



Bethe-Salpeter studies of mesons beyond rainbow-ladder

*5th International Conference on Quarks and Nuclear Physics,
21st–26th Sept., Beijing 2009*

Richard Williams and Christian Fischer

C. S. Fischer and RW, Phys. Rev. D **78** 2008 074006,
C. S. Fischer and RW, Phys. Rev. Lett. **103** 2009 122001,
C. S. Fischer and RW,

[arXiv:0808.3372]

[arXiv:0905.2291]

unpublished

24. September 2009



Motivation

Desire: description of hadronic bound-states in QCD

- Non-perturbative
- Formulated in the continuum
- Poincaré covariant
- Applicable to chiral/heavy quarks

Framework that satisfies all of these requirements:

- Bethe-Salpeter/Faddeev equations (mesons/baryons)
 - Input Green's functions from Dyson-Schwinger equations.

Description is

- *ab initio*, modulo necessary *truncations*. Work in Landau gauge.
- complementary to lattice QCD, effective theories, quark models etc.

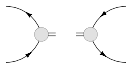
- Mesons: described by Bethe-Salpeter equation
- Baryons: quark-diquark model/Faddeev equation.

[G. Eichmann, PhD thesis (KFU Graz) [arXiv:0909.0703]]

Most common truncation: Rainbow-Ladder – *Time to do better?*

[P. Maris, P. C. Tandy, PRC **60** (1999) 055214]

Introduction: Bound-state amplitudes



$$(2\pi)^4 (P^2 + M^2)$$

- Meson bound-states: poles in two-body propagator.
- Depend on two momenta: total P and relative q .
- Amplitude defined on-shell $P^2 = -M^2$ (Euclidean Space).



$$= \Gamma^{(\mu)}(q, P) = \sum_i \tau_i(q, P) T_i^{(\mu)}(q, P)$$

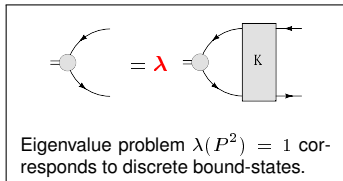
Relativistic covariants

[C. H. Llewellyn-Smith, Annals Phys. **53** (1969) 521.]

Component		0 ⁻⁺	0 ⁺⁺	Component	1 ⁻⁻	1 ⁺⁺	1 ⁺⁻
T_1	$\mathbb{1}$	○	○	T_1^μ	$i\gamma_T^\mu$	○	
T_2	$-i\not{P}$	○	●	T_2^μ	$\gamma_T^\mu \not{P}$	○	●
T_3	$-i\not{p}$	●	○	T_3^μ	$-\gamma_T^\mu \not{p} + p_T^\mu \mathbb{1}$	●	○
T_4	$[\not{p}, \not{P}]$	○	○	T_4^μ	$i\gamma_T^\mu [\not{P}, \not{p}] + 2ip_T^\mu \not{P}$	○	○
				T_5^μ	$p_T^\mu \mathbb{1}$	○	○
				T_6^μ	$ip_T^\mu \not{P}$	●	○
				T_7^μ	$-ip_T^\mu \not{p}$	○	●
				T_8^μ	$p_T^\mu [\not{P}, \not{p}]$	○	○

$J^{PC} = 0^{-+}, 1^{++}, 1^{+-}$:
prefixed by γ_5 .

Introduction: Bound-state amplitudes

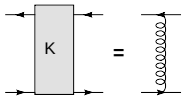


Solve homogeneous Bethe-Salpeter Equation:

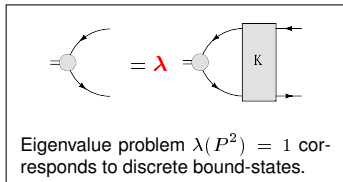
- Obtain Bound-state mass
- Electroweak decay constants
- Electromagnetic form factors
 - pion charge radius
 - magnetic moments
- transition form factors (hadronic decays)
- structure functions

Need:

- quark propagator from Dyson-Schwinger equation
- Normalisation (provided by an auxiliary condition)
- Dynamical information – Bethe-Salpeter kernel K



Introduction: Bound-state amplitudes

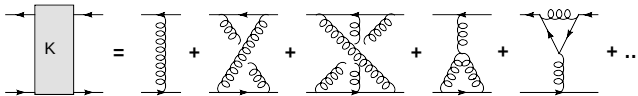


Solve homogeneous Bethe-Salpeter Equation:

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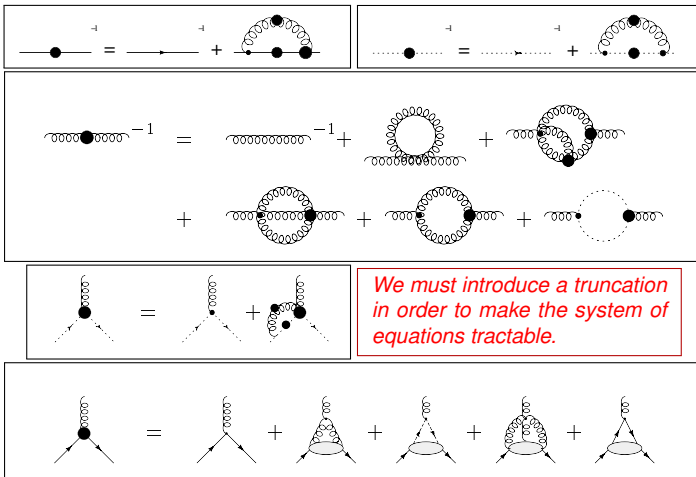
- quark propagator from Dyson-Schwinger equation
- Normalisation (provided by an auxiliary condition)
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Introduction: Dyson-Schwinger equations

infinite tower of coupled integral equations relating the Greens functions of QCD

- Propagators, Vertices and higher n -point Greens functions



Truncation: Minimum requirements

$$-i \not{p} A(p^2) + B(p^2) =$$



$$\text{gluon loop} = \left(\delta^{\mu\nu} - \frac{q^\mu q^\nu}{q^2} \right) \frac{Z(q^2)}{q^2}$$

$$= \sum_{i=1}^{12} \lambda_i(k, p; q) L_i^\mu(k, p; q)$$

- Quark propagator essential for the Bethe-Salpeter equation
- Specification of:
 - **Gluon propagator** and the **Quark-gluon vertex**
- sufficient to close the system and provide input for BSE.

That means provision of:

- One scalar dressing for the gluon
- Twelve scalar dressings for the quark, dependent upon the incoming and outgoing quark momenta.

Simplest truncation: the Rainbow

$$-i \not{p} A(p^2) + B(p^2) =$$



$$\text{gluon loop on quark line} = \text{gluon line} \frac{Z(q^2)}{q^2}$$

$$\text{gluon vertex} = \text{tree-level vertex} \lambda_1 ((k-p)^2)$$

Simplifications:

- Construct a reduced quark-gluon vertex:
 - Consider only the tree-level component, $\Gamma(k, p)^\mu = \lambda_1(k, p) \gamma^\mu$
 - $\lambda_1(k, p) \rightarrow \lambda_1(k - p)$ (important for BS equation)
- Provide a gluon dressing by
 - either **solve** Dyson-Schwinger equation for gluon
 - provide **Ansatz** that provides dynamical chiral symmetry breaking (IR strength, matching of perturbation theory etc.)

Simplest truncation: the Ladder

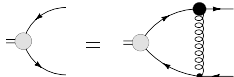
Axial-vector Ward–Takahashi identity

$$\begin{aligned} P_\mu \Gamma_{5\mu}^a(k; P) &= S^{-1}(k_+) \frac{1}{2} \lambda_f^a i \gamma_5 + \frac{1}{2} \lambda_f^a i \gamma_5 S^{-1}(k_-) \\ &\quad - M_\zeta i \Gamma_5^a(k; P) - i \Gamma_5^a(k; P) M_\zeta \end{aligned}$$

Satisfaction of this gives rise to

- Gell-Mann–Oakes–Renner
- Massless pion in the chiral limit

$$m_\pi^2 f_\pi \simeq 2m_q \langle \bar{q}q \rangle$$



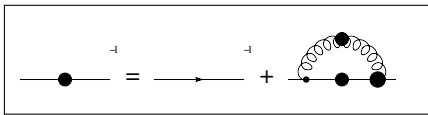
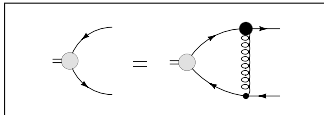
Symmetry preserving truncation

- Replace kernel by single gluon-exchange
- Quark-gluon vertex minimally dressed:

Require consistency between quark DSE and BS kernel

- Gluon dressing function $Z(q^2)$ must match
- Vertex dressing $\lambda_1(q)$ must match

Philosophy of Rainbow-Ladder



- Gluon dressing function $Z(q^2)$
- Quark-gluon vertex $\Gamma^\mu = \lambda_1(q^2) \gamma^\mu$

Phenomenological Model

- Provide model interaction
 $\alpha_{eff} = \alpha Z(q^2) \lambda_1(q^2)$
- Dynamical chiral symmetry breaking
- Chiral condensate
- Meson observables

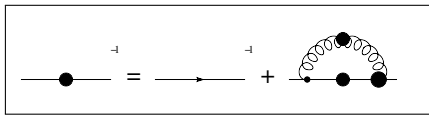
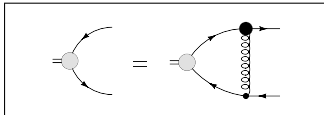
Entirely an Ansatz.

Towards *ab initio*

- Solve Gluon DSE
- Solve Quark-Gluon DSE
- Discard tensor structure
- restrict momentum dependence.

Rainbow-ladder but
Guided by calculations.

Philosophy of Rainbow-Ladder



Phenomenological models successful

Interaction features enhancement in infrared

→ Dynamical chiral symmetry breaking

→ Reasonable description of

- light-meson masses
- leptonic decay constants
- hadronic decays
- electromagnetic form factors

Result? Majority of properties are governed by chiral symmetry breaking, not the details of the interaction.

Limitations? Intrinsic to rainbow-ladder: restriction to vector-vector coupling limits important differences in different meson channels.

Interaction has no quark mass dependence (unless we model it)

Beyond Rainbow-Ladder

Any scheme must

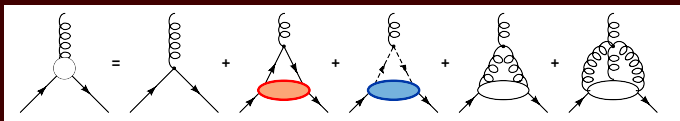
- Respect the axial-vector Ward-Takahashi identity
- Intimate relation between BS kernel and quark DS kernel

[H. J. Munczek, Phys. Rev. D **52** (1995) 4736.], [A. Bender, C. D. Roberts, L. von Smekal, Phys. Rev. Lett. B **380** (1996) 7]

Symmetry preserving truncation: obtained by cutting internal quark-lines in the quark Dyson-Schwinger equation

Look at the quark-gluon vertex

Dyson-Schwinger equation involves several unknown Green's functions



- Perform dressed skeleton expansion to perform truncation
- IR power counting, $1/N_c$ expansion yields dominant contributions.

Beyond Rainbow-Ladder: what has been done?

Any scheme must

- Respect the axial-vector Ward-Takahashi identity
- Intimate relation between BS kernel and quark DS kernel

[H. J. Munczek, Phys. Rev. D **52** (1995) 4736.], [A. Bender, C. D. Roberts, L. von Smekal, Phys. Rev. Lett. B **380** (1996) 7]

Symmetry preserving truncation may be obtained by cutting internal quark-lines in the quark Dyson-Schwinger equation

Semi-perturbative expansion (one-loop)

The diagram shows the semi-perturbative expansion of the quark Dyson-Schwinger equation at one-loop order. On the left is the full quark propagator, represented by a solid line with a black dot. This is equal to the sum of several terms: 1) the tree-level propagator (a simple solid line), 2) a term with coefficient $\frac{N_c}{2}$ representing a one-loop correction with a gluon loop (two wavy lines forming a triangle), and 3) a term with coefficient $-\frac{1}{2N_c}$ representing a one-loop correction with a ghost loop (two dashed lines forming a triangle). The equation ends with $+\dots$ indicating higher-order terms.

$$\text{[Diagram: Full propagator]} = \text{[Diagram: Tree-level propagator]} + \frac{N_c}{2} \text{[Diagram: Gluon loop]} - \frac{1}{2N_c} \text{[Diagram: Ghost loop]} + \dots$$

Beyond Rainbow-Ladder: what has been done?

Any scheme must

- Respect the axial-vector Ward-Takahashi identity

→ Intimate relation between BS kernel and quark DS kernel

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Symmetry preserving truncation may be obtained by cutting internal quark-lines in the quark Dyson-Schwinger equation

Semi-perturbative expansion (one-loop)

$$\text{Vertex} = \text{Vertex} + \frac{N_c}{2} \text{Triangle (gluons)} - \frac{1}{2N_c} \text{Triangle (quarks)} + \dots$$

Trivialization of the gluon momentum dependence (reduces # integrals):

$$D_{\mu\nu}(q) \propto \left(\delta_{\mu\nu} - \frac{q_\mu q_\nu}{q^2} \right) \delta^{(4)}(q) \quad \sim (\text{Munczek-Nemirovsky})$$

[A. Bender, C. D. Roberts, L. von Smekal, Phys. Rev. Lett. B **380** (1996) 7]

[P. Watson, W. Cassing and P. C. Tandy, Few Body Syst. **35** (2004) 129]

[H. H. Matevosyan, A. W. Thomas and P. C. Tandy, Phys. Rev. C **75** (2007) 045201]

Beyond Rainbow-Ladder: what **must** be done?

To investigate dominant contributions we need:

- a gluon with non-trivial momentum dependence
- appropriate internal vertex dressings

We must bite the bullet and tackle two-loop integrals over non-perturbative propagators and vertex dressings. *Challenging on a technical level.*

First study: *compare to existing literature*

Investigate the new truncation with a **model gluon** that provide $D\chi SB$:

- no trivialization of the gluon momentum

→ study qualitative impact of vertex corrections

To perform a quantitative investigation: replace the gluon Ansatz with something more appropriate, and specify the internal vertex dressings.

$$\alpha\lambda_1(q^2)Z(q^2) = \frac{\pi D}{\omega^2} q^4 e^{-q^2/\omega^2}$$

- Enhanced at intermediate momentum
- Damped in UV (no renormalisation)

Beyond Rainbow-Ladder: what **must** be done?

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We must bite the bullet and tackle two-loop integrals over non-perturbative propagators and vertex dressings. *Challenging on a technical level.*

What do we gain?

- Truncation does not restrict what mesons we can study:

$$(J^{PC}) \quad \frac{\pi}{0^{-+}} \quad \frac{\sigma}{0^{++}} \quad \frac{\rho}{1^{--}} \quad \frac{a_1}{1^{++}} \quad \frac{b_1}{1^{+-}}$$

- **Quark-gluon vertex:**

- all twelve components – dynamically generate scalar parts
- full momentum dependence

→ imparts quark-mass dependence on interaction.

- **Different interactions:**

- vector $\times$ vector
- vector $\times$ scalar
- different degrees of attraction/repulsion in different meson channels.

Beyond Rainbow-Ladder: what **must** be done?

To investigate dominant contributions we need:

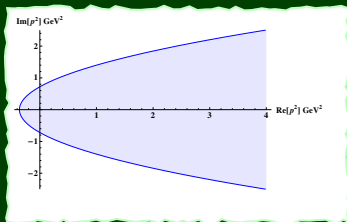
- a gluon with non-trivial momentum dependence
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Technical challenges?

Auxiliary normalisation condition: (*Leon–Cutkosky*)

- momentum dependent kernel complicates evaluation



Bound state in Eucl. space, $P^0 = iM$
quark and vertex momenta $(q \pm P/2)^2$:

- must be evaluated for complex momenta
- region bounded by parabola dependent upon meson mass.

$$\delta^{ij} = \frac{\partial}{\partial P^2} \text{tr} \int_k \left[\left(\bar{\Gamma}_\pi^i S \Gamma_\pi^j S \right) + \int_q \left([\bar{\chi}_\pi^i]_{sr} K_{tu;rs} [\chi_\pi^j]_{ut} \right) \right]$$

$$\delta^{ij} = \frac{\partial}{\partial P^2} \left[\begin{array}{c} \text{diagram 1} + \text{diagram 2} + \text{diagram 3} \\ + \text{diagram 4} + \text{diagram 5} + \text{diagram 6} \end{array} \right]$$

lines

dashed lines

wiggles

quark propagator.

quark propagator fixed wrt. derivative.

gluon propagator.

[R. E. Cutkosky and M. Leon, Phys. Rev. **135** (1964) B1445.]

[P. Maris and P. C. Tandy, Nucl. Phys. Proc. Suppl. **161** (2006) 136.]

$$\left(\frac{d \ln(\lambda)}{dP^2} \right)^{-1} = \text{tr} \int_k \bar{\Gamma} S \Gamma S$$

$$\left(\frac{d \ln(\lambda)}{dP^2} \right)^{-1} = \left[\text{Diagram} \right]$$

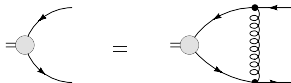
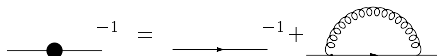
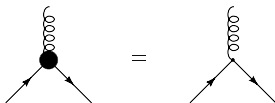
Where λ is the eigenvalue obtained via $\Gamma = \lambda K \Gamma$.

- considerably simpler
- valid for all truncations
- first time applied beyond rainbow-ladder

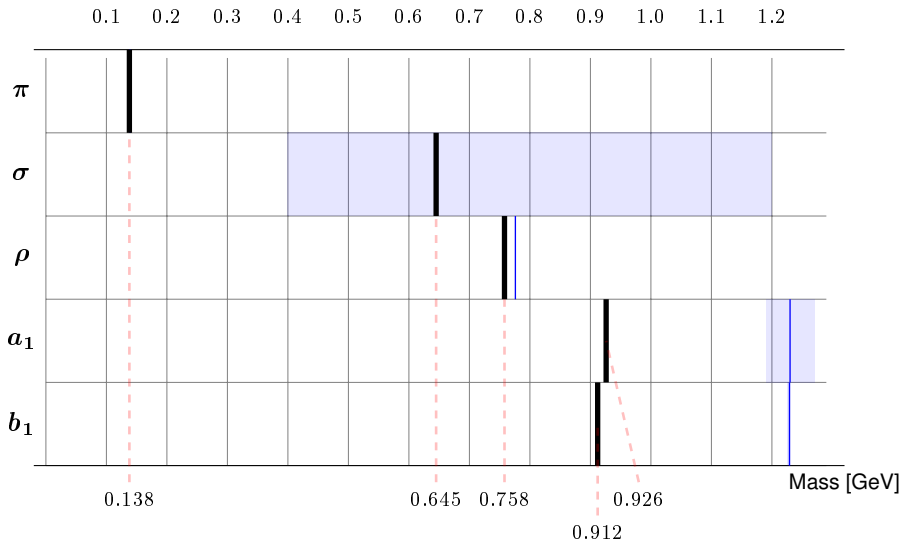
[N. Nakanishi, Phys. Rev. **138**, B1182 (1965)]

Results: Rainbow-Ladder benchmark

Calculate bound-state masses with bare vertex:



Results: Rainbow-Ladder benchmark



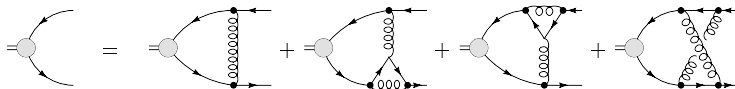
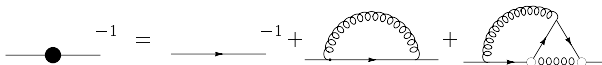
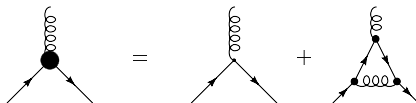
Results: Abelian corrections

Consider sub-leading Abelian corrections to vertex

[A. Bender, C. D. Roberts, L. von Smekal, Phys. Rev. Lett. B **380** (1996) 7]

[A. Bender, W. Detmold, C. D. Roberts and A. W. Thomas, Phys. Rev. C **65** (2002) 065203]

[P. Watson, W. Cassing and P. C. Tandy, Few Body Syst. **35** (2004) 129]

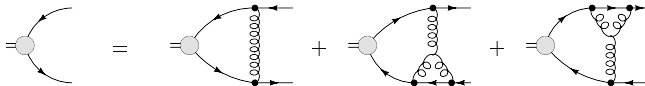
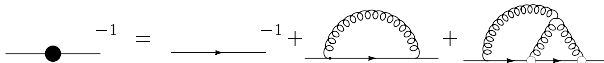
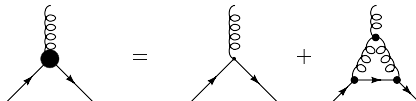


Results: Abelian corrections (unpublished)

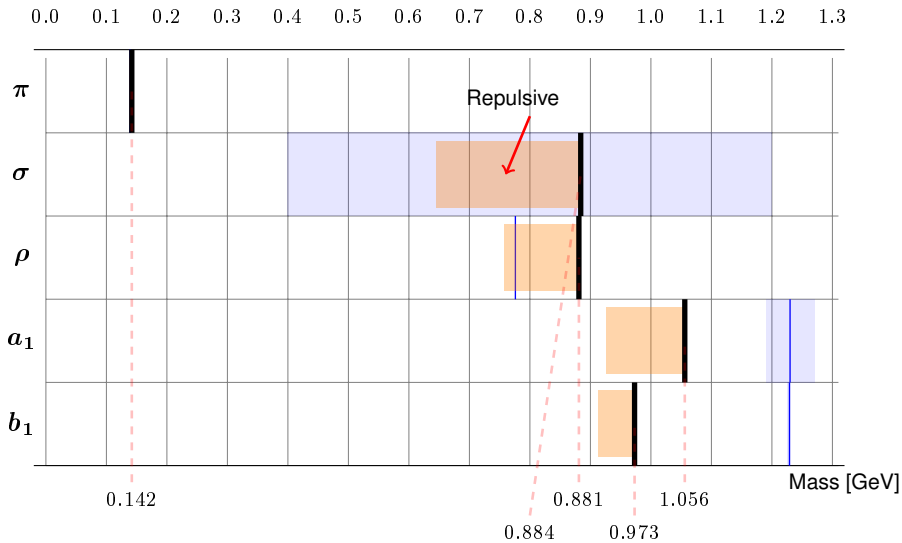


Results: non-Abelian corrections

Consider dominant non-Abelian corrections to vertex

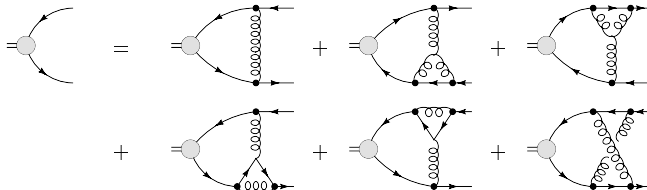
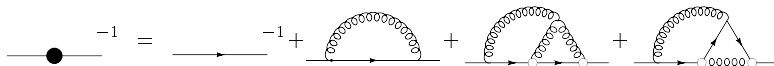
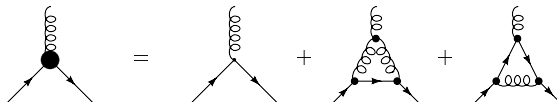


Results: non-Abelian corrections

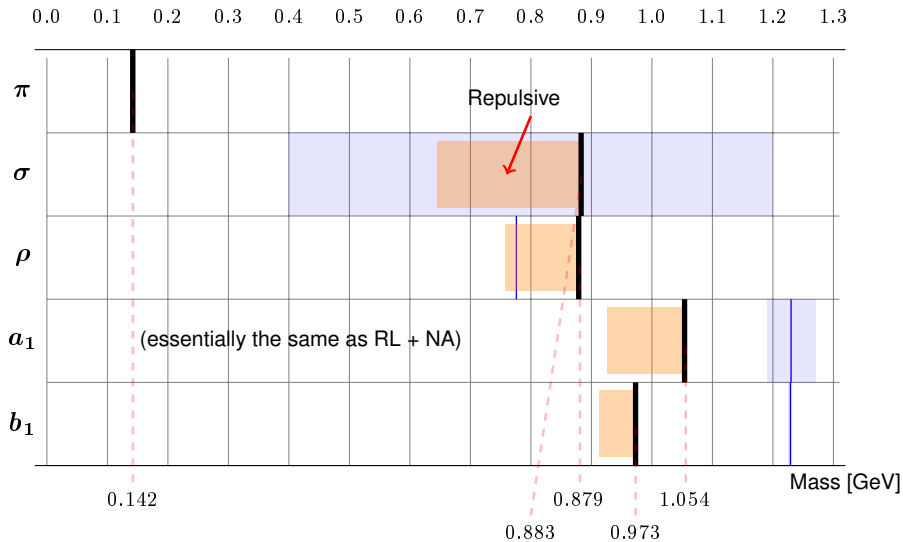


Results: non-Abelian and Abelian corrections

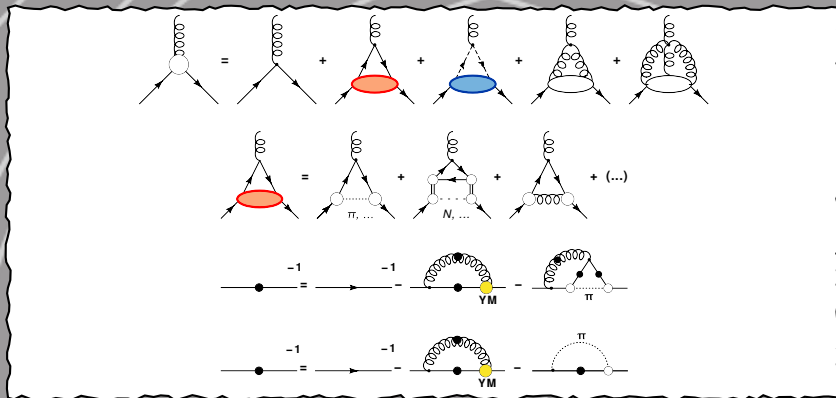
Put them all-together!



Results: non-Abelian and Abelian corrections (unpublished)



Contributions from the pion cloud



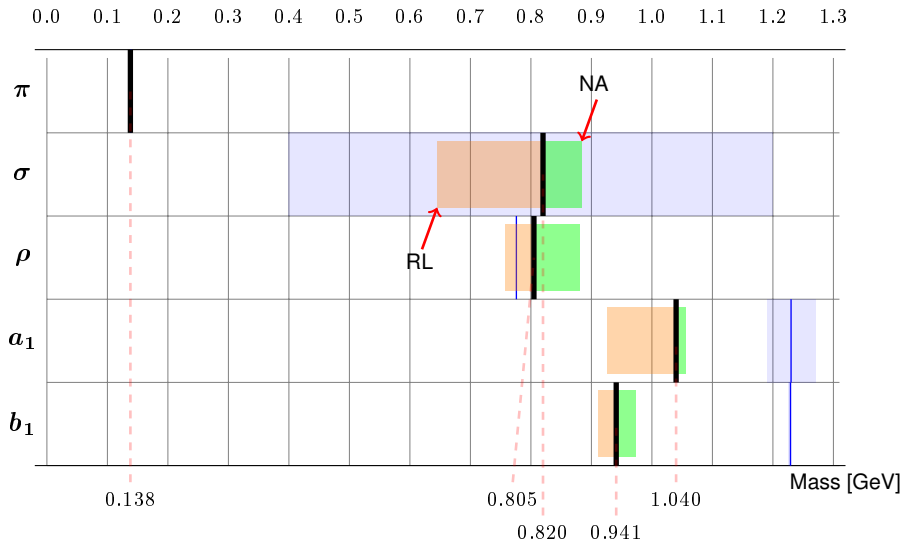
In addition to non-resonant diagrams

- resonant contributions from, *e.g.* pion exchange

→ *approximate and include*

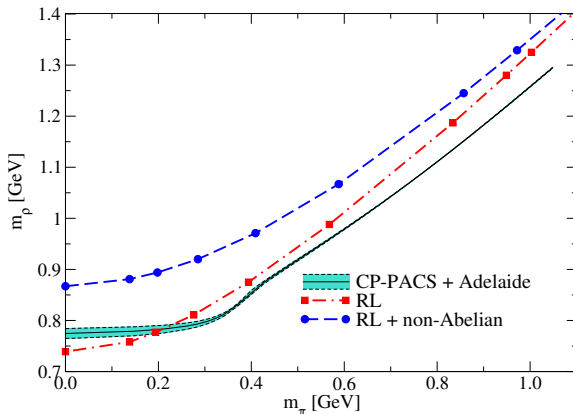
(Ignore contribution from sub-leading Abelian diagrams)

Results: non-Abelian and pion-cloud effects



Results: Mass dependence of Vector Meson

- Implicit mass-dependence of vertex corrections
- inc. pion cloud: shifts towards corrected lattice results.



[C. R. Allton, W. Armour, D. B. Leinweber, A. W. Thomas and R. D. Young, Phys. Lett. B **628** (2005) 125.]

Conclusions

Summary

Reached an era where rainbow-ladder is no longer the only feasible truncation

- first study **no trivialization of momenta**:
 - leading non-Abelian corrections
 - sub-leading Abelian corrections
 - *also* hadronic resonance contributions
- full calculation:
 - solve the full system of equations without overt simplifications
- Not just a toy-model
 - foundation on which to build future *ab initio* bound-state studies
 - just change the inputs to something realistic!

Outlook or *work in progress*

- incorporate solutions from gluon DSE.
- pion/electromagnetic form factors, decays
- do *exotics* and *diquarks* exist beyond rainbow-ladder – Bound/Unbound?
- excited states, tensor mesons, (and more quarks!).