

Complementary studies on N^* from e^+e^- , pp and $\bar{p}p$ collisions

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Outline:

- **Introduction**
- **N^* from $e^+e^- \rightarrow \psi \rightarrow \bar{N} N^*$**
- **N^* from $pp \rightarrow N N^*$**
- **N^* from $\bar{p}p \rightarrow \bar{N} N^*$**
- **Summary and Prospects**

1. Introduction

Two well-known problems for the classical 3q model

- Mass order reverse problem for the lowest excited baryons

uud (L=1) $1/2^- \sim N^*(1535)$ **should be the lowest**

uud (n=1) $1/2^+ \sim N^*(1440)$

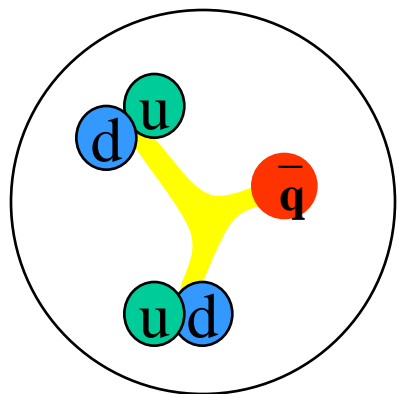
uds (L=1) $1/2^- \sim \Lambda^*(1405)$

harmonic oscillator ($2n + L + 3/2$) $\hbar\omega$

- The number of predicted states is much less than observed
“missing” baryon states : non-existence / to be observed ?

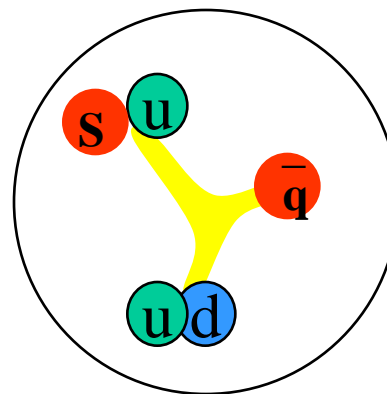
A new scheme to explain the mass reverse problem

A.Zhang et al, HEPNP29(2005)250 ; B.S.Zou, EPJA35(2008)325



$$\bar{q} \quad 1/2^+$$

$$\left. \begin{array}{l} [ud] \\ [ud] \end{array} \right\} L=1$$



$$\bar{q} \quad 1/2^-$$

$$\left. \begin{array}{l} [ud] \\ [us] \end{array} \right\} L=0$$

$$N^*(1535) \sim uud (L=1) + \varepsilon [ud][us] \bar{s} + \dots$$

$$N^*(1440) \sim uud (n=1) + \xi [ud][ud] \bar{d} + \dots$$

$$\Lambda^*(1405) \sim uds (L=1) + \varepsilon [ud][su] \bar{u} + \dots$$

$N^*(1535)$: $[ud][us] \bar{s} \rightarrow$ larger coupling to $N\eta, N\eta', N\phi$ & $K\Lambda$, weaker to $N\pi$ & $K\Sigma$, and heavier !

The new picture for the $1/2^-$ nonet predicts:

$$\Lambda^* \quad [us][ds] \quad \bar{s} \quad \sim \quad 1575 \text{ MeV}$$

$$\Sigma^* \quad [us][du] \quad \bar{d} \quad \sim \quad 1360 \text{ MeV}$$

$$\Xi^* \quad [us][ds] \quad \bar{u} \quad \sim \quad 1520 \text{ MeV}$$

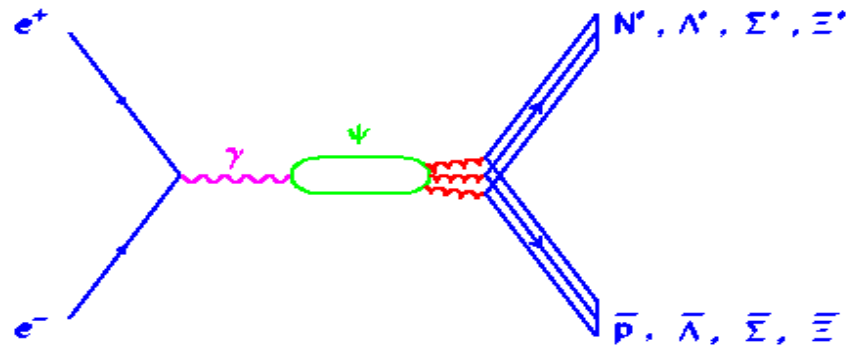
Where to look for “missing” baryon states ?

High statistics γp and $K p$ experiments

+ e^+e^- , pp & $\bar{p}p$ experiments

2. N^* from $e^+e^- \rightarrow \psi \rightarrow \bar{N} N^*$

$$\Psi \rightarrow \bar{B} B M \Rightarrow N^*, \Lambda^*, \Sigma^*, \Xi^*$$

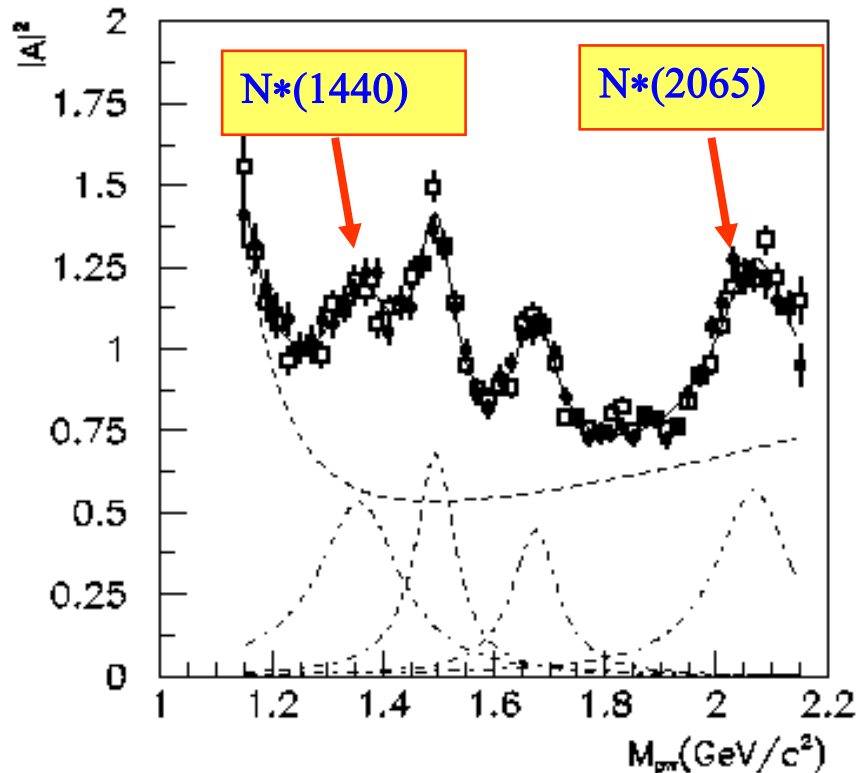


New mechanism for baryon production & an ideal isospin filter

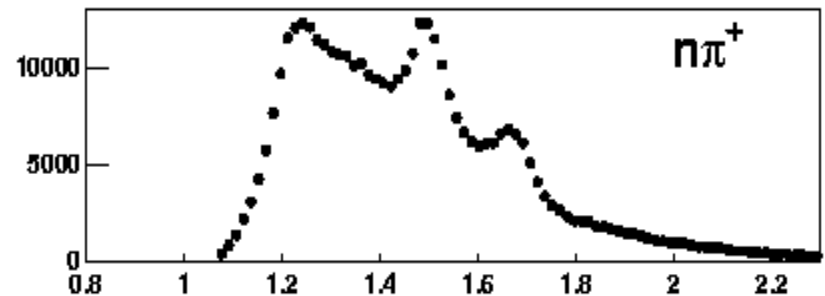
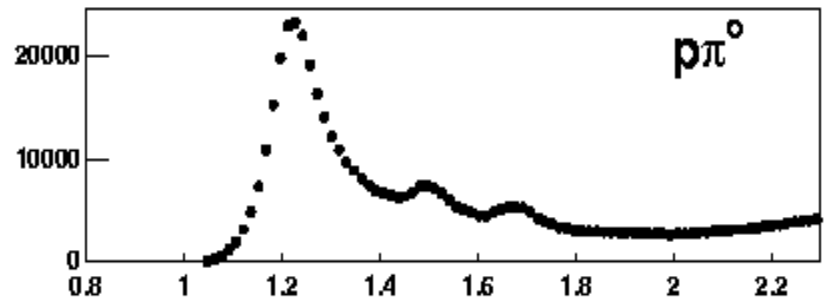
1) N^* in $\psi \rightarrow \bar{p}n\pi^+ \text{ \& \ } \bar{n}\pi^-p$

$J/\psi \rightarrow \bar{p}n\pi^+ \text{ \& \ } \bar{n}\pi^-p$

V. Burkert



CLAS E=4 GeV $ep \rightarrow eX$



Observation of two new N^* peaks in πN mass spectrum

BES, Phys.Rev.Lett.97(2006)062001

The new $N^*(2065)$ peak in πN mass spectrum

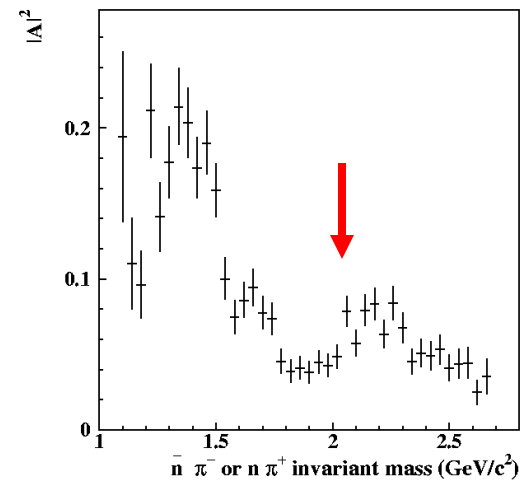
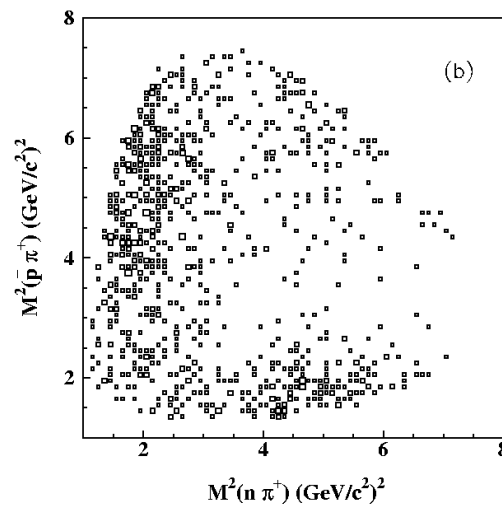
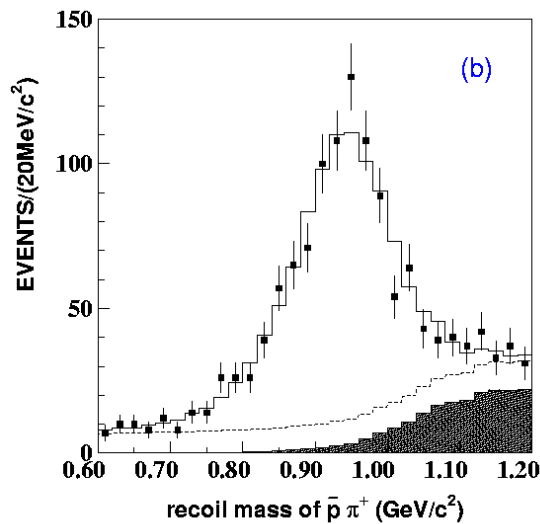
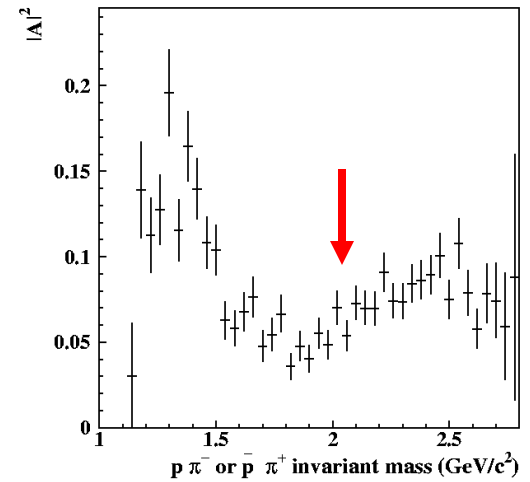
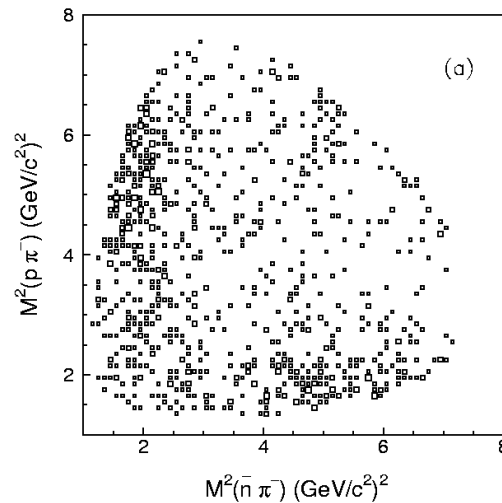
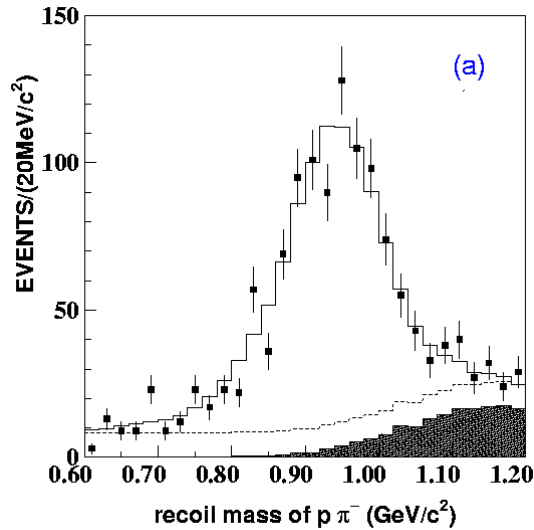
$$M = 2068 \pm 3^{+15}_{-40} \text{ MeV}/c^2, \Gamma = 165 \pm 14 \pm 40 \text{ MeV}/c^2$$

$J/\psi \rightarrow \bar{n} N^*(2065)$ with $L=0$ (small excess energy)
limits its J^P to be $3/2^+$ or $1/2^+$ (spin filter)

PWA with event-based Maximum-Likelihood method
 \rightarrow Both $3/2^+$ and $1/2^+$ are there.

$$\psi(2S) \rightarrow p\pi^-\bar{n} + c.c.$$

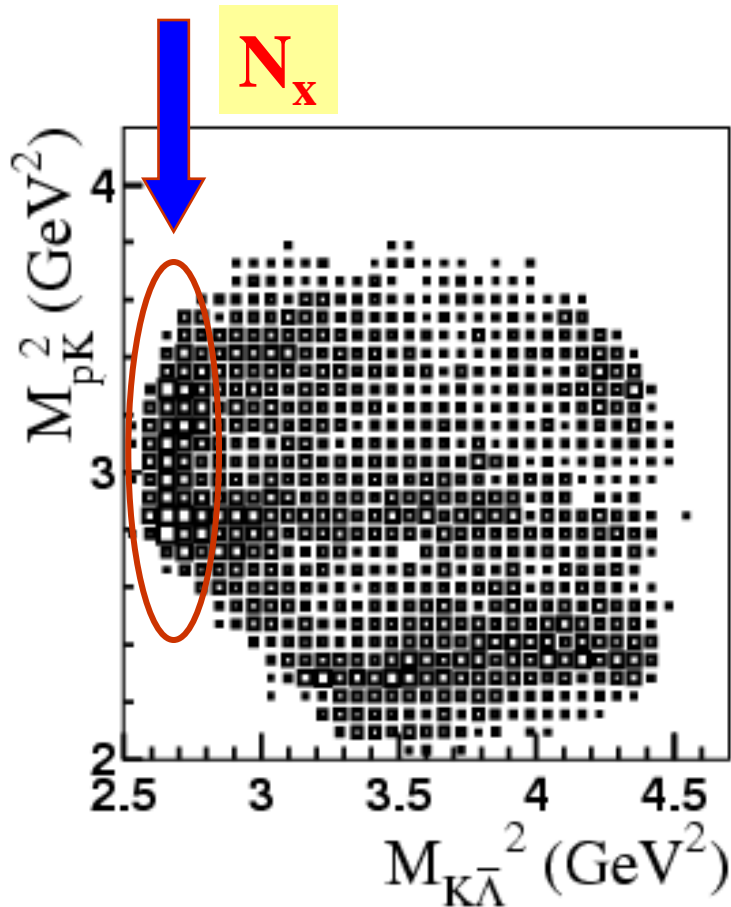
BES, Phys. Rev. D74, 012004 (2006)



Clear $N^*(1440)$ peak and evidence for N^* above 2 GeV

2) N^* in $\psi \rightarrow \bar{\Lambda} K N$

N^* in $J/\psi \rightarrow p K^- \bar{\Lambda} + \text{c.c.}$



Mass **1500~1650MeV**

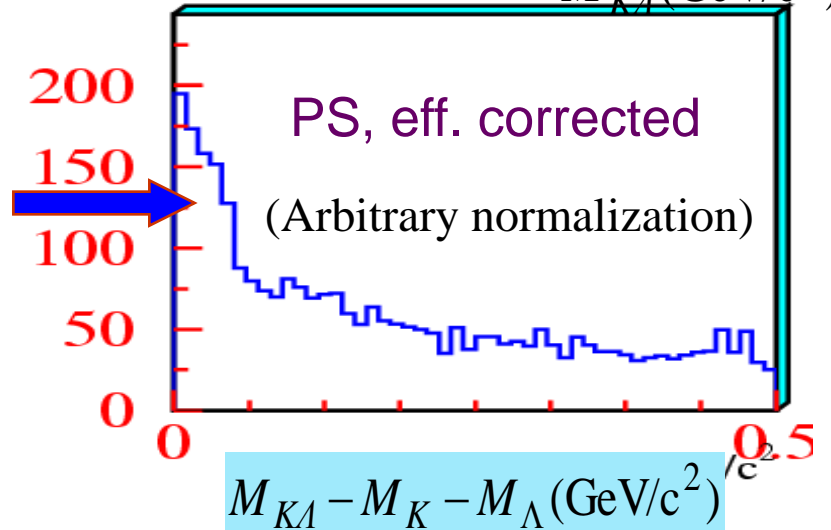
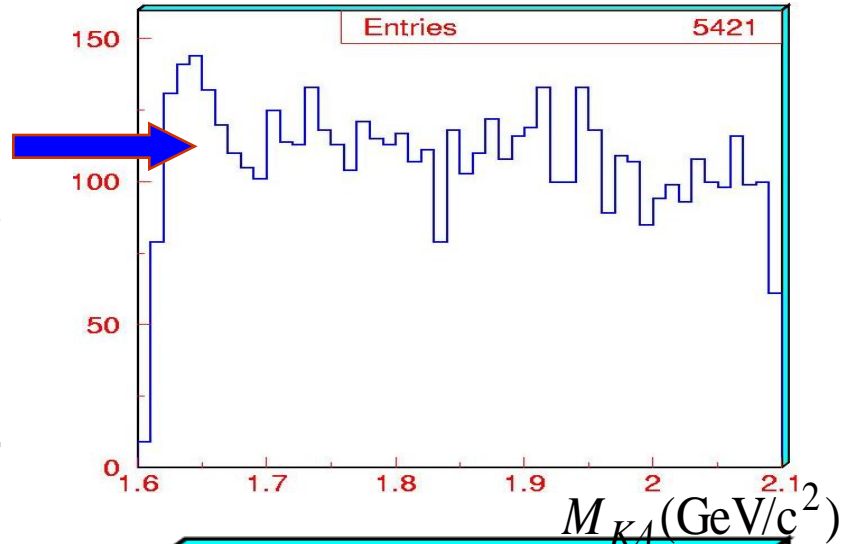
Width **70~110MeV**

J^P **favors 1/2-**

N_x

Events/10MeV

N_x



a) Assuming N_x to be purely $N^*(1535)$:

B.C. Liu, B.S. Zou, PRL96 (2006) 042002; PRL98 (2007) 039102

From relative branching ratios of
 $J/\psi \rightarrow p \bar{N}^* \rightarrow p (K^- \bar{\Lambda}) / p (\bar{p}\eta)$



$$g_{N^*K\Lambda} / g_{N^*p\eta} / g_{N^*N\pi} \sim 2 : 2 : 1$$

b) N_x as dynamical generated with unitary chiral theory:

$N^*(1535)$ + non-resonant part

L.S.Geng, E.Oset, B.S. Zou, M.Doring, PRC79 (2009) 025203

$$g_{N^*K\Lambda} / g_{N^*p\eta} / g_{N^*N\pi} \sim 1.2 : 2 : 1$$

Phenomenology : Large $g_{N^*K\Lambda} \rightarrow$ large $\bar{s}s$ in $N^*(1535)$

$\bar{s}[su][ud]$ or $K\Lambda$ - $K\Sigma$ state

Many more channels:

$J/\psi \rightarrow \bar{p}p\eta$: Phys. Lett. B510 (2001) 75

$J/\psi \rightarrow \bar{\Lambda} K_S^0$: Phys. Lett. B659 (2008) 789

$J/\psi \rightarrow \bar{p}p\omega$: Eur. Phys. J. C53 (2008) 15

$J/\psi \rightarrow \bar{p}p\pi^0$: Hongxun Yang's talk

$\psi(2S) \rightarrow \bar{p}p\pi^0, \bar{p}p\eta$: Phys. Rev. D71 (2005) 072006

J/ψ decay branching ratios (BR*10³)

$p\bar{n}\pi^-$	$p\bar{p}\pi^0$	$p\bar{p}\pi^+\pi^-$	$p\bar{p}\eta$	$p\bar{p}\eta'$	$p\bar{p}\omega$
2.4 ± 0.2	1.1 ± 0.1	6.0 ± 0.5	2.1 ± 0.2	0.9 ± 0.4	1.3 ± 0.3
$\Lambda\bar{\Sigma}^-\pi^+$	$pK^-\bar{\Lambda}$	$pK^-\bar{\Sigma}^0$	$\bar{p}p\phi$	$\Delta(1232)^{++}\bar{p}\pi^-$	$pK^-\bar{\Sigma}(1385)^0$
1.1 ± 0.1	0.9 ± 0.2	0.3 ± 0.1	0.045 ± 0.015	1.6 ± 0.5	0.51 ± 0.32

BES3 : 10⁸ ~ 10⁹ $\psi(3686)$ & 10⁹ ~ 10¹⁰ $J/\psi(3097)$

N* ~ 2740 MeV, Λ^* ~ 2570 MeV,

Σ^* ~ 2490 MeV, Ξ^* ~ 2360 MeV

3. N^* from $pp \rightarrow N N^*$

complimentary e^+e^- and pp experiments

$$\psi \rightarrow P K^- \bar{\Lambda} \quad \text{vs} \quad P P \rightarrow P K^+ \Lambda$$

$P \bar{\Lambda}$ & $P \Lambda$ **the same t-channel interaction**

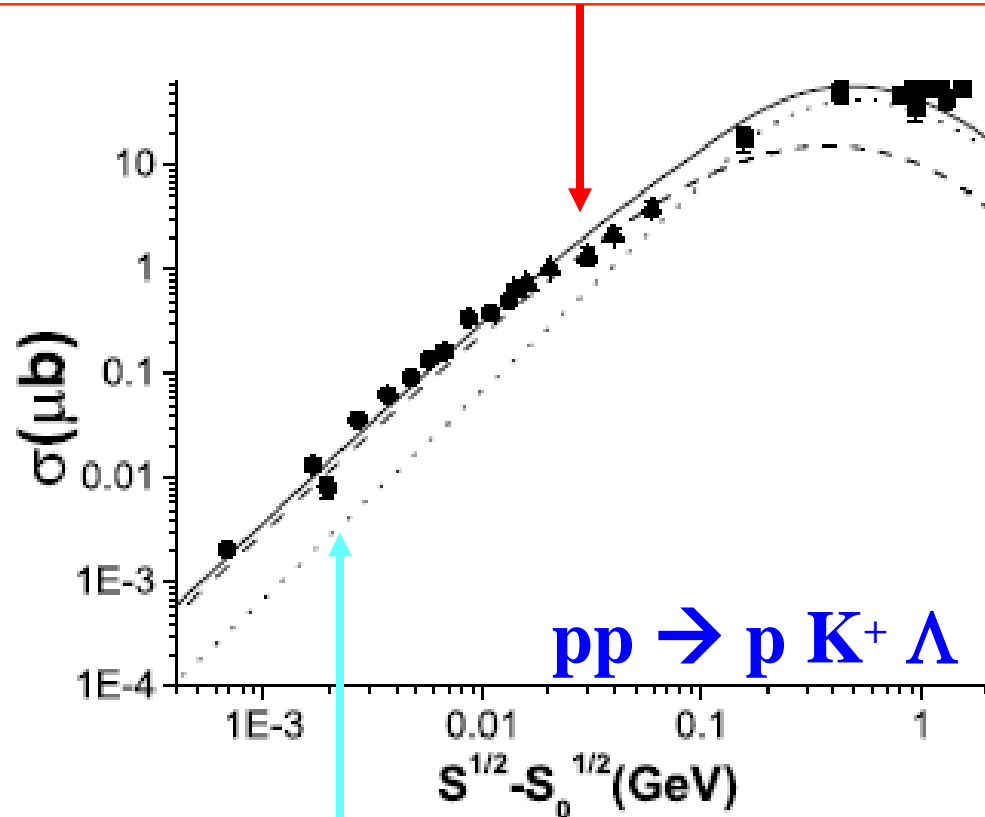
$K^- \bar{\Lambda}$ & $K^+ \Lambda$ **the same interaction**

$P K^-$ for Λ^* $P K^+$ for pentaquarks

Evidence for large $g_{N^*(1535)K\Lambda}$ from $pp \rightarrow p K^+ \Lambda$

**Total cross section and theoretical results with
 $N^*(1535)$, $N^*(1650)$, $N^*(1710)$, $N^*(1720)$**

B.C.Liu, B.S.Zou, Phys. Rev. Lett. 96 (2006) 042002



Tsushima, Sibirtsev, Thomas, PRC59 (1999) 369, without including $N^*(1535)$

FSI vs $N^*(1535)$ contribution in $pp \rightarrow p K^+ \Lambda$

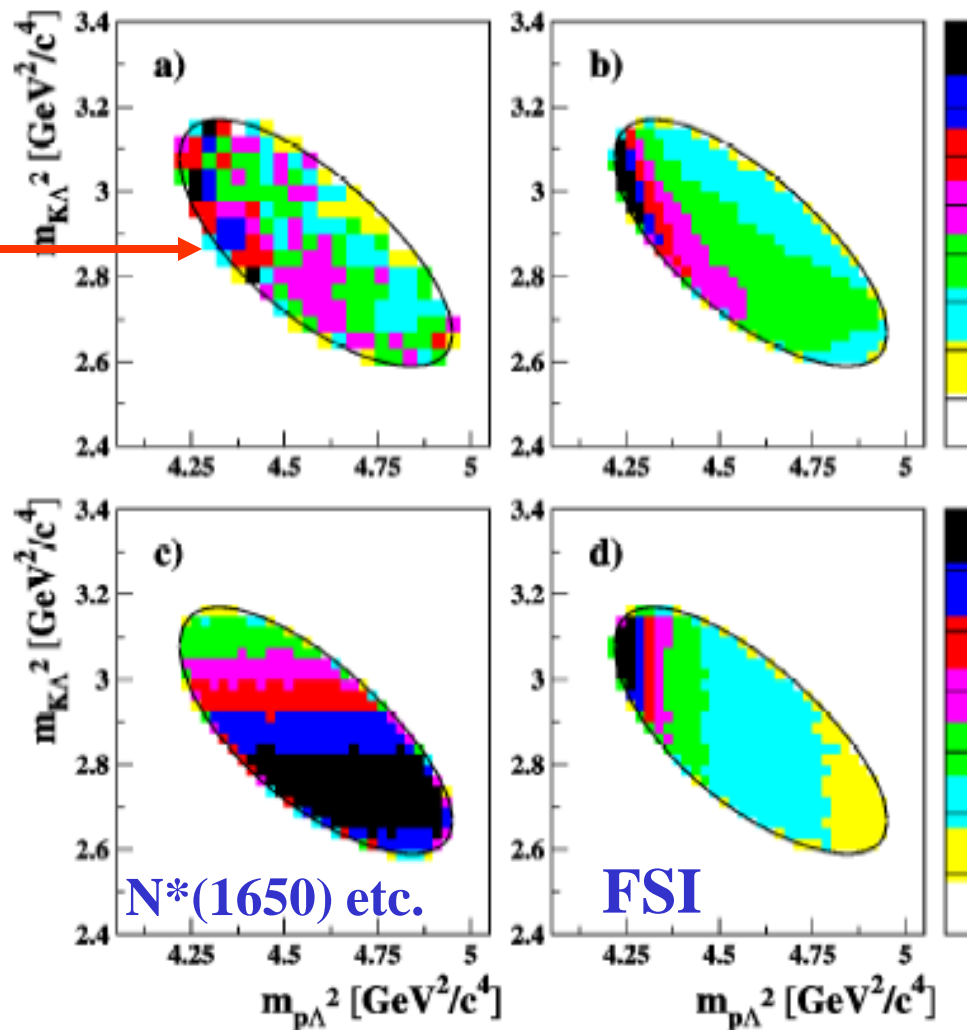
B.C.Liu & B.S.Zou, Phys. Rev. Lett. 98 (2007) 039102 (reply)

A.Sibirtsev et al., Phys. Rev. Lett. 98 (2007) 039101 (comment)

COSY-TOF data
S. Abdel-Samad *et al.*,
Phys.Lett.B632:27(2006)



**Both FSI & $N^*(1535)$
are needed !**



$N^*(1440)$ in $pp \rightarrow pn\pi^+$, consistent with $J/\psi \rightarrow \bar{p}n\pi^+$

CELSIUS-WASA Collaboration, nucl-ex/0612015

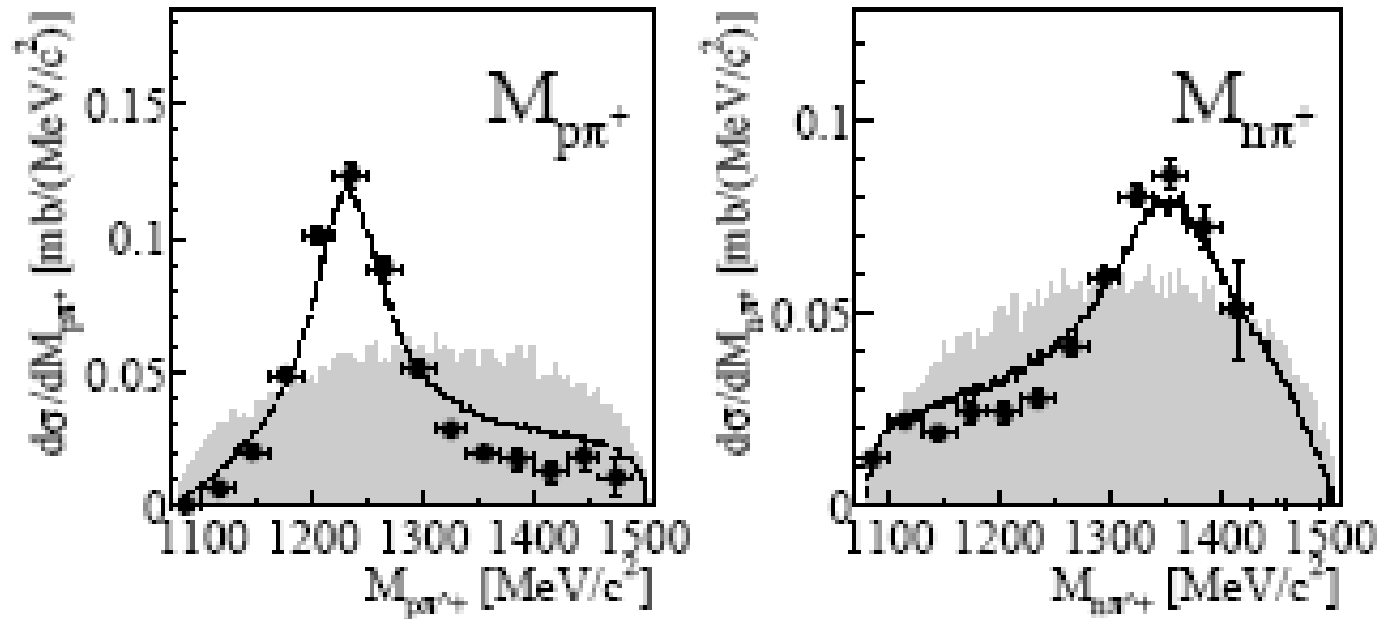
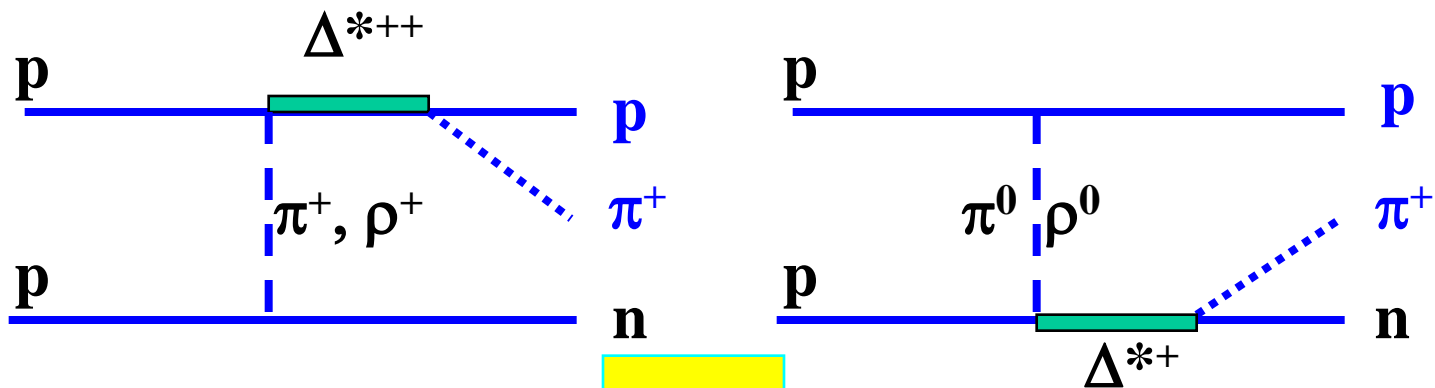
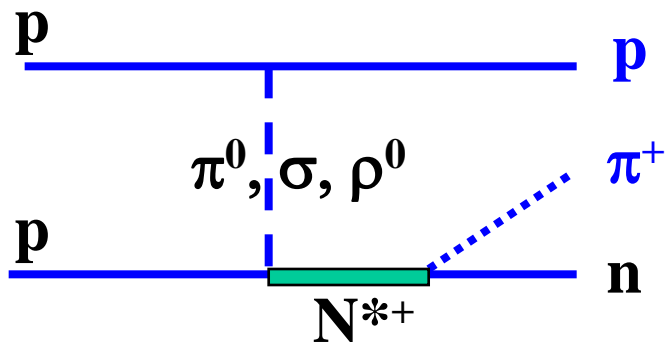


Figure 2: Invariant mass spectra $M_{pn\pi^+}$ and $M_{n\pi^+}$ obtained from the measurement of the $pp \rightarrow pn\pi^+$ reaction at $T_p = 1300$ MeV. The shaded areas show the pure phase space distributions, whereas the dashed lines show calculations assuming BW shapes for both the Δ^{++} and the Roper excitation [19].

Higher energies and Dalitz plots are needed at COSY&CSR !



9:1

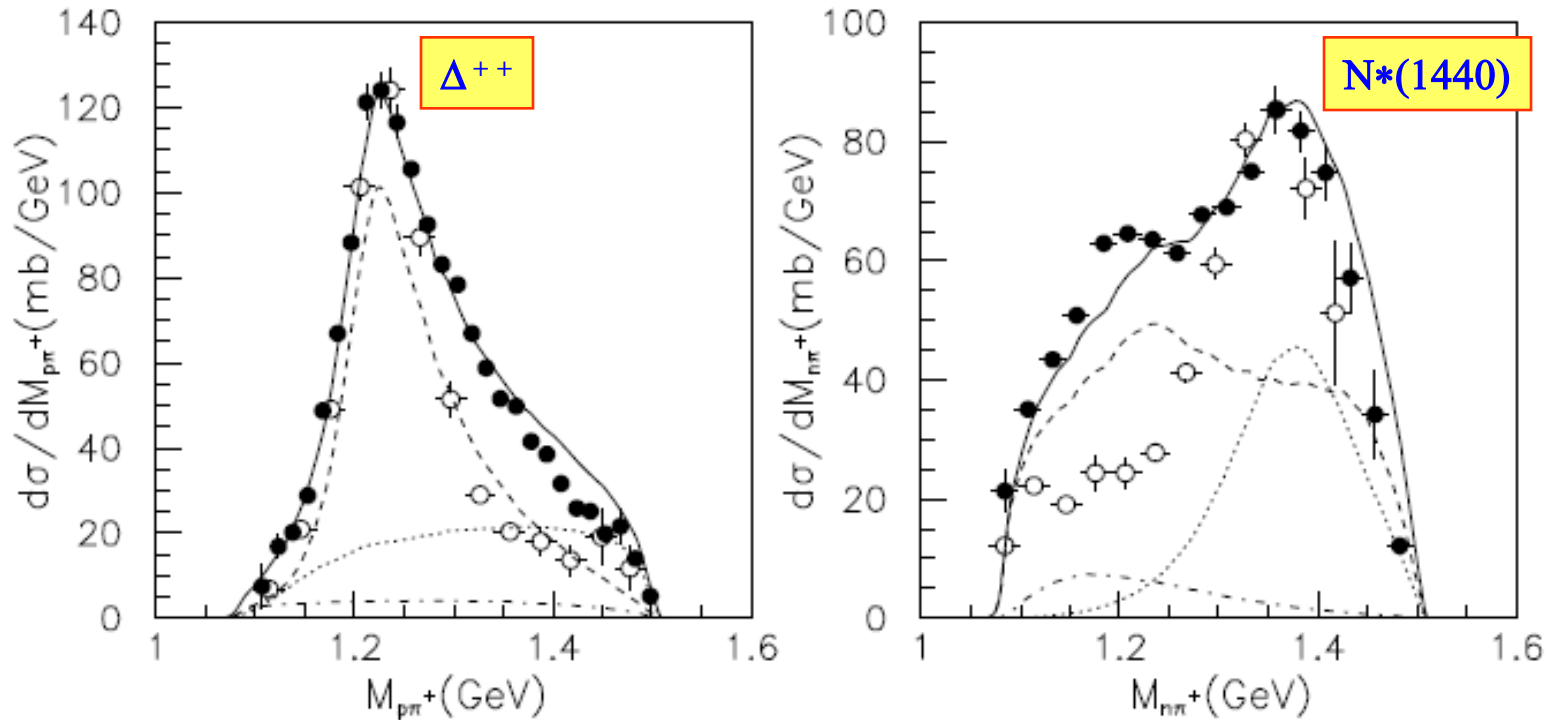


A quite good isospin filter !

Look for “missing” N^* by σp production !

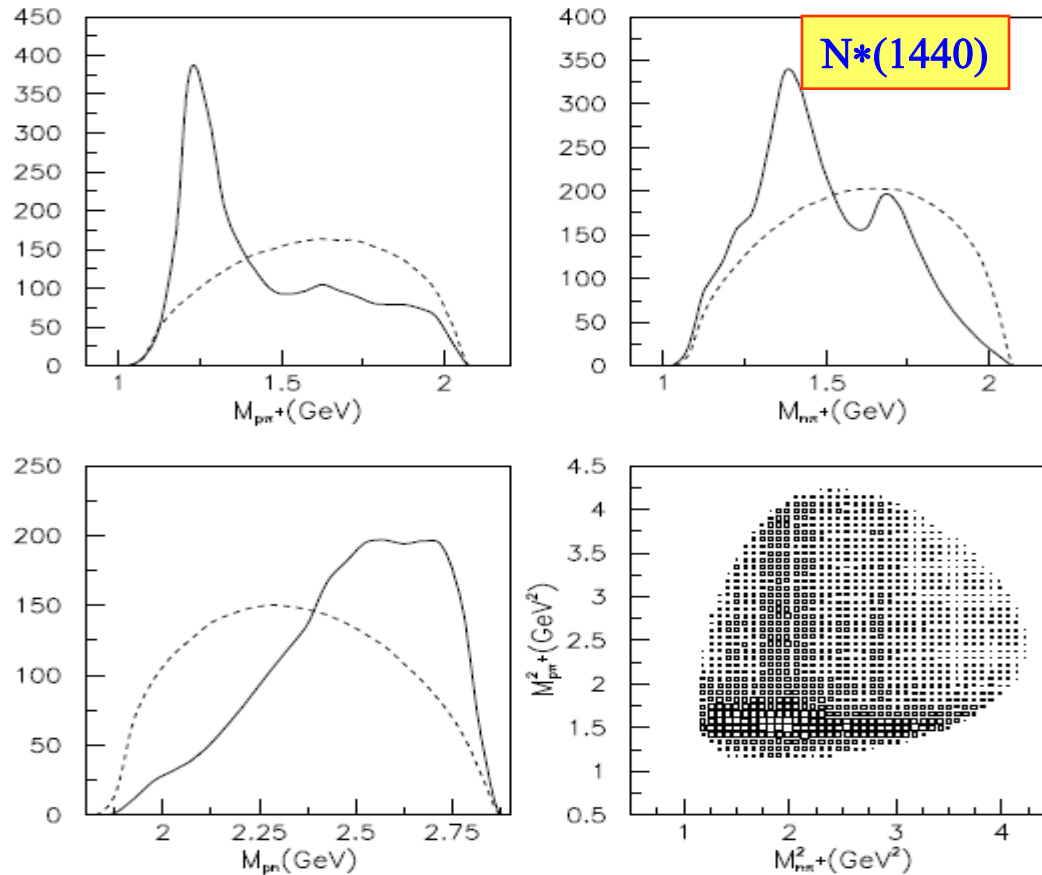
Theoretical study on $pp \rightarrow pn\pi^+$

Z.Ouyang, J.J.Xie, B.S.Zou & H.S.Xu, Nucl.Phys.A 821(2009)220; IJMPE(2009)



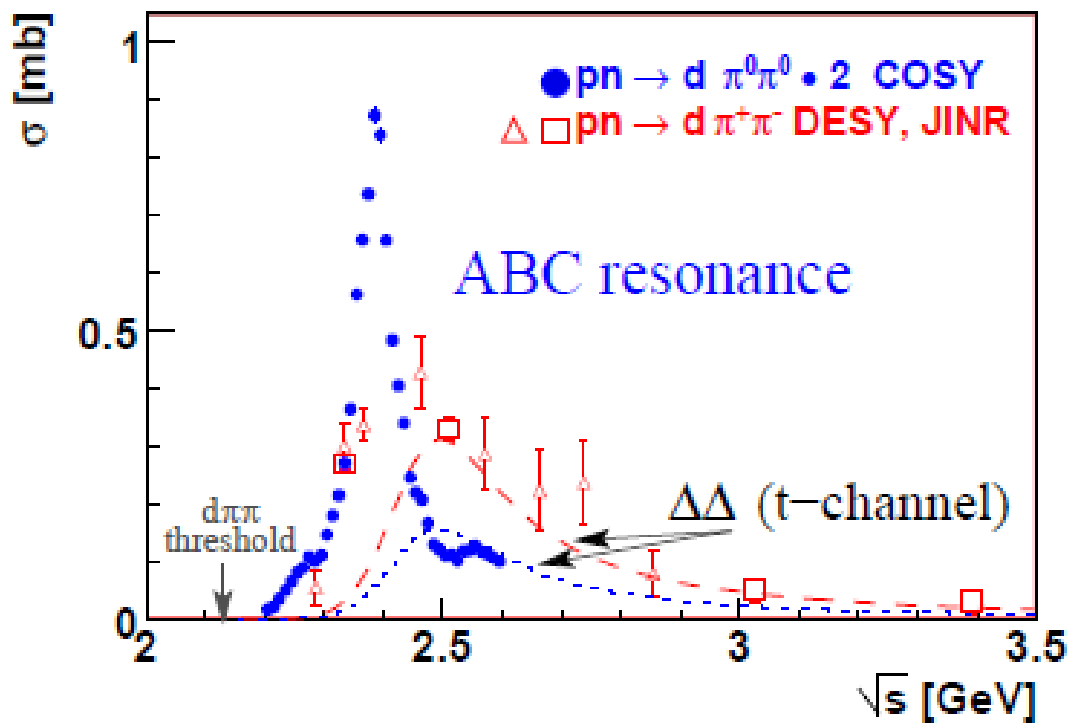
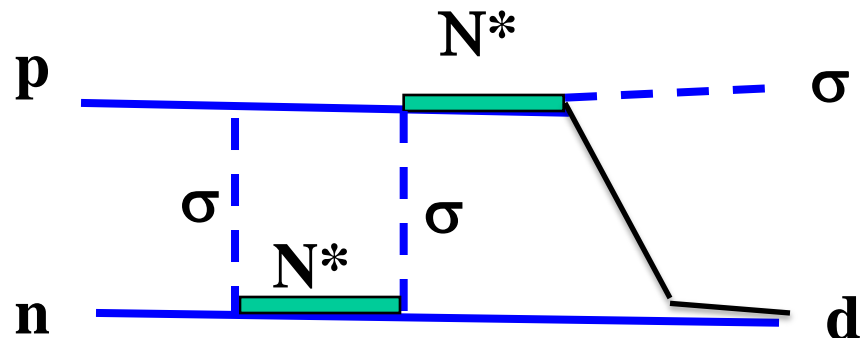
Including $N^*(1440)$ by σ exchange is crucial to reproduce data

Prediction for $pp \rightarrow pn\pi^+$ at $T_p = 2.88$ GeV



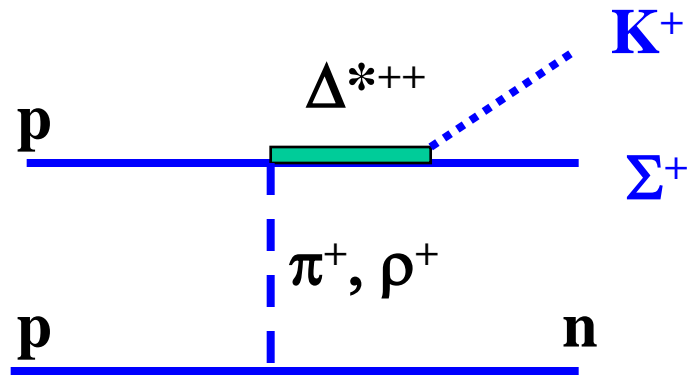
to be checked by COSY & CSR

a $NN^*(1440) \rightarrow d\sigma$ state ?



$pp \rightarrow nK^+\Sigma^+$: the best place for $\rho^+p \rightarrow \Delta^{*++} (uuu)$

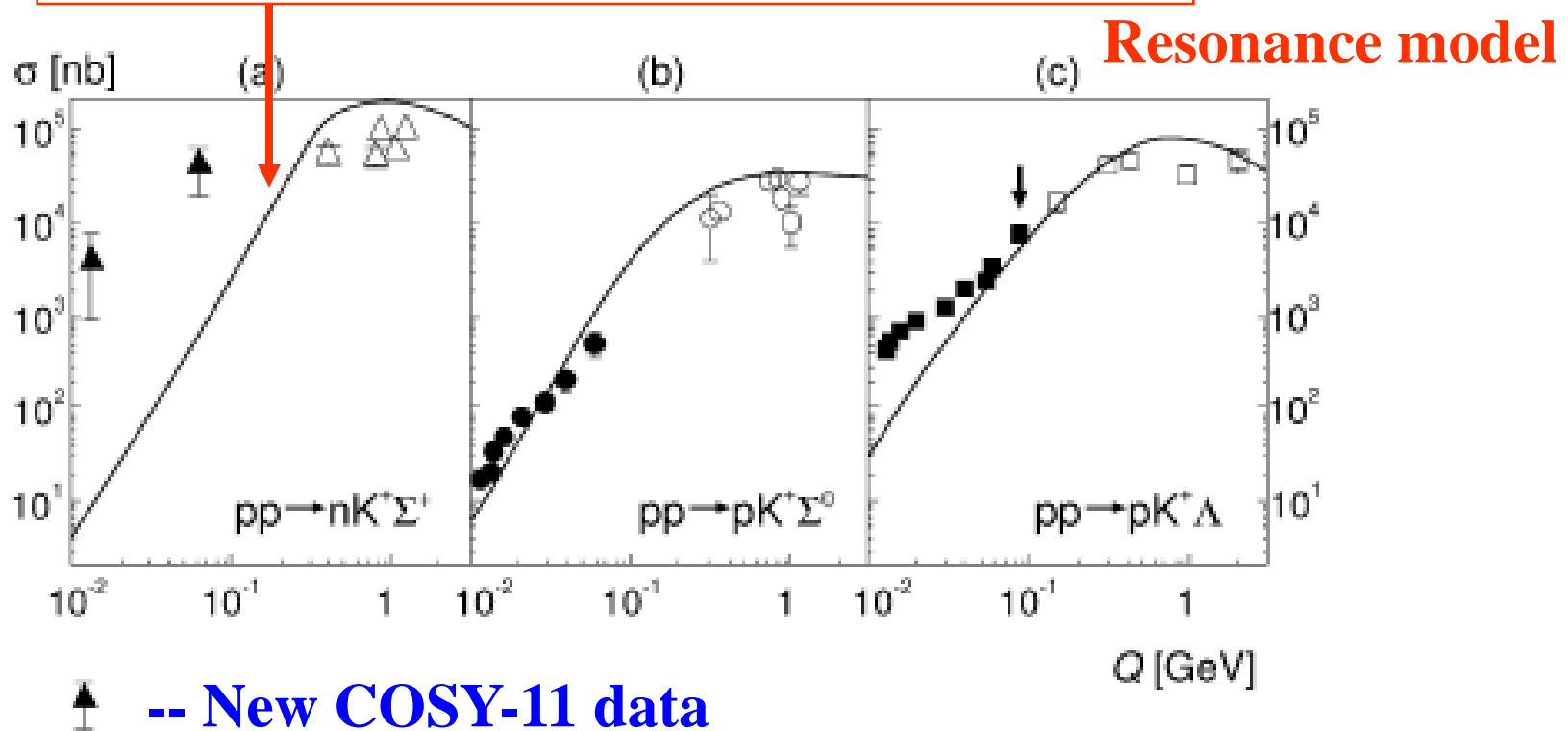
J.J.Xie & B.S.Zou, Phys. Lett. B649 (2007) 405



- $\Delta^{*++} (uuu)$ -- the most accessible system with 3 identical valence quarks
- our present knowledge from π^+p data is still very poor
- accessible at COSY and CSR

Status on the study of $pp \rightarrow nK^+\Sigma^+$

N. Tsushima, et al., Phys. Rev. C 59 (1999) 369.



T. Rożek et al. / Physics Letters B 643 (2006) 251–256

Incoherent π/K exchange model also fails to reproduce these data

A. Sibirtsev, et al., Nucl. Phys. A 646 (1999) 427.

Our calculation for $pp \rightarrow nK^+\Sigma^+$

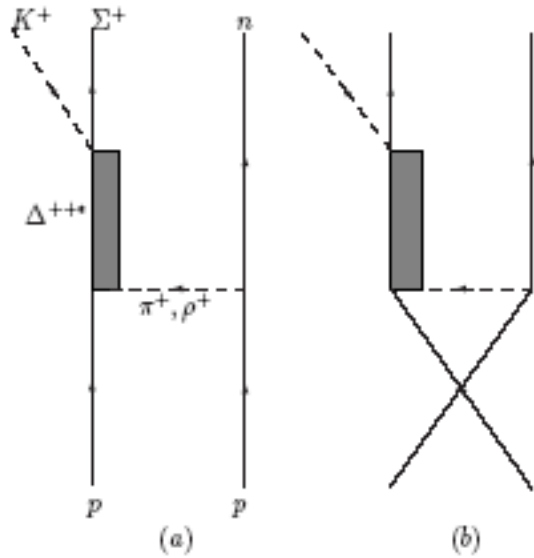
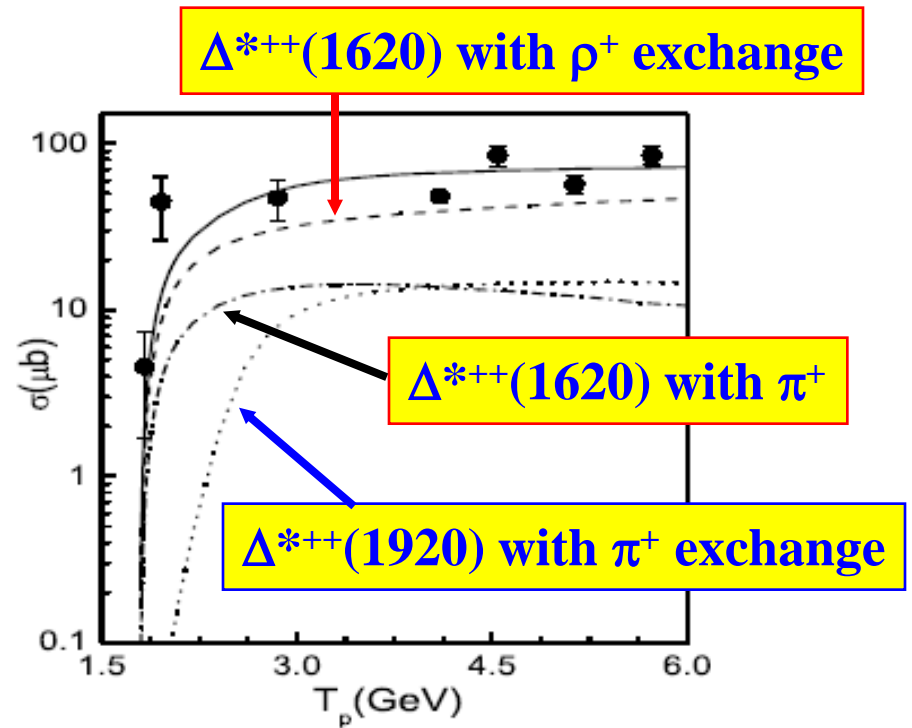


Figure 1: Feynman diagrams for $pp \rightarrow nK^+\Sigma^+$ reaction.



t-channel ρ -exchange plays important role !

Dalitz plot and projections: predictions to be checked by COSY&CSR

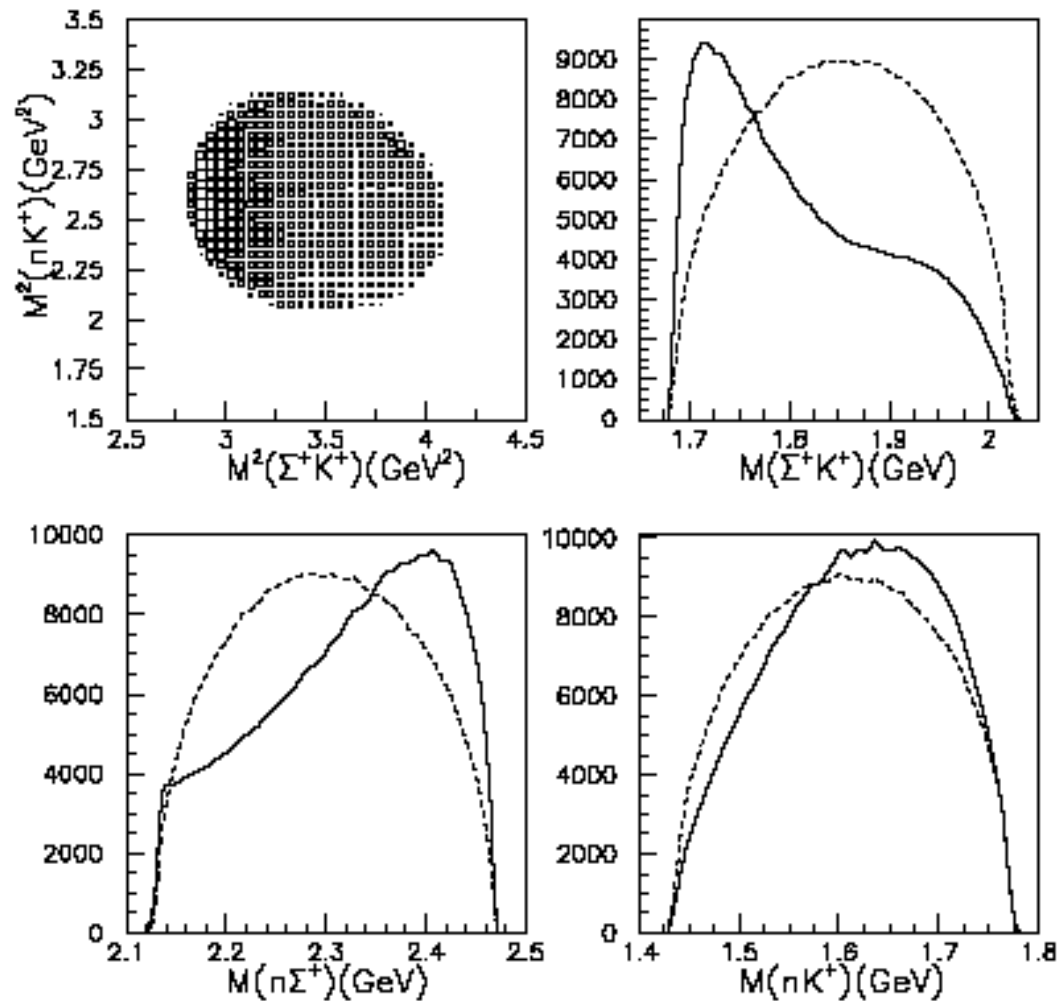


Figure 3: The Dalitz Plot and invariant mass spectra for the $pp \rightarrow nK^+\Sigma^+$ at $T_p=2.8\text{GeV}$, compared with pure phase space distributions (dashed curves)

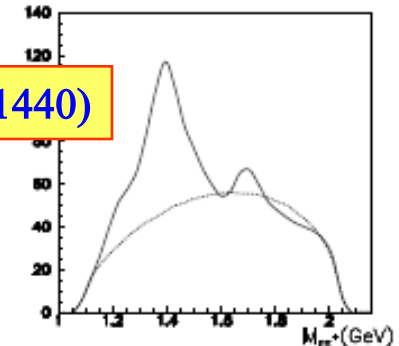
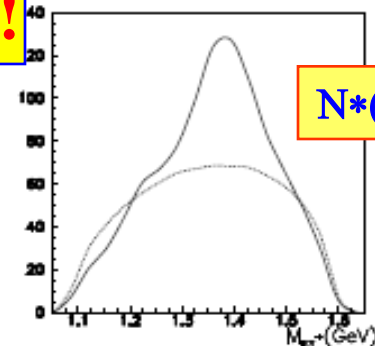
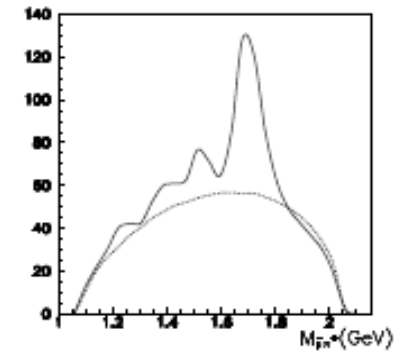
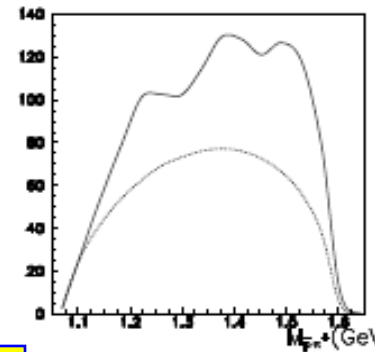
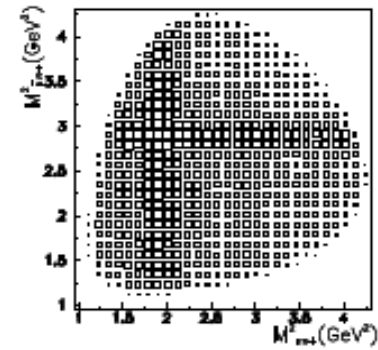
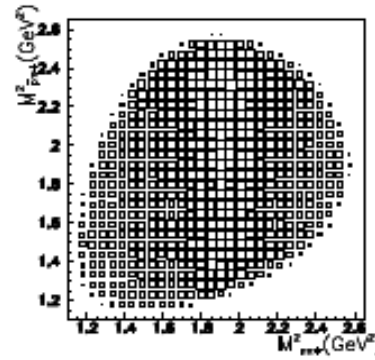
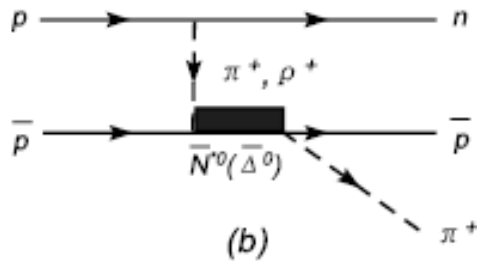
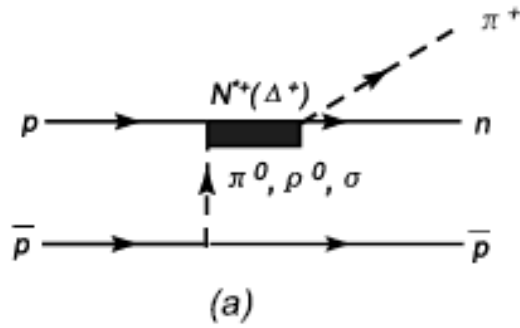
4. N^* from $\bar{p}p \rightarrow \bar{N} N^*$

$T_b \sim 2\text{GeV}$	$p p$	vs	$\bar{p} p$
σ_{tot} (mb)	50		90
σ_{el} (mb)	25		35
$\sigma_{pn\pi^+}$ (mb)	15		10

All final states of e^+e^- are accessible by $\bar{p}p$
A large portion of $\bar{p}p$ final states contain baryons
& should not be wasted at PANDA/FAIR

$\bar{p}p \rightarrow \bar{p}n\pi^+$: the best place for studying $N^*(1440)$ and other $N^* - N\sigma$

J.J.Wu, Z.Ouyang, B.S.Zou, ArXiv: 0902.2295 [hep-ph]



$T_b \sim 1.55\text{GeV}$

$T_b \sim 2.88\text{GeV}$

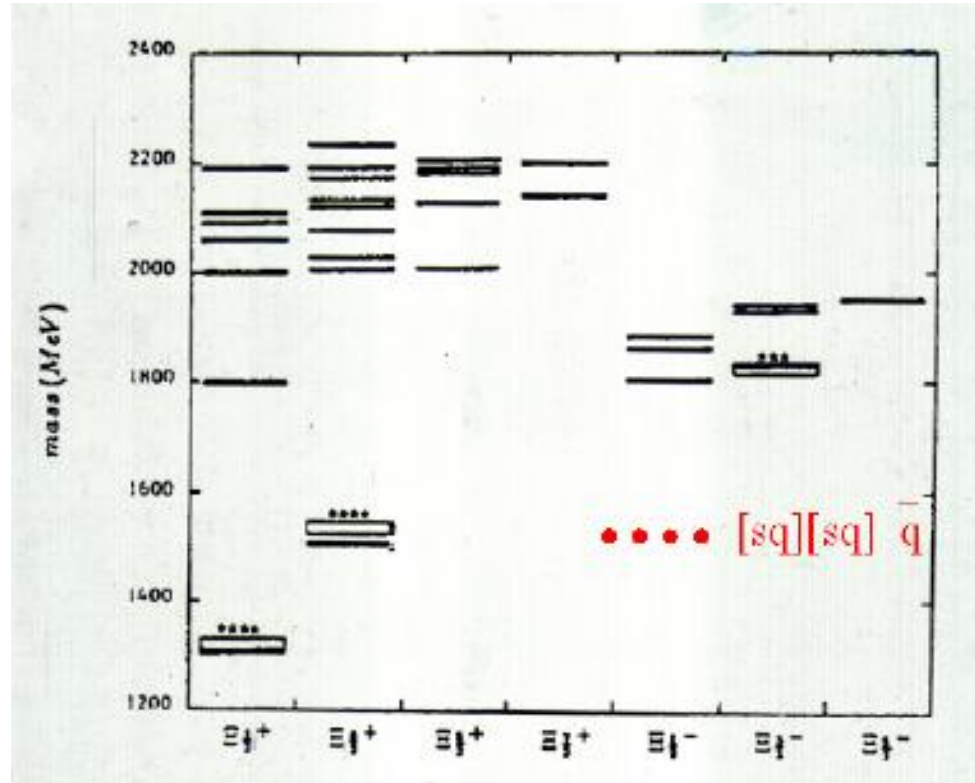
t-channel σ -exchange dominates !

Δ -production suppressed !

π, ρ -exchange $\rightarrow \bar{N}^* \sim 4N^*$

5. Summary and Prospects

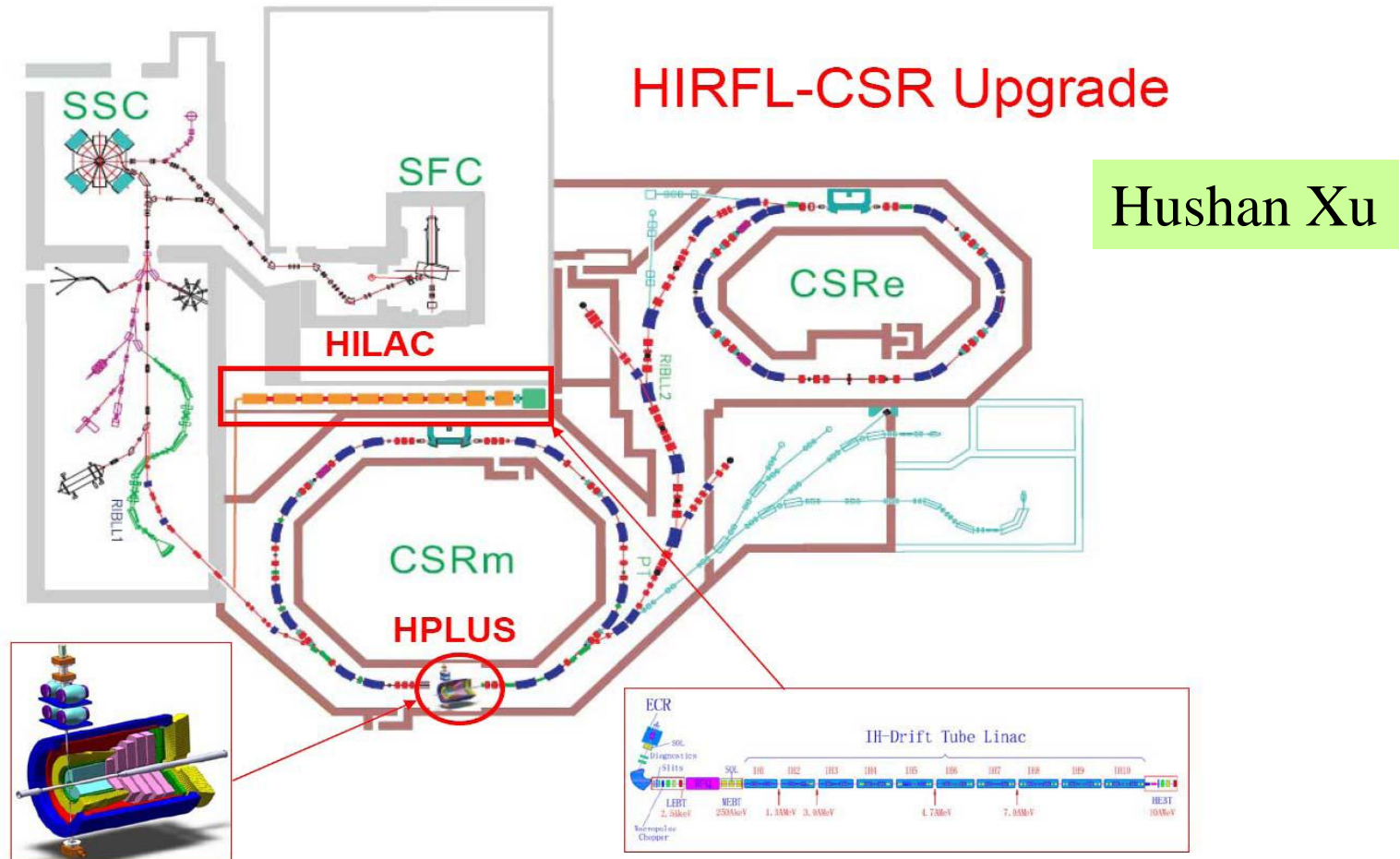
1) ψ decays : a nice filter for N^* , Δ^* , Λ^* , Σ^* , Ξ^*



BES3 : $10^8 \sim 10^9$ $\psi(3686)$ & $10^9 \sim 10^{10}$ $J/\psi(3097)$

Many new baryon resonances should be observed !

2) pp : nice place for $\rho^+p \rightarrow \Delta^{*++}$ & $\sigma p \rightarrow N^*$



Important for HPLUS/CSR to continue the effort at COSY

3) $\bar{p}p$: best place for $\sigma p \rightarrow N^*$ & maybe Ω^*

PANDA should also play important role on baryon spectroscopy.

+ γp at CEBAF, ELSA, Spring-8 + Kp at JPARC + ...



new era of baryon spectroscopy

Atomic Spectroscopy \Rightarrow Atomic Quantum Theory

**Nuclear Spectroscopy \Rightarrow Shell Model &
Collective motion Model**

Baryon Spectroscopy \Rightarrow ? Important discoveries

Thanks !