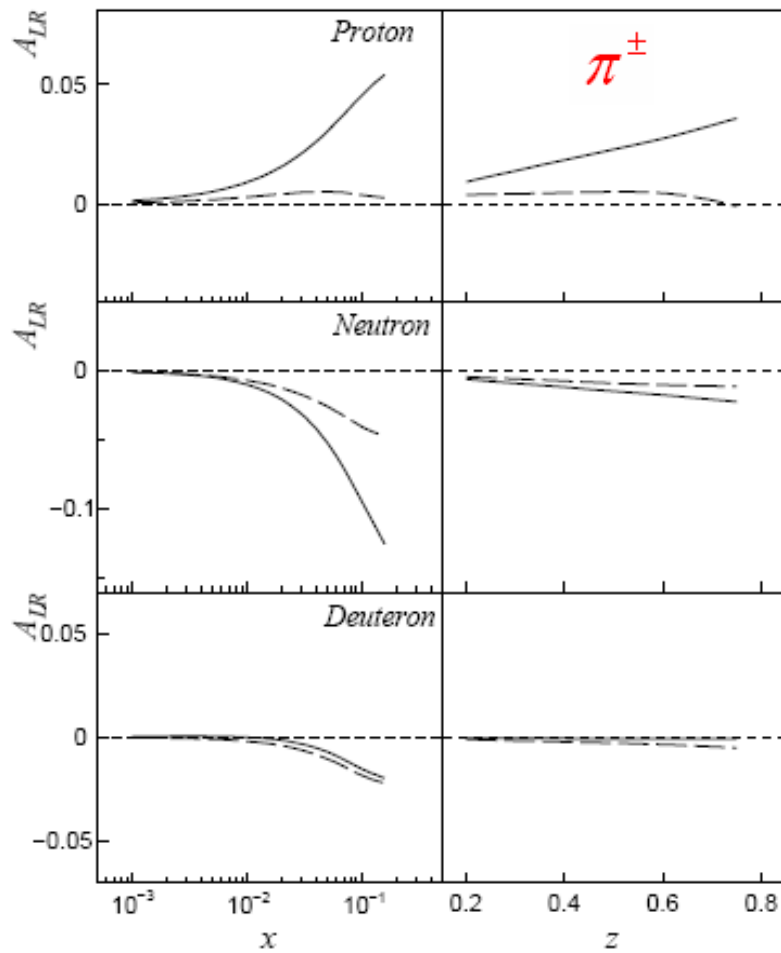
- Parameterization for fragmentation function:

$$D_i^H(z, \mu_0) = \frac{N_i z^{\alpha_i} (1-z)^{\beta_i} [1 + \gamma_i (1-z)^{\delta_i}]}{\text{B}[2 + \alpha_i, \beta_i + 1] + \gamma_i \text{B}[2 + \alpha_i, \beta_i + \delta_i + 1]}$$

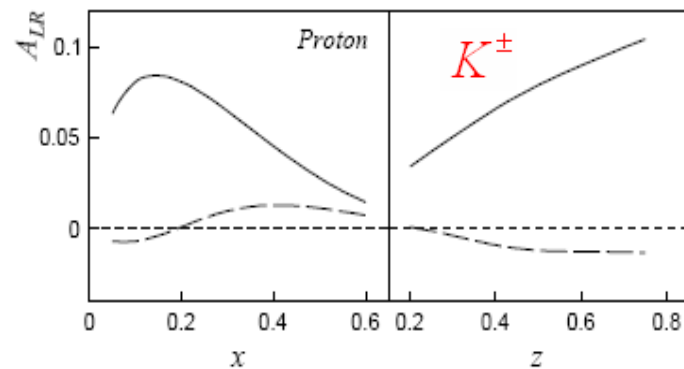
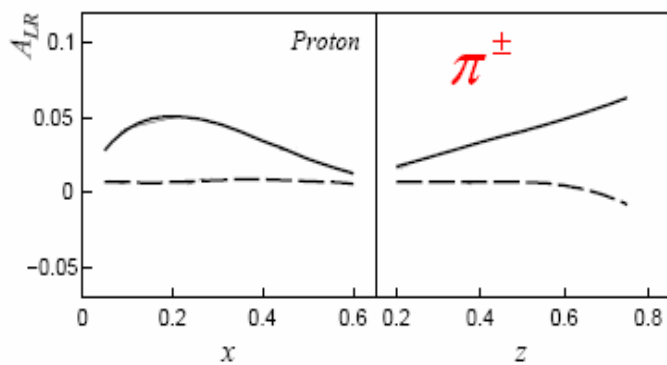


D. de Florian, R. Sassot, Phys. Rev. D 75, 114010 (2007).

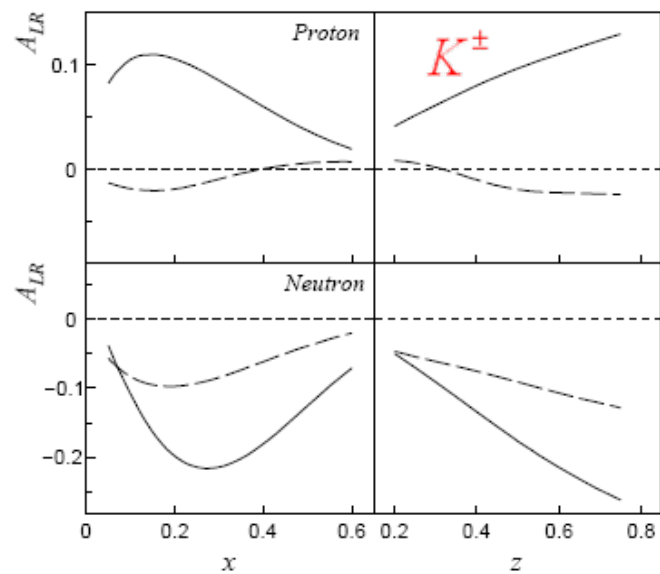
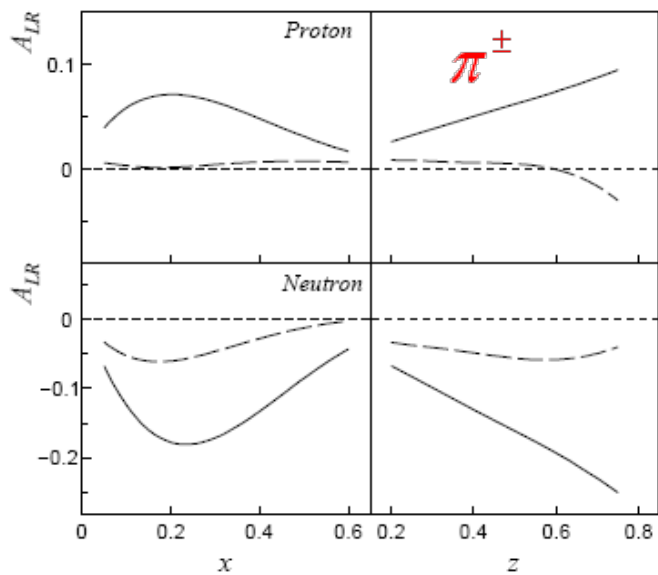
Results on COMPASS kinematics



HERMES



JLab



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Summary

- E704 method can be used in SIDIS process. Asymmetry can be obtained in this way.
- We provide the prediction for pions and kaons production on HERMES, JLab and COMPASS kinematics on different targets. Be caution that results for K^- production may not accurate enough.
- With the current understanding, Sivers effect dominant here, Collins effect is suppressed.

We suggest relevant collaboration apply data analysis under this way to provide more information for theoretical studies.