Baryon-Baryon Interactions from Lattice QCD

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HAL QCD Collaboration

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1. Introduction

What binds protons and neutrons inside a nuclei?



gravity: too weak Coulomb: repulsive between pp no force between nn, np

New force (nuclear force)?

1935 H. Yukawa

introduced virtual particles (mesons) to explain the nuclear force





1949 Nobel prize

Modern nucleon-nucleon potential





Origin of RC: "The most fundamental problem in Nuclear physics."

Note: Pauli principle is not essential for the "RC".

A challenge in (lattice) QCD !



Plan of my talk

- 1. Introduction
- 2. Strategy in (Lattice) QCD
 - 1) Previous attempt
 - 2) Strategy
 - 3) 1st (quenched) result
- 3. Recent Developments
 - 1) Velocity Dependence
 - 2) Quark mass dependence
 - 3) Tensor potential
 - 4) Full QCD
 - 5) Hyperon-Nucleon interaction
- 4. Conclusion

2. Strategy in (Lattice) QCD

Definition of "Potential" in (lattice) QCD ?

Previous attempt

Takahashi-Doi-Suganuma, AIP Conf.Proc. 842,249(2006)

calculate energy of Qqq +Qqq as a function of r between 2Q. Q:static quark, q: light quark Qqq qqQ Qqq qqQ Qqq qqQ Qqq qqQ Q (b) (c)(d) (a) Quenched result +(b)+(c)all (a) (**κ=0.1650**) (**κ=0.1650**) (k=0.1650) 0.5 0.5 0.5 $V_{BB} - 2V_{B}$ [GeV] $V_{BB} - 2V_{B}$ [GeV] $V_{BB} - 2V_{B}$ [GeV] 0 ŢŢŢ 0 -0.5 -0.5 -0.5 0.2 0.2 0.4 0.6 0.8 0 0.4 0.6 0.8 1 0 0.2 0.4 0.6 0.8 0 Distance r [fm] Distance r [fm] Distance r [fm]

Almost no dependence on r !

Quantum Field Theoretical consideration

S-matrix below inelastic threshold

$$S = e^{2i\delta}$$

Bethe-Salpeter (BS) Wave function $\varphi_E(\mathbf{r}) = \langle 0 | N(\mathbf{x} + \mathbf{r}, 0) N(\mathbf{x}, 0) | 6q, E \rangle$ 6 quark state with energy E $N(x) = \varepsilon_{abc} q^a(x) q^b(x) q^c(x)$: local operator Asymptotic behavior $| \quad r = |{f r}| o \infty$ $\varphi_E^l(r) \longrightarrow A_l \frac{\sin(kr - l\pi/2 + \delta_l(k))}{kr} \qquad E = \frac{k^2}{2\mu_N} = \frac{k^2}{m_N}$

partial wave



 $\delta_l(k)$ is the scattering phase shift



1. Define a (non-local) potential from the BS wave function as

$$[E - H_0]\varphi_E(\mathbf{x}) = \int d^3y \, U(\mathbf{x}, \mathbf{y})\varphi_E(\mathbf{y}) \qquad H_0 = \frac{-\nabla^2}{2\mu_N}$$

2. Expand the potential as $U(\mathbf{x}, \mathbf{y}) = V(\mathbf{x}, \nabla)\delta^3(\mathbf{x} - \mathbf{y})$

$$V(\mathbf{x}, \nabla) = V_C(r) + V_T(r)S_{12} + V_{\mathrm{LS}}(r)\mathbf{L} \cdot \mathbf{S} + \{V_D(r), \nabla^2\} + \cdots$$

$$\mathsf{Dkubo-Marshak} \text{ (1958)} \qquad \mathsf{tensor operator} \quad S_{12} = \frac{3}{r^2}(\sigma_1 \cdot \mathbf{x})(\sigma_2 \cdot \mathbf{x}) - (\sigma_1 \cdot \sigma_2) \qquad r = |\mathbf{x}|$$

- 3. Successive determination using BS wave function at different E
 - similar to the expansion of EFT
- 4. Calculate observables (phase shift, binding energy etc.)

First (quenched) results



Qualitative features of NN potential are reproduced !

Ishii-Aoki-Hatsuda, PRL90(2007)0022001

This paper has been selected as one of 21 papers in Nature Research Highlights 2007

3. Recent developments

velocity dependence of V

$$V(\mathbf{x}, \nabla) = V_C(r) + V_T(r)S_{12} + V_{LS}(r)\mathbf{L} \cdot \mathbf{S} + \{V_D(r), \nabla^2\} + \cdots$$

K. Murano, S. Aoki, T. Hatsuda, N. Ishii, H. Nemura



velocity-dep. terms can be determined form E-dependence of V.
 E-dep. turns out to be small at low energy in our choice of N(x).

Quark mass dependence of V





as quark mass decreases

stronger repulsive core at short distance

stronger attraction at intermediate distance





Ishii, Aoki, Hatsuda (in progress)

Triplet (S=1) ${}^{3}S_{1}$ $J = 1, L = J \pm 1 = 2, 0$ mix with ${}^{3}D_{1}$ $(E - H_{0})\psi = [V_{C} + V_{T}S_{12}]\psi$ $P\psi \equiv \frac{1}{24}\sum_{R}\psi(R\vec{r})$ "projection" to L=0 ${}^{3}S_{1}$ Q = 1 - P "projection" to L=2 ${}^{3}D_{1}$

$$\begin{pmatrix} P\psi & PS_{12}\psi \\ Q\psi & QS_{12}\psi \end{pmatrix} \times \begin{pmatrix} V_{\rm C} \\ V_T \end{pmatrix} = (E - H_0) \begin{pmatrix} P\psi \\ Q\psi \end{pmatrix}$$



No repulsive core in tensor



Quark mass dependence



Fit: pi+rho with gaussian form-factors

Tensor forces becomes stronger as quark mass decreases.

Preliminary

Full QCD calculations







* Larger repulsive core than quenched* Larger tensor force than quenched

 $V_C(r)$ and $V_T(r)$ in full QCD ($m_{\pi} = 570$ MeV, L=2.9 fm)



Hyperon-Nucleon interactions

Almost no information from experiments

Experimetal data for ΛN

Only total cross section. No phase-shift analysis. Spin-dependence is unclear.



J-PARC

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Structure of the neutron-star core



 ΛN in full QCD ($m_{\pi} = 415$ MeV, L=2.9 fm)



Preliminary

- * Weaker repulsive core than NN
- * Stronger spin dependence than NN
- * Weaker tensor force than NN

4. Conclusion



- 1. Nuclear force from lattice QCD
 - BS wave function -> NN, NY, YY potentials -> observables
- 2. NN potential in quenched QCD: good "shape"
 - repulsive core, intermediate attraction, tensor force

"The achievement is both a computational *tour de force* and a triumph for theory."

Nature Research Highlights 2007





various channel -> input to hyper- nuclear physics

- Full QCD with 140 MeV pion is our ultimate goal.
 - current: PACS-CS config. with L=2.9 fm & pion mass = 156 -701 MeV
 - in 1-2 years: PACS-CS config. wiyj L=5.8 fm & pion mass = 140 MeV

in 5 years: new config. on 20 PFlops machine (2011-)

Current and Future target of HAL QCD

- tensor force and deuteron binding
- origin of the repulsive core
- LS force
- YN and YY forces
- 3 body forces
- light nuclei from lattice QCD
- relation to EFT





O NN force in quenched QCD:

Ishii, Aoki & Hatsuda, Phys. Rev. Lett. 99 (2007) 022001 (nucl-th/0611.096).

O Introductory review:

Aoki, Hatsuda & Ishii, Comput. Sci. Disc. I (2008) 015009 (arXiv:0805.2462[hep-lat]).

O YN force in quenched QCD:

Nemura, Ishii, Aoki & Hatsuda, Phys. Lett. B673 (2009) 136 (arXiv:0806.1094[nucl-th]).

O NN force in full QCD:

Ishii, Aoki & Hatsuda, arXiv: 0903.5497[hep-lat].

O YN force in full QCD:

Nemura, Ishii, Aoki & Hatsuda, arXiv:0902.12251[hep-lat].