

Comparing AdS/CFT inspired dipole model to HERA F_2 data

Zhun Lu

Southeast University/Universidad Técnica Federico Santa María

The 5-th International Conference on Quarks and Nuclear Physics

Beijing, September 21-26, 2009

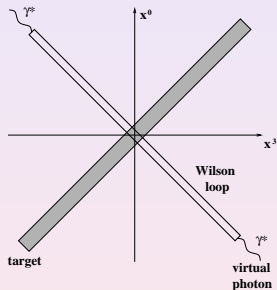
- 1 Brief introduction to AdS/CFT dipole model
- 2 Fitting AdS/CFT inspired dipole model to HERA F_2 data
- 3 Predictions
 - Inclusion of charm and prediction for F_2^c
 - prediction for F_L
 - prediction for total photoproduction cross section
- 4 Summary

Reference:

Yuri Kovchegov, ZL, Amir Rezaeian, arxiv:0906.4197

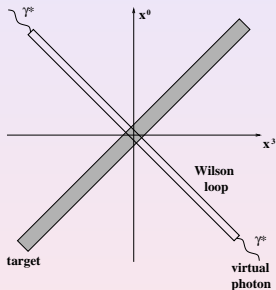
- Total deeply inelastic cross-section at large s

$$\sigma_{tot}^{\gamma^* A}(Q^2, x_{Bj}) \propto \int \frac{d^2 \mathbf{r} dz}{2\pi} \Psi(r, z; Q^2) N(\mathbf{r}, Y), \quad Y = \ln \left(\frac{1}{x_{Bj}} \right)$$



- Total deeply inelastic cross-section at large s

$$\sigma_{tot}^{\gamma^* A}(Q^2, x_{Bj}) \propto \int \frac{d^2 \mathbf{r} dz}{2\pi} \Psi(r, z; Q^2) N(\mathbf{r}, Y), \quad Y = \ln \left(\frac{1}{x_{Bj}} \right)$$



- N related to the expectation value of Wilson loop $\langle W \rangle$:

$$N(\mathbf{r}, Y) = 1 - S(\mathbf{r}, Y)$$

$$S(\mathbf{r}, \mathbf{b}, Y) = \frac{1}{N_c} \text{Re} \left\langle W \left(\mathbf{b} + \frac{1}{2} \mathbf{r}, \mathbf{b} - \frac{1}{2} \mathbf{r}, Y \right) \right\rangle$$

- An AdS/CFT based solution — $\mathcal{N} = 4$ Super Yang Mills $q\bar{q}$ dipole \iff open string in the AdS_5 space whose end points (q and \bar{q}) are located at the boundary of the AdS_5 space (J. Albacete et.al, [arxiv:0806.1484](#)) :

- An AdS/CFT based solution — $\mathcal{N} = 4$ Super Yang Mills $q\bar{q}$ dipole \iff open string in the AdS_5 space whose end points (q and \bar{q}) are located at the boundary of the AdS_5 space (J. Albacete et.al, [arxiv:0806.1484](#)) :

$$N(r, s) = 1 - S(r, s) = 1 - \exp \left[- \frac{\lambda_{YM}}{\sqrt{2}\pi s} \left(\frac{c_0^2 r^2}{\rho^3} + \frac{2}{\rho} - 2\sqrt{s} \right) \right]$$

- An AdS/CFT based solution — $\mathcal{N} = 4$ Super Yang Mills $q\bar{q}$ dipole \iff open string in the AdS_5 space whose end points (q and \bar{q}) are located at the boundary of the AdS_5 space (J. Albacete et.al, arxiv: 0806.1484) :

$$N(r, s) = 1 - S(r, s) = 1 - \exp \left[- \frac{\lambda_{YM}}{\sqrt{2}\pi s} \left(\frac{c_0^2 r^2}{\rho^3} + \frac{2}{\rho} - 2\sqrt{s} \right) \right]$$

- $\lambda_{YM} = g_{YM}^2 N_C$: 't Hooft coupling.

- An AdS/CFT based solution — $\mathcal{N} = 4$ Super Yang Mills $q\bar{q}$ dipole \iff open string in the AdS_5 space whose end points (q and \bar{q}) are located at the boundary of the AdS_5 space (J. Albacete et.al, [arxiv:0806.1484](#)) :

$$N(r, s) = 1 - S(r, s) = 1 - \exp \left[- \frac{\lambda_{YM}}{\sqrt{2}\pi s} \left(\frac{c_0^2 r^2}{\rho^3} + \frac{2}{\rho} - 2\sqrt{s} \right) \right]$$

- $\lambda_{YM} = g_{YM}^2 N_C$: 't Hooft coupling. In large N_C , λ_{YM} is large \Rightarrow strong coupling

- An AdS/CFT based solution — $\mathcal{N} = 4$ Super Yang Mills $q\bar{q}$ dipole \iff open string in the AdS_5 space whose end points (q and \bar{q}) are located at the boundary of the AdS_5 space (J. Albacete et.al, [arxiv:0806.1484](#)) :

$$N(r, s) = 1 - S(r, s) = 1 - \exp \left[- \frac{\lambda_{YM}}{\sqrt{2}\pi s} \left(\frac{c_0^2 r^2}{\rho^3} + \frac{2}{\rho} - 2\sqrt{s} \right) \right]$$

- $\lambda_{YM} = g_{YM}^2 N_C$: 't Hooft coupling. In large N_c , λ_{YM} is large \Rightarrow strong coupling
- g_{YM} : $\mathcal{N} = 4$ SYM coupling constant

- An AdS/CFT based solution — $\mathcal{N} = 4$ Super Yang Mills $q\bar{q}$ dipole \iff open string in the AdS_5 space whose end points (q and \bar{q}) are located at the boundary of the AdS_5 space (J. Albacete et.al, [arxiv: 0806.1484](#)) :

$$N(r, s) = 1 - S(r, s) = 1 - \exp \left[- \frac{\lambda_{YM}}{\sqrt{2}\pi s} \left(\frac{c_0^2 r^2}{\rho^3} + \frac{2}{\rho} - 2\sqrt{s} \right) \right]$$

- $\lambda_{YM} = g_{YM}^2 N_C$: 't Hooft coupling. In large N_C , λ_{YM} is large \Rightarrow strong coupling
- g_{YM} : $\mathcal{N} = 4$ SYM coupling constant
- $c_0 = \frac{\Gamma^2(\frac{1}{4})}{(2\pi)^{3/2}}$: a constant which relates the transverse dipole size r , the collision energy \sqrt{s} and the maximum extent of the string in the z -direction

- An AdS/CFT based solution — $\mathcal{N} = 4$ Super Yang Mills $q\bar{q}$ dipole \iff open string in the AdS_5 space whose end points (q and \bar{q}) are located at the boundary of the AdS_5 space (J. Albacete et.al, arxiv: 0806.1484) :

$$N(r, s) = 1 - S(r, s) = 1 - \exp \left[- \frac{\lambda_{YM}}{\sqrt{2}\pi s} \left(\frac{c_0^2 r^2}{\rho^3} + \frac{2}{\rho} - 2\sqrt{s} \right) \right]$$

- $\lambda_{YM} = g_{YM}^2 N_C$: 't Hooft coupling. In large N_C , λ_{YM} is large \Rightarrow strong coupling
- g_{YM} : $\mathcal{N} = 4$ SYM coupling constant
- $c_0 = \frac{\Gamma^2(\frac{1}{4})}{(2\pi)^{3/2}}$: a constant which relates the transverse dipole size r , the collision energy \sqrt{s} and the maximum extent of the string in the z-direction
- ρ is the function of r, s and c_0

s -dependent dipole form:

$$N(r, s) = 1 - \exp \left[- \frac{\lambda_{YM}}{\sqrt{2\pi s}} \left(\frac{c_0^2 r^2}{\rho^3} + \frac{2}{\rho} - 2\sqrt{s} \right) \right]$$

s -dependent dipole form:

$$N(r, s) = 1 - \exp \left[- \frac{\lambda_{YM}}{\sqrt{2\pi} s} \left(\frac{c_0^2 r^2}{\rho^3} + \frac{2}{\rho} - 2\sqrt{s} \right) \right]$$

Rewrite $N(r, s)$ to $N(r, x)$:

$$x = \frac{Q^2}{s + Q^2} \Rightarrow s = \frac{Q^2(1 - x)}{x}$$

$$Q \sim \frac{b_0}{r}$$

s -dependent dipole form:

$$N(r, s) = 1 - \exp \left[- \frac{\lambda_{YM}}{\sqrt{2\pi} s} \left(\frac{c_0^2 r^2}{\rho^3} + \frac{2}{\rho} - 2\sqrt{s} \right) \right]$$

Rewrite $N(r, s)$ to $N(r, x)$:

$$x = \frac{Q^2}{s + Q^2} \Rightarrow s = \frac{Q^2(1-x)}{x}$$

$$Q \sim \frac{b_0}{r}$$

$$s \approx \frac{b_0^2(1-x)}{r^2 x}$$

s-dependent dipole form:

$$N(r, s) = 1 - \exp \left[- \frac{\lambda_{YM}}{\sqrt{2}\pi s} \left(\frac{c_0^2 r^2}{\rho^3} + \frac{2}{\rho} - 2\sqrt{s} \right) \right]$$

Rewrite $N(r, s)$ to $N(r, x)$:

$$\begin{aligned} x &= \frac{Q^2}{s + Q^2} \Rightarrow s = \frac{Q^2(1-x)}{x} \\ Q &\sim \frac{b_0}{r} \\ s &\approx \frac{b_0^2(1-x)}{r^2 x} \end{aligned}$$

x-dependent dipole form:

$$\Rightarrow N(r, x) = 1 - \exp \left[- \frac{\lambda_{YM} x r}{\mathcal{M}_0^2(1-x)\pi\sqrt{2}} \left(\frac{1}{\rho_m^3} + \frac{2}{\rho_m} - 2\mathcal{M}_0 \sqrt{\frac{1-x}{x}} \right) \right]$$

$$\mathcal{M}_0 = b_0 c_0$$

x -dependent dipole form:

$$N(r, x) = 1 - \exp \left[- \frac{\lambda_{YM} x r}{\mathcal{M}_0^2 (1-x) \pi \sqrt{2}} \left(\frac{1}{\rho_m^3} + \frac{2}{\rho_m} - 2 \mathcal{M}_0 \sqrt{\frac{1-x}{x}} \right) \right]$$

x -dependent dipole form:

$$N(r, x) = 1 - \exp \left[- \frac{\lambda_{YM} x r}{\mathcal{M}_0^2 (1-x) \pi \sqrt{2}} \left(\frac{1}{\rho_m^3} + \frac{2}{\rho_m} - 2 \mathcal{M}_0 \sqrt{\frac{1-x}{x}} \right) \right]$$

$$\rho_m = \begin{cases} \left(\frac{1}{3m} \right)^{1/4} \sqrt{2 \cos(\frac{\theta}{3})} & : m \leq \frac{4}{27} \\ \sqrt{\frac{1}{3m\Delta} + \Delta} & : m > \frac{4}{27} \end{cases}$$

$$\theta = \arccos \left(\sqrt{\frac{27m}{4}} \right)$$

x -dependent dipole form:

$$N(r, x) = 1 - \exp \left[- \frac{\lambda_{YM} x r}{\mathcal{M}_0^2 (1-x) \pi \sqrt{2}} \left(\frac{1}{\rho_m^3} + \frac{2}{\rho_m} - 2\mathcal{M}_0 \sqrt{\frac{1-x}{x}} \right) \right]$$

$$\rho_m = \begin{cases} \left(\frac{1}{3m} \right)^{1/4} \sqrt{2 \cos(\frac{\theta}{3})} & : m \leq \frac{4}{27} \\ \sqrt{\frac{1}{3m\Delta} + \Delta} & : m > \frac{4}{27} \end{cases}$$

$$\theta = \arccos \left(\sqrt{\frac{27m}{4}} \right)$$

$$\Delta = \left[\frac{1}{2m} - \sqrt{\frac{1}{4m^2} - \frac{1}{27m^3}} \right]^{1/3}$$

$$m = \frac{\mathcal{M}_0^4 (1-x)^2}{x^2}$$

- Dipole cross-section:

$$\sigma_{q\bar{q}}(r, x) = \sigma_0 N(r, x)$$

- Parameters in AdS/CFT-based dipole model: $\sigma_0, M_0, \lambda_{YM}$

- Dipole cross-section:

$$\sigma_{q\bar{q}}(r, x) = \sigma_0 N(r, x)$$

- Parameters in AdS/CFT-based dipole model: $\sigma_0, M_0, \lambda_{YM}$

$$\sigma_{L,T}^{\gamma^*P}(Q^2, x) = \sum_f \int d^2r \int_0^1 dz |\Psi_{L,T}^{(f)}(r, z; Q^2)|^2 \sigma_{q\bar{q}}(r, x)$$

- Dipole cross-section:

$$\sigma_{q\bar{q}}(r, x) = \sigma_0 N(r, x)$$

- Parameters in AdS/CFT-based dipole model: $\sigma_0, M_0, \lambda_{YM}$

$$\sigma_{L,T}^{\gamma^*P}(Q^2, x) = \sum_f \int d^2r \int_0^1 dz |\Psi_{L,T}^{(f)}(r, z; Q^2)|^2 \sigma_{q\bar{q}}(r, x)$$

$$\begin{aligned} |\Psi_T^{(f)}(r, z; Q^2)|^2 &= \frac{\alpha_{EM} N_c}{2\pi^2} \sum_f e_f^2 \{ a_f^2 [K_1(r a_f)]^2 [z^2 + (1-z)^2] \\ &\quad + m_f^2 [K_0(r a_f)]^2 \} \\ |\Psi_L^{(f)}(r, z; Q^2)|^2 &= \frac{\alpha_{EM} N_c}{2\pi^2} \sum_f e_f^2 \{ 4 Q^2 z^2 (1-z)^2 [K_0(r a_f)]^2 \} \end{aligned}$$

- Dipole cross-section:

$$\sigma_{q\bar{q}}(r, x) = \sigma_0 N(r, x)$$

- Parameters in AdS/CFT-based dipole model: $\sigma_0, M_0, \lambda_{YM}$

$$\sigma_{L,T}^{\gamma^*P}(Q^2, x) = \sum_f \int d^2r \int_0^1 dz |\Psi_{L,T}^{(f)}(r, z; Q^2)|^2 \sigma_{q\bar{q}}(r, x)$$

$$\begin{aligned} |\Psi_T^{(f)}(r, z; Q^2)|^2 &= \frac{\alpha_{EM} N_c}{2\pi^2} \sum_f e_f^2 \{ a_f^2 [K_1(r a_f)]^2 [z^2 + (1-z)^2] \\ &+ m_f^2 [K_0(r a_f)]^2 \} \end{aligned}$$

$$|\Psi_L^{(f)}(r, z; Q^2)|^2 = \frac{\alpha_{EM} N_c}{2\pi^2} \sum_f e_f^2 \{ 4 Q^2 z^2 (1-z)^2 [K_0(r a_f)]^2 \}$$

$$F_2(Q^2, x) = \frac{Q^2}{4\pi^2 \alpha_{EM}} \left[\sigma_L^{\gamma^*P}(Q^2, x) + \sigma_T^{\gamma^*P}(Q^2, x) \right]$$

AdS/CFT-based fit to F_2 data

$$\sigma_{q\bar{q}}(r, x) = \sigma_0 \left(1 - \exp \left[- \frac{\lambda_{YM} x r}{\mathcal{M}_0^2 (1-x) \pi \sqrt{2}} \left(\frac{1}{\rho_m^3} + \frac{2}{\rho_m} - 2\mathcal{M}_0 \sqrt{\frac{1-x}{x}} \right) \right] \right)$$

AdS/CFT dipole model	λ_{YM}	$\mathcal{M}_0/10^{-3}$	$\sigma_0[\text{mb}]$	$\chi^2/\text{d.o.f.}$
$x \in [6.2 \times 10^{-7}, 10^{-4}]$	5	9.85	31.164	110.70/78 = 1.42
$x \in [6.2 \times 10^{-7}, 10^{-4}]$	20	6.36	22.65	141.12/78 = 1.81
$x \in [6.2 \times 10^{-7}, 6 \times 10^{-5}]$	5	10.114	30.97	44.24/60 = 0.74
$x \in [6.2 \times 10^{-7}, 6 \times 10^{-5}]$	10	8.16	26.08	49.22/60 = 0.82
$x \in [6.2 \times 10^{-7}, 6 \times 10^{-5}]$	20	6.54	22.47	55.195/60 = 0.92

Table: Parameters of the AdS/CFT dipole model determined from a fit to F_2 data reported by ZEUS in two Bjorken x bins. The value of quark mass $m_f = 140$ MeV ($f = u, d, s$) is taken for all fits. The data for the first two rows and the rest are within $Q^2/\text{GeV}^2 \in [0.045, 6.5]$ and $Q^2/\text{GeV}^2 \in [0.045, 2.5]$ respectively.

GBW dipole model fit to the same data bins

- Golec-Biernat-Wüsthoff (GBW) dipole model (Phys. Rev. D59 (1999) 014017; D60 (1999) 114023):

$$\sigma_{q\bar{q}}^{\text{GBW}}(r, x) = \sigma_0 \left(1 - e^{-r^2(\frac{x_0}{x})^\lambda/4}\right)$$

GBW dipole model	$x_0/10^{-4}$	λ	$\sigma_0[\text{mb}]$	$\chi^2/\text{d.o.f.}$
$x \in [6.2 \times 10^{-7}, 10^{-4}]$	2.225	0.299	22.77	63.09/78 = 0.81
$x \in [6.2 \times 10^{-7}, 6 \times 10^{-5}]$	2.371	0.368	21.13	39.35/60 = 0.66

Table: Parameters of the GBW color dipole model determined from a fit to F_2 data from ZEUS in two Bjorken x bins. The data for the first two rows and the rest are within $Q^2/\text{GeV}^2 \in [0.045, 6.5]$ and $Q^2/\text{GeV}^2 \in [0.045, 2.5]$

AdS/CFT-based fit to F_2 data

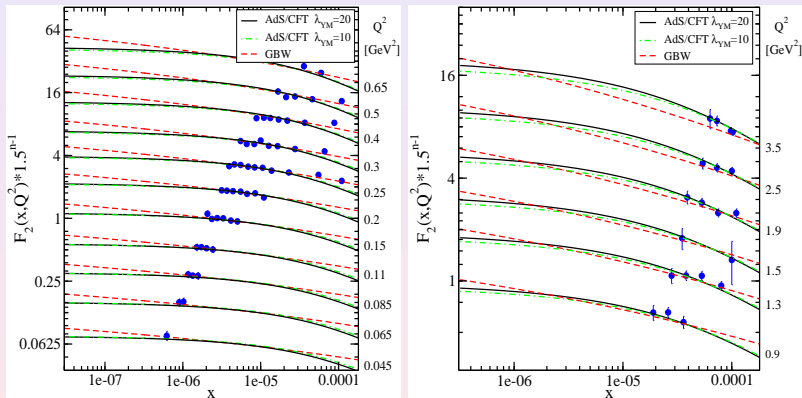
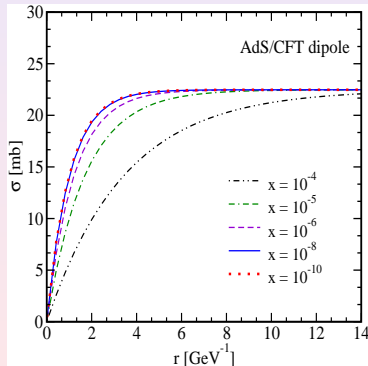


Figure: Results of our AdS/CFT-based fit to the proton structure function F_2 .

AdS/CFT-based dipole cross-section

$$\sigma_{q\bar{q}}(r, x) = \sigma_0 \left(1 - \exp \left[- \frac{\lambda_{YM} x r}{\mathcal{M}_0^2 (1-x) \pi \sqrt{2}} \left(\frac{1}{\rho_m^3} + \frac{2}{\rho_m} - 2\mathcal{M}_0 \sqrt{\frac{1-x}{x}} \right) \right] \right)$$



The AdS/CFT dipole cross-section obtained from the fit for $\lambda_{YM} = 20$ and $m_q = 140$ MeV at various fixed Bjorken- x as a function of r

Saturation scale

$$N(r_s = \sqrt{2}/Q_s, x) = \mathcal{N}_0 \equiv 1 - e^{-1/2} \approx 0.4$$

Saturation scale

$$N(r_s = \sqrt{2}/Q_s, x) = \mathcal{N}_0 \equiv 1 - e^{-1/2} \approx 0.4$$

$$Q_s^{\text{GBW}}(x) = \left(\frac{x_0}{x}\right)^{\lambda/2}$$

Saturation scale

$$N(r_s = \sqrt{2}/Q_s, x) = \mathcal{N}_0 \equiv 1 - e^{-1/2} \approx 0.4$$

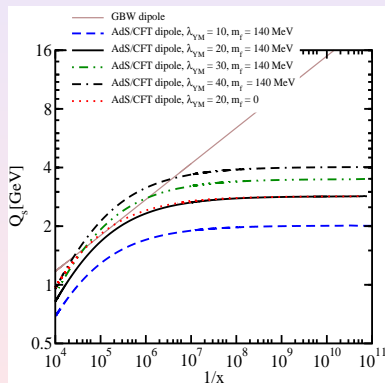
$$Q_s^{\text{GBW}}(x) = \left(\frac{x_0}{x}\right)^{\lambda/2} \quad Q_s^{\text{AdS}}(x) = \frac{2\lambda_{YM}x}{\mathcal{M}_0^2(1-x)\pi} \left(\frac{1}{\rho_m^3} + \frac{2}{\rho_m} - 2\mathcal{M}_0\sqrt{\frac{1-x}{x}} \right)$$

Saturation scale

$$N(r_s = \sqrt{2}/Q_s, x) = \mathcal{N}_0 \equiv 1 - e^{-1/2} \approx 0.4$$

$$Q_s^{\text{GBW}}(x) = \left(\frac{x_0}{x}\right)^{\lambda/2}$$

$$Q_s^{\text{AdS}}(x) = \frac{2\lambda_{\text{YM}} x}{\mathcal{M}_0^2(1-x)\pi} \left(\frac{1}{\rho_m^3} + \frac{2}{\rho_m} - 2\mathcal{M}_0 \sqrt{\frac{1-x}{x}} \right)$$

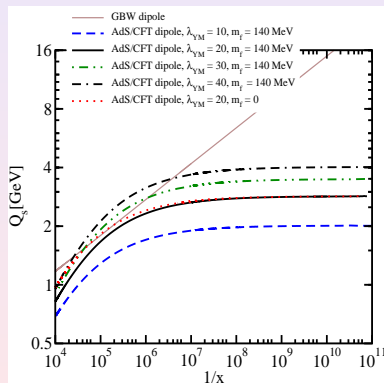


Saturation scale

$$N(r_s = \sqrt{2}/Q_s, x) = \mathcal{N}_0 \equiv 1 - e^{-1/2} \approx 0.4$$

$$Q_s^{\text{GBW}}(x) = \left(\frac{x_0}{x}\right)^{\lambda/2}$$

$$Q_s^{\text{AdS}}(x) = \frac{2\lambda_{\text{YM}} x}{\mathcal{M}_0^2(1-x)\pi} \left(\frac{1}{\rho_m^3} + \frac{2}{\rho_m} - 2\mathcal{M}_0 \sqrt{\frac{1-x}{x}} \right)$$



x -independence of $Q_s^{\text{AdS}}(x)$ at small x : saturation of saturation
 Conjectured in [arxiv:0707.0811](https://arxiv.org/abs/0707.0811),
 D. Kharzeev et.al.

s-dependent dipole fit

$$\sigma_{q\bar{q}}(r, s) = \sigma_0 \left(1 - \exp \left[- \frac{\lambda_{YM}}{s\pi\sqrt{2}} \left(\frac{c_0^2 r^2}{\rho^3} + \frac{2}{\rho} - 2\sqrt{s} \right) \right] \right)$$

λ_{YM}	c_0	$\sigma_0[\text{mb}]$	$\chi^2/\text{d.o.f.}$
5	0.00583	40.55	62.61/60 = 1.04
10	0.00440	36.30	77.17/60 = 1.29
20	0.00324	33.58	92.11/60 = 1.53

Table: Parameters of the s -dependent AdS/CFT dipole model determined from a fit to F_2 data from ZEUS. The data are within $x \in [6.2 \times 10^{-7}, 6 \times 10^{-5}]$ and $Q^2/\text{GeV}^2 \in [0.045, 2.5]$.

- In $\mathcal{N} = 4$ SYM: $c_0 = \frac{\Gamma^2(\frac{1}{4})}{(2\pi)^{3/2}} = 0.834$.
- our fitting result is two orders of magnitude smaller: difference between QCD and $\mathcal{N} = 4$ SYM.

- 1 Brief introduction to AdS/CFT dipole model
- 2 Fitting AdS/CFT inspired dipole model to HERA F_2 data
- 3 Predictions**
 - Inclusion of charm and prediction for F_2^c
 - prediction for F_L
 - prediction for total photoproduction cross section
- 4 Summary

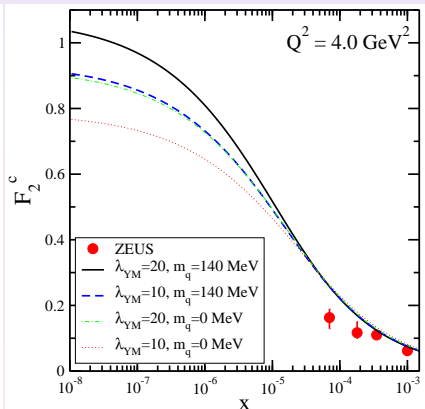
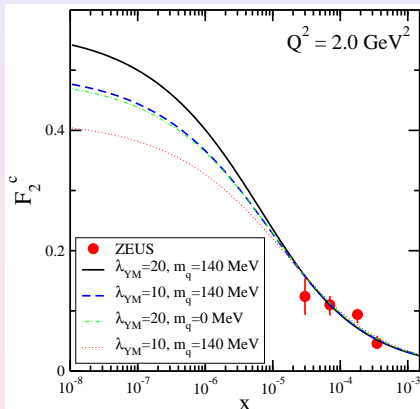
Inclusion of charm and prediction for F_2^c

$$F_2 = F_2^u + F_2^d + F_2^s + F_2^c$$

$m_c[\text{GeV}]$	λ_{YM}	$\mathcal{M}_0/10^{-3}$	$\sigma_0[\text{mb}]$	$\chi^2/\text{d.o.f.}$
-	10	8.16	26.08	49.22/60 = 0.82
-	20	6.54	22.47	55.195/60 = 0.92
1.4	10	7.66	24.72	61.66/60 = 1.03
1.4	20	6.16	21.31	70.99/60 = 1.18

Table: Parameters of the AdS/CFT dipole model determined from a fit to F_2 data reported by ZEUS. The value of quark mass $m_f = 140$ MeV is taken for all fits. The data used are within $x \in [6.2 \times 10^{-7}, 6 \times 10^{-5}]$ and $Q^2/\text{GeV}^2 \in [0.045, 2.5]$ respectively.

Inclusion of charm and prediction for F_2^c



prediction for F_L

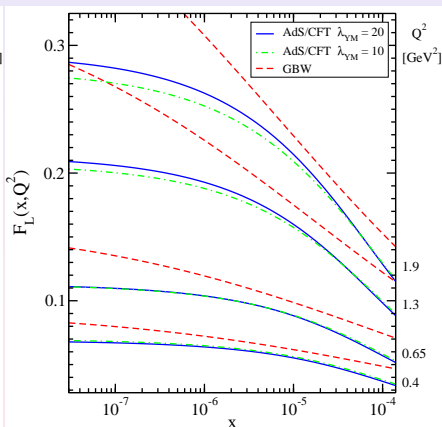
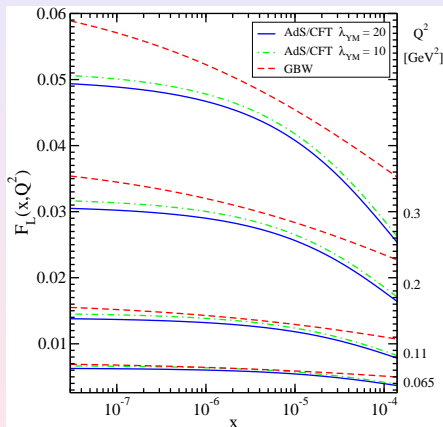


Figure: F_L predicted from our fit

photoproduction

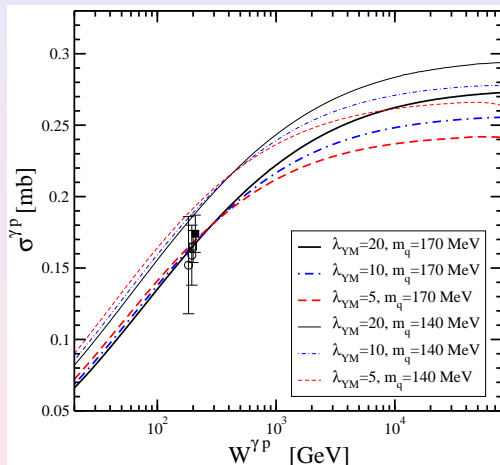


Figure: Photoproduction cross section at large \sqrt{s} predicted from our fit compared with ZEUS (square) and H1 (circle) data.

Summary

Summary

- HERA data for inclusive structure function $F_2(x, Q^2)$ at small Bjorken- x and Q^2 can be reasonably well described by an AdS/CFT-inspired dipole model with three parameters.

Summary

- HERA data for inclusive structure function $F_2(x, Q^2)$ at small Bjorken- x and Q^2 can be reasonably well described by an AdS/CFT-inspired dipole model with three parameters.
- The saturation scale in AdS/CFT-inspired dipole model varies in the range of $1 \div 3$ GeV becoming independent of energy (Bjorken- x) at very small x (10^{-8}), leads to the prediction of x -independent of the F_2 and F_L structure function at very small x .

Summary

- HERA data for inclusive structure function $F_2(x, Q^2)$ at small Bjorken- x and Q^2 can be reasonably well described by an AdS/CFT-inspired dipole model with three parameters.
- The saturation scale in AdS/CFT-inspired dipole model varies in the range of $1 \div 3$ GeV becoming independent of energy (Bjorken- x) at very small x (10^{-8}), leads to the prediction of x -independent of the F_2 and F_L structure function at very small x .
- Predictions for F_2 , F_2^c , F_L and photoproduction in the kinematical regions of future experiments.

Summary

- HERA data for inclusive structure function $F_2(x, Q^2)$ at small Bjorken- x and Q^2 can be reasonably well described by an AdS/CFT-inspired dipole model with three parameters.
- The saturation scale in AdS/CFT-inspired dipole model varies in the range of $1 \div 3$ GeV becoming independent of energy (Bjorken- x) at very small x (10^{-8}), leads to the prediction of x -independent of the F_2 and F_L structure function at very small x .
- Predictions for F_2 , F_2^c , F_L and photoproduction in the kinematical regions of future experiments.
- AdS/CFT-based model of non-perturbative physics at large coupling could be viewed as complimentary to the perturbative description of data based on saturation/Color Glass Condensate physics.

Backup slides

